

Derivation of general representation of angular momentum algebra

Let $\vec{J} = (J_x, J_y, J_z)$ be a set of hermitian operators obeying

$$[J_x, J_y] = i\hbar J_z$$

$$[J_y, J_z] = i\hbar J_x \quad (*)$$

$$[J_z, J_x] = i\hbar J_y$$

Recall $J^2 = J_x^2 + J_y^2 + J_z^2$

[earlier we saw that for $sp = \frac{1}{2}$, S^2 and S_z are compatible because $|\pm\rangle$ are eigenstates of both.]

Using (*), you showed $[J_z, J^2] = 0$

$\therefore J^2$ and J_z are compatible and have a complete set of mutual eigenstates $|\alpha, m\rangle$

$$J^2 |\alpha, m\rangle = \alpha \hbar^2 |\alpha, m\rangle$$

$$J_z |\alpha, m\rangle = m\hbar |\alpha, m\rangle$$

[NB we do not assume α, m are integers]

Let $|\alpha, m\rangle$ form an orthonormal basis

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Recall $J_{\pm} = J_x \pm iJ_y$

You showed $[J_z, J_{\pm}] = \pm \hbar J_{\pm}$

so

$$J_z (J_{\pm} |\alpha, m\rangle) = (m \pm 1) \hbar (J_{\pm} |\alpha, m\rangle)$$

i.e. $J_{\pm}$ raises/lowers the J_z eigenvalue by 1

You showed

$$\begin{aligned} [J^2, J_x] &= 0 \\ [J^2, J_y] &= 0 \\ \Rightarrow [J^2, J_{\pm}] &= 0 \end{aligned}$$

$$J^2 (J_{\pm} |\alpha, m\rangle) = \alpha \hbar^2 (J_{\pm} |\alpha, m\rangle)$$

i.e. $J_{\pm}$ does not change the J^2 eigenvalue

$$J_{\pm} |\alpha, m\rangle \sim |\alpha, m \pm 1\rangle$$

To prove: $\alpha \geq m^2$

$$J^2 = J_z^2 + J_y^2 + J_x^2$$

$$\langle \alpha, m | J^2 | \alpha, m \rangle = \langle \alpha, m | J_z^2 | \alpha, m \rangle + \langle \alpha, m | J_y^2 | \alpha, m \rangle + \langle \alpha, m | J_x^2 | \alpha, m \rangle$$

$$\alpha \hbar^2 = m^2 \hbar^2 + \quad ? \quad + \quad ??$$

Lemma: for any hermitian operator A , $\langle \psi | A^2 | \psi \rangle \geq 0$

$$\langle \psi | A^2 | \psi \rangle = \langle \psi | A^\dagger A | \psi \rangle = \langle A\psi | A\psi \rangle \geq 0$$

Proof: $\langle \psi | A^2 | \psi \rangle = \langle A\psi | A\psi \rangle$

Let $\phi = A\psi$ then $\langle \psi | A^2 | \psi \rangle = \langle \phi | \phi \rangle$

$\langle \phi | \phi \rangle \geq 0$ (norm squared)

J_x and J_y are hermitian so ? and ?? are ≥ 0

$$\Rightarrow \alpha \geq m^2$$

Hence J_+ , J_- cannot operate indefinitely

$\exists$ a maximal value of m , call it $j \Rightarrow J_+ | \alpha, j \rangle = 0$

and a minimum value, call it $\bar{j} \Rightarrow J_- | \alpha, \bar{j} \rangle = 0$

$$\# \text{ of states} = j - \bar{j} + 1$$

$$\begin{array}{c} J_+ \\ \uparrow \\ | \alpha, j \rangle \\ J_- \downarrow \\ | \alpha, j-1 \rangle \\ \vdots \\ | \alpha, \bar{j} \rangle \\ J_- \downarrow \end{array}$$

This finite set of states is a ~~the~~ representation of angular momentum algebra

(m 's not necessarily integers)

We previously showed $J^2 = J_z^2 + \hbar J_z + J_- J_+$

Act on top state

$$J^2 |\alpha, j\rangle = (J_z^2 + \hbar J_z + J_- J_+) |\alpha, j\rangle$$

$$\alpha \hbar^2 |\alpha, j\rangle = (j(j+1)\hbar^2 + 0) |\alpha, j\rangle$$

$$\Rightarrow \alpha = j(j+1)$$

We also showed $J^2 = J_z^2 - \hbar J_z + J_+ J_-$

Act on lowest state

$$\alpha \hbar^2 |\alpha, \bar{j}\rangle = ((\bar{j}\hbar)^2 - \hbar(+\bar{j}\hbar) + 0) |\alpha, \bar{j}\rangle$$

$$\alpha = \bar{j}(\bar{j}-1)$$

$$j(j+1) = \bar{j}(\bar{j}-1)$$

$$j^2 - \bar{j}^2 + j + \bar{j} = 0$$

$$(j+\bar{j})(j-\bar{j}) = 0$$

$$(j-\bar{j})(j-\bar{j}+1) = 0$$

$$|\alpha, j\rangle$$

$$\bar{j} = \begin{cases} -j \\ j+1 \end{cases} \rightarrow \text{became } \bar{j} < j$$

$$\# \text{ of states} = j - \bar{j} + 1 = 2j + 1 \in \mathbb{Z}$$

j can be either
integer (orbital angular momentum / spin of boson)

half-integer (spin of fermion)

"Spin j representation"

~~Most general state of the rep~~

$$\sim 2j+1$$

Change in notation:

instead of $|\alpha, m\rangle$ w $\alpha = j(j+1)$

we write $|j, m\rangle$ where $j = m_{\max}$

$$J^2 |j, m\rangle = j(j+1) \hbar^2 |j, m\rangle$$

$$J_z |j, m\rangle = m \hbar |j, m\rangle$$

Assuming that $|j, m\rangle$ is an orthonormal basis,

$$J_+ |j, m\rangle = c_+(j, m) |j, m+1\rangle$$

$$J_- |j, m\rangle = c_-(j, m) |j, m-1\rangle$$

(Problem: calculate $c_+(j, m)$ and $c_-(j, m)$)

$$\left[\begin{array}{l} \text{Ans: } c_+ = \sqrt{j(j+1) - m(m+1)} \quad \hbar = \sqrt{(j-m)(j+m+1)} \hbar \\ c_- = \sqrt{j(j+1) - m(m-1)} \quad \hbar = \sqrt{(j+m)(j-m+1)} \hbar \end{array} \right]$$

Most general state of spin j is $|\psi\rangle = \sum_{m=-j}^j a_m |j, m\rangle$

has well-defined J^2 but not J_z