

[After they've done the problem]

33-1

Harmonic oscillator summary

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2 \quad \text{where } \omega = \sqrt{\frac{k}{m}}$$

operator approach (Dirac)

$$\begin{cases} a = \sqrt{\frac{m\omega}{2\hbar}} x + \frac{i}{\sqrt{2\hbar m\omega}} p \\ a^\dagger = \sqrt{\frac{m\omega}{2\hbar}} x - \frac{i}{\sqrt{2\hbar m\omega}} p \end{cases}$$

$$\Rightarrow [a, a^\dagger] = 1$$

$$\Rightarrow H = \hbar\omega (N + \frac{1}{2}1) \quad \text{where } N = a^\dagger a.$$

Let $|n\rangle$ be eigenstates of N w/ eigenvalue n
+ therefore eigenstates of H w/ eigenvalue $\hbar\omega(n + \frac{1}{2})$

You showed $a^\dagger + a$ are raising + lowering operators.

$$a|n\rangle = \sqrt{n}|n-1\rangle$$

$$a|n\rangle = \sqrt{n+1}|n+1\rangle$$

remember: $\sqrt{\text{larger integer}}$

$$\Rightarrow a|0\rangle = 0$$

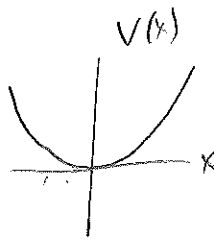
$$N|0\rangle = 0|0\rangle$$

$$H|0\rangle = \frac{1}{2}\hbar\omega|0\rangle$$

grd st energy

$$|n\rangle = \frac{1}{\sqrt{n!}} (a^\dagger)^n |0\rangle$$

Harmonic oscillator
(eg diatomic molecule)



$$E = \frac{p^2}{2m} + \frac{1}{2} k_s x^2$$

Natural frequency $\omega = \sqrt{\frac{k_s}{m}} \Rightarrow E = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 \hat{x}^2 = \hbar \omega \left[\frac{m \omega}{2 \hbar} \hat{x}^2 + \frac{1}{2 m \hbar \omega} \hat{p}^2 \right]$$

operator approach

Define $\hat{a} = \sqrt{\frac{m \omega}{2 \hbar}} \hat{x} + \frac{i}{\sqrt{2 m \hbar \omega}} \hat{p}$

non hermitian! (not an observable)

$$\hat{a}^\dagger = \sqrt{\frac{m \omega}{2 \hbar}} \hat{x} - \frac{i}{\sqrt{2 m \hbar \omega}} \hat{p}$$

$[\hat{a}, \hat{a}^\dagger] = 1$ using $[\hat{x}, \hat{p}] = i \hbar 1$

Define $\hat{N} = \hat{a}^\dagger \hat{a} = \frac{m \omega}{2 \hbar} \hat{x}^2 + \frac{1}{2 m \hbar \omega} \hat{p}^2 - \frac{1}{2} 1 = \frac{\hat{H}}{\hbar \omega} - \frac{1}{2} 1$

$$\hat{H} = \hbar \omega \left(\hat{N} + \frac{1}{2} 1 \right)$$

$$[\hat{N}, \hat{a}^\dagger] = \hat{a}^\dagger \Rightarrow \hat{N} \hat{a}^\dagger = \hat{a}^\dagger (\hat{N} + 1)$$

$$[\hat{N}, \hat{a}] = -\hat{a} \Rightarrow \hat{N} \hat{a} = \hat{a} (\hat{N} - 1)$$

Let $\hat{N} |\lambda\rangle = \lambda |\lambda\rangle$ be eigenstates of $\hat{N}$ (normalized)

Then $\hat{N} (\hat{a}^\dagger |\lambda\rangle) = (\lambda + 1) (\hat{a}^\dagger |\lambda\rangle)$ i.e.
 $\hat{N} (\hat{a} |\lambda\rangle) = (\lambda - 1) (\hat{a} |\lambda\rangle)$ i.e.

$\hat{a}^\dagger |\lambda\rangle \sim |\lambda + 1\rangle$ raising
 $\hat{a} |\lambda\rangle \sim |\lambda - 1\rangle$ lowering

$$\begin{array}{c} \hat{a}^\dagger \uparrow |\lambda + 1\rangle \\ \hat{a} \downarrow |\lambda - 1\rangle \end{array}$$

not necessary
a done - proba

$|\lambda\rangle$ are also energy eigenstates

$$\hat{H}|\lambda\rangle = \underbrace{\hbar\omega(\lambda + \frac{1}{2})}_E |\lambda\rangle$$

From analysis of Schrodinger eqn $\Rightarrow E \geq V_{\min} = 0$
 $\Rightarrow \lambda \geq -\frac{1}{2}$

$$\exists \lambda_{\min} \ni a|\lambda_{\min}\rangle = 0$$

$$\langle \lambda_{\min} | a^\dagger = 0$$

$$0 = \langle \lambda_{\min} | a^\dagger a | \lambda_{\min} \rangle = \langle \lambda_{\min} | N | \lambda_{\min} \rangle = \lambda_{\min}$$

$$E_{\min} = \hbar\omega(\lambda_{\min} + \frac{1}{2}) = \frac{1}{2}\hbar\omega \leftarrow \text{ground state energy of harmonic oscillator}$$

$|0\rangle = \text{ground state}$

All other states obtained by raising λ by 1 $\Rightarrow |n\rangle$
 $a^\dagger |1\rangle = |2\rangle$
 $n = \text{non-neg integers}$

$$\text{Eigenstate } |n\rangle \text{ has energy } E = \hbar\omega(n + \frac{1}{2})$$

one can show:

$$a^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle$$

$$a |n\rangle = \sqrt{n} |n-1\rangle$$

mnemonic = $\sqrt{\text{larger}}$

$$|1\rangle = a^\dagger |0\rangle$$

$$|2\rangle = \frac{1}{\sqrt{2}} a^\dagger |1\rangle = \frac{1}{\sqrt{2}} a^\dagger a^\dagger |0\rangle$$

$$|3\rangle = \frac{1}{\sqrt{3}} a^\dagger |2\rangle$$

$$\Rightarrow |n\rangle = \frac{1}{\sqrt{n!}} (a^\dagger)^n |0\rangle$$

Harmonic oscillator in position space

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$$\hat{a} = \sqrt{\frac{m\omega}{2\hbar}} x + \sqrt{\frac{\hbar}{2m\omega}} \frac{d}{dx}$$

$$= \frac{1}{\sqrt{2}} \left(y + \frac{d}{dy} \right)$$

$$y = \sqrt{\frac{m\omega}{\hbar}} x$$

$$\therefore \hat{a}^\dagger = \frac{1}{\sqrt{2}} \left(y - \frac{d}{dy} \right)$$

Ground state setup $a|0\rangle = 0$

In position space $\left(y + \frac{d}{dy} \right) u_0(y) = 0$

$$\frac{du_0}{dy} = -y u_0$$

$$\frac{du_0}{u_0} = -y dy$$

$$\ln u_0 = -\frac{1}{2} y^2 + \text{const}$$

$$u_0 = C e^{-\frac{1}{2} y^2}$$

$$= C e^{-\frac{m\omega x^2}{2\hbar}}$$

$$\text{Normalized: } u_0 = \left(\frac{m\omega}{\pi\hbar} \right)^{\frac{1}{4}} e^{-\frac{m\omega x^2}{2\hbar}}$$

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^\dagger)^n |0\rangle$$

In part a)

$$\begin{aligned} u_n &= \frac{1}{\sqrt{n!}} \left(\frac{1}{\sqrt{2}} \left[y - \frac{d}{dy} \right] \right)^n \left(\frac{m\omega}{\pi\hbar} \right)^{\frac{1}{4}} e^{-\frac{1}{2}y^2} \\ &= \left(\frac{m\omega}{\pi\hbar} \right)^{\frac{1}{4}} \frac{1}{\sqrt{n!} 2^n} \underbrace{\left(y - \frac{d}{dy} \right)^n e^{-\frac{1}{2}y^2}}_{\text{this is } H_n(y) e^{-\frac{1}{2}y^2}} \end{aligned}$$

eg $\left(y - \frac{d}{dy} \right) e^{-\frac{1}{2}y^2}$

$$= \left[y - (-y) \right] e^{-\frac{1}{2}y^2} = 2y e^{-\frac{1}{2}y^2}$$

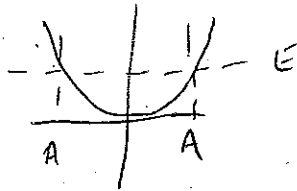
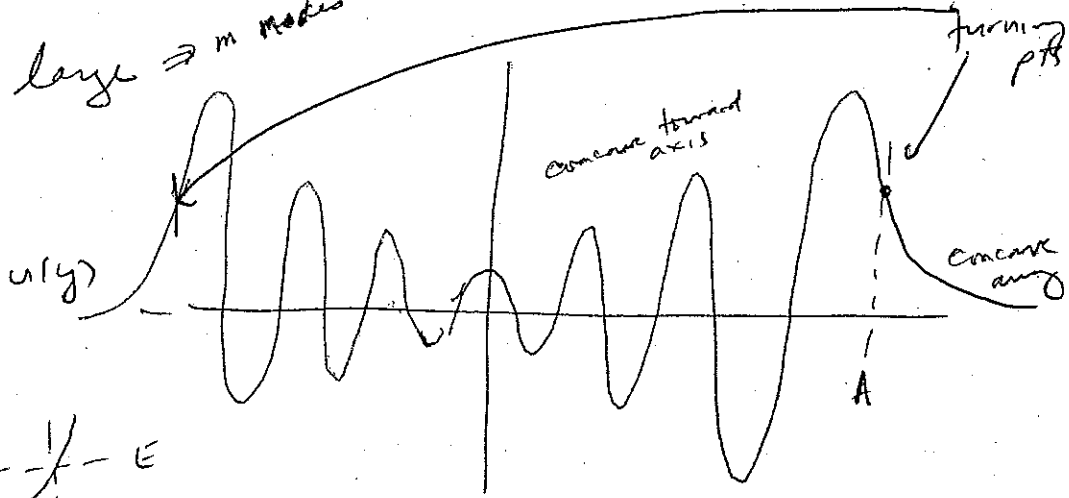
$$H_1(y) = 2y$$

n	$H_n(y)$
0	1
1	$2y$
2	$4y^2 - 2$
3	$8y^3 + 2y$
4	$16y^4 - 48y^2 + 12$

$$u_n = \left(\frac{m\omega}{\pi\hbar} \right)^{\frac{1}{4}} \frac{1}{\sqrt{n!} 2^n} H_n(y) e^{-\frac{1}{2}y^2}$$

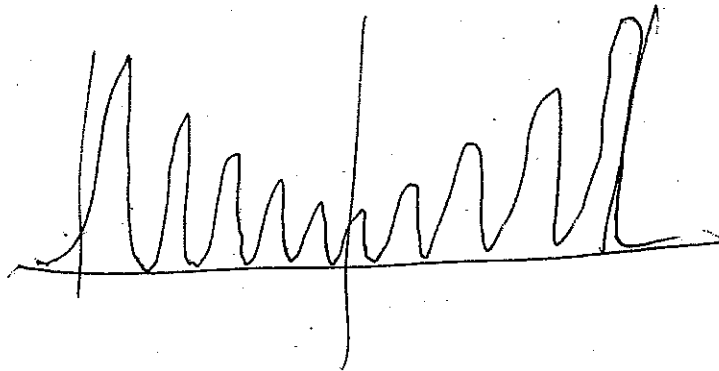
33-4
~~1st~~
 (circled)
 II
 26

$m \text{ large} \Rightarrow m \text{ modes}$



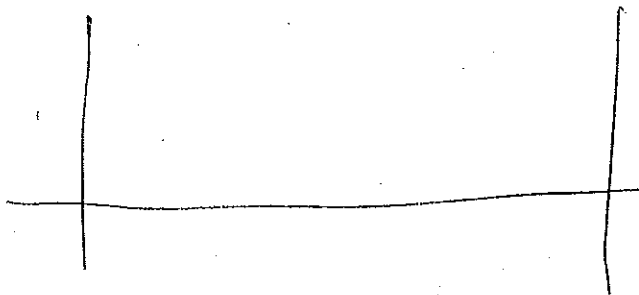
~~Waves~~ beyond turning points, $|u|^2 \neq 0$

$|u|^2$



Why is it more probable to find ~~the~~ particle near turning points? Because particle moves ~~slow~~ more slowly there!

~~As the potential increases, the probability density increases.~~



Shyp

$$\text{Prob} \sim \text{time spent} \sim \frac{1}{v}$$

point of time spent
 b/w x & $x + \Delta x$
 $\Delta t = \frac{\Delta x}{v}$

$$x = A \cos \omega t$$

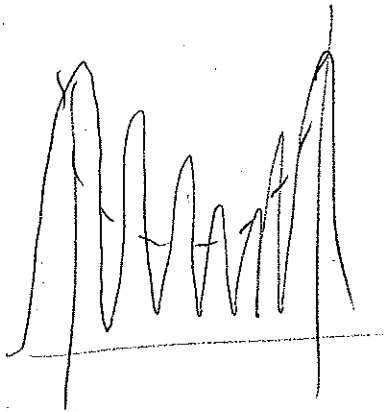
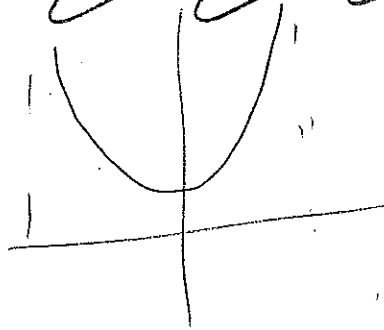
$$v = -A \omega \sin \omega t = -\omega \sqrt{A^2 - x^2}$$

$$v = -\omega \sqrt{A^2 - x^2}$$

$$\rho(x) \sim \frac{1}{v} \sim \frac{1}{\sqrt{A^2 - x^2}}$$

$$\frac{1}{2}mv^2 = E - V(x) = \frac{1}{2}m\omega^2(A^2 - x^2)$$

$$V = \pm \omega \sqrt{A^2 - x^2}$$



Om

do a problem

Correspondence finished

$$Q_m \rightarrow C_m \text{ as } n \rightarrow \infty$$

