

Momentum space wavefunction

Use completeness $\int dp |p\rangle \langle p| = 1$ to write

$$|\psi\rangle = \int dp |p\rangle \langle p|\psi\rangle$$

Define $\phi(p) = \langle p|\psi\rangle$ = wavefunction of ψ in momentum space
prob. amp.

$$|\psi\rangle = \int dp \phi(p) |p\rangle$$

$$1 = \langle\psi|\psi\rangle = \int dp \langle\psi|p\rangle \langle p|\psi\rangle = \int dp |\phi(p)|^2$$

ϕ square integrable in p -space.

$|\phi(p)|^2$ = probability density in momentum space

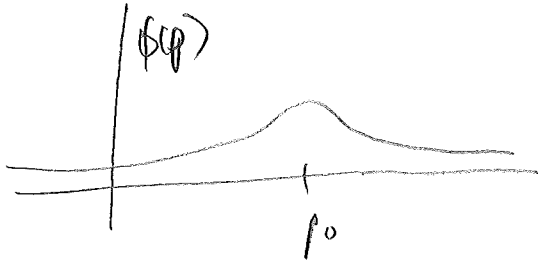
$$\langle p \rangle = \langle\psi|\hat{p}|\psi\rangle = \int dp \langle\psi|\underbrace{\hat{p}|p\rangle}_{p|p\rangle} \langle p|\psi\rangle = \int_{-\infty}^{\infty} dp |\phi(p)|^2 p$$

$|p\rangle$ is an idealization; perfectly well-defined p
 try for something more realistic: range of momenta $\rightarrow$ why spread out in position space

29-2

gaussian

Consider a particle γ momentum space distribution $\phi(p) = N e^{-\beta(p-p_0)^2}$



[expect to measure p_0 , γ some spread]

$$1 = \int dp |\phi(p)|^2 = N^2 \int_{-\infty}^{\infty} dp e^{-2\beta(p-p_0)^2} = N^2 \underbrace{\int_{-\infty}^{\infty} dp' e^{-2\beta(p')^2}}_{I(\beta)}$$

$$\Rightarrow N = \frac{1}{\sqrt{I(\beta)}}$$

$$I = \int_{-\infty}^{\infty} dx e^{-2\beta x^2}$$

$$I^2 = \int_{-\infty}^{\infty} dx e^{-2\beta x^2} \int_{-\infty}^{\infty} dy e^{-2\beta y^2} = \int \underbrace{dx dy}_{r dr d\theta} e^{-2\beta \underbrace{(x^2+y^2)}_{r^2}}$$

$$= \underbrace{\int_0^{2\pi} d\theta}_{2\pi} \underbrace{\int_0^{\infty} r e^{-2\beta r^2} dr}_{\frac{1}{4\beta} \int_0^{\infty} du e^{-u}}$$

$$u = 2\beta r^2 \\ du = 4\beta r dr$$

$$\frac{1}{4\beta} \underbrace{\int_0^{\infty} du e^{-u}}_1$$

$$= \frac{\pi}{2\beta}$$

$$I = \sqrt{\frac{\pi}{2\beta}}$$

$$\Rightarrow \phi(p) = \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} e^{-\beta(p-p_0)^2}$$

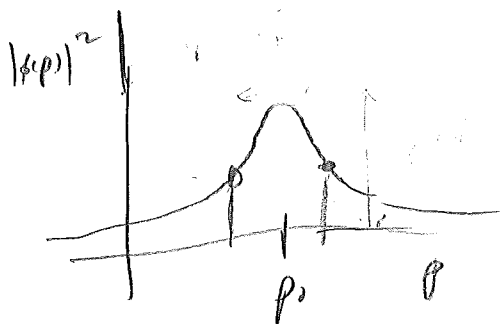
$$\begin{aligned}
\langle p \rangle &= \int dp \, |\psi(p)|^2 p \\
&= \left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dp \, e^{-2\beta(p-p_0)^2} p \quad \text{let } p = p_0 + p' \\
&= \left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dp' \, e^{-2\beta p'^2} (p_0 + p') \\
&= p_0 \left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dp' \, e^{-2\beta p'^2} + \underbrace{\left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dp' \, e^{-2\beta p'^2} p'}_{\text{odd integrand}} \\
&= p_0
\end{aligned}$$

$$\begin{aligned}
\langle p^2 \rangle &= \int_{-\infty}^{\infty} dp \, |\psi(p)|^2 p^2 \\
&= \left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dp' \, e^{-2\beta p'^2} (p_0^2 + 2p_0 p' + p'^2) \\
&= p_0^2 + 0 + \underbrace{\left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} \int_{-\infty}^{\infty} dp' \, e^{-2\beta p'^2} p'^2}_{\text{even integrand}}
\end{aligned}$$

$$I(\beta) = \int dp' \, e^{-2\beta p'^2} = \sqrt{\frac{\pi}{2\beta}} \quad \longrightarrow \quad -\frac{1}{2} \frac{dI}{d\beta} = \frac{1}{4} \sqrt{\frac{\pi}{2}} \frac{1}{\beta^{3/2}}$$

$$\langle p^2 \rangle = p_0^2 + \frac{1}{4\beta}$$

[problem: $\langle p^{2n} \rangle$]

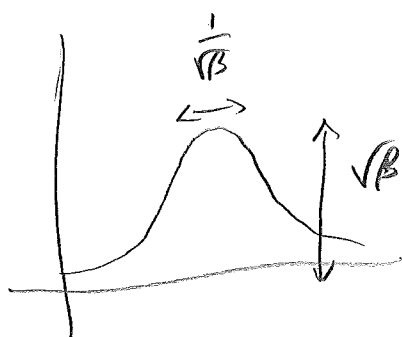


$$|\psi(p)|^2 = \left(\frac{2\beta}{\pi}\right)^{\frac{1}{2}} e^{-2\beta(p-p_0)^2}$$

$$= e^{-1}$$

$$2\beta(p-p_0)^2 = 1$$

$$|p-p_0| = \frac{1}{\sqrt{2\beta}}$$



max spread: $\Delta p = \sqrt{\langle p^2 \rangle - \langle p \rangle^2} = \frac{1}{\sqrt{2\beta}}$

$$\begin{aligned}
 |\psi\rangle &= \int dx \, \psi(x) |x\rangle \\
 &= \int dp \, \phi(p) |p\rangle
 \end{aligned}$$

How are $\phi(p)$ & $\psi(x)$ related? Fourier transforms!

$$\psi(x) = \langle x | \psi \rangle = \int dp \, \phi(p) \langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} \int dp \, \phi(p) e^{i\frac{px}{\hbar}}$$

$\psi(x)$ = superposition of plane waves weighted by $\phi(p)$

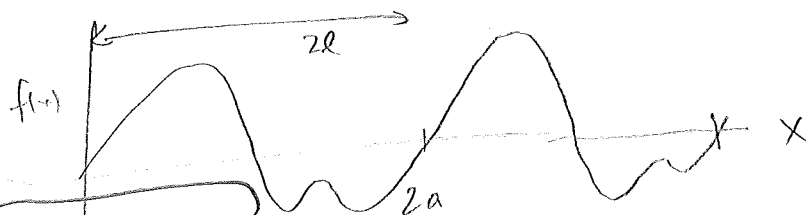
$$\phi(p) = \langle p | \psi \rangle = \int dx \, \psi(x) \langle p | x \rangle = \frac{1}{\sqrt{2\pi\hbar}} \int dx \, \psi(x) e^{-i\frac{px}{\hbar}}$$

Quick review

Periodic function

$$f(x+2a) = f(x)$$

These are part of Griffiths problem
I then assign part of it



Some information

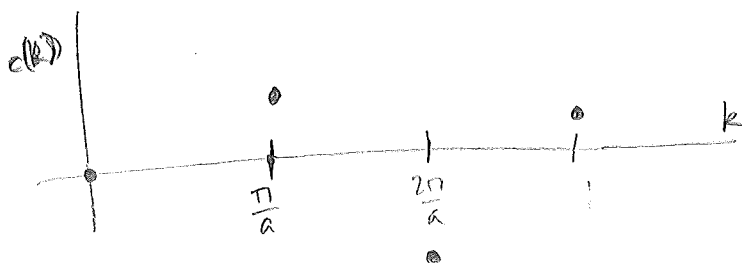
Recall from Phys 360

$$f(x) = \sum (a_n \cos \frac{n\pi x}{a} + b_n \sin \frac{n\pi x}{a})$$

(or a, b complex since γ complex) $i k x$

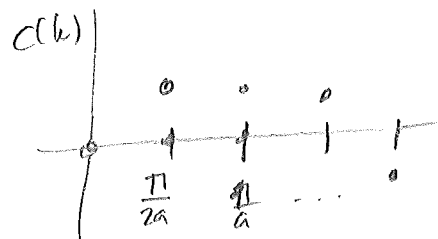
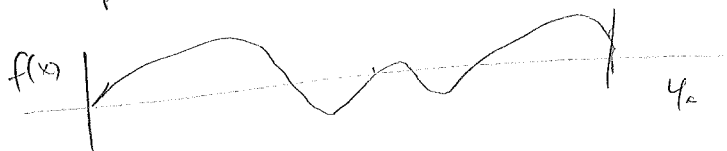
$$f(x) = \sum_{-\infty}^{\infty} c_n e^{i k_n x} = \sum_{k \in \frac{\pi}{a} \mathbb{Z}} c(k) e^{i k x}$$

[Fourier series]



[2nd plot is discrete but goes to ∞ ; first is continuous but limited to $0 \leq x \leq 2a$]

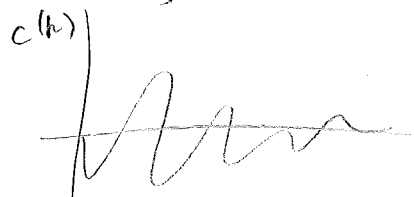
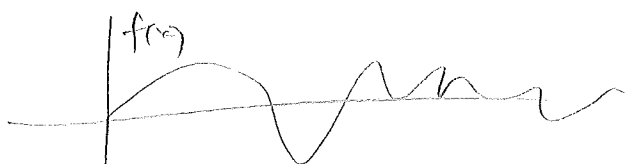
Double period



twice as much into each.

Let $a \rightarrow \infty$

Fourier transform



Can calc $\langle p \rangle$ using either $\phi(p)$ or $\psi(x)$

$$\langle p \rangle = \langle \psi | \hat{p} | \psi \rangle = \int dp \phi^*(p) p \phi(p)$$

Now Fourier transform

$$= \int dp \phi^*(p) p \frac{1}{\sqrt{2\pi\hbar}} \int dx \psi(x) e^{-ipx/\hbar}$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \int dp \phi^*(p) \int_{-\infty}^{\infty} dx \psi(x) \underbrace{p e^{-ipx/\hbar}}_{-\frac{\hbar}{i} \frac{d}{dx} e^{-ipx/\hbar}}$$

$$\int u \frac{dv}{dx} = \underbrace{\int dx \frac{d}{dx}(uv)}_{uv} - \int dx \frac{du}{dx} v$$

$$\underbrace{-\frac{\hbar}{i} \psi(x) e^{-ipx/\hbar}}_0 - \int dx \left(-\frac{\hbar}{i}\right) \frac{d\psi}{dx} e^{-ipx/\hbar}$$

$$= \int_{-\infty}^{\infty} dx \frac{\hbar}{i} \frac{d\psi}{dx} \underbrace{\frac{1}{\sqrt{2\pi\hbar}} \int dp \phi^*(p) e^{-ipx/\hbar}}_{\psi^*(x)}$$

$$= \int_{-\infty}^{\infty} dx \psi^*(x) \left(\frac{\hbar}{i} \frac{d}{dx}\right) \psi(x)$$

$$\langle A \rangle = \int_{-\infty}^{\infty} dx \psi^*(x) \hat{A}_{pos.} \psi(x)$$

in position space $\left\{ \begin{array}{l} \hat{p}_{pos.} \\ \hat{p}_{pos} \end{array} \right\} \quad \frac{\hbar}{i} \frac{d}{dx}$

(Note: show $\hat{p}$ is hermitian)

$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int dp e^{i p x / \hbar} \phi(p)$$

what is the position space wavefunction of a gaussian momentum distrib.?

$$\phi(p) = \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} e^{-\beta(p-p_0)^2} \Rightarrow \Delta p = \frac{1}{2\sqrt{\beta}}$$

$$\psi(x) = \left(\frac{1}{2\pi\hbar}\right)^{\frac{1}{2}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \int dp e^{-\beta(p-p_0)^2 + i p x / \hbar}$$

Let $q = p - p_0 \Rightarrow p = q + p_0$

$$e^{i p_0 x / \hbar} \int dq e^{-\beta q^2 + i q x / \hbar}$$

Consider $\int dq e^{-aq^2 + bq} = \int dq e^{-a(q - \frac{b}{2a})^2 + \frac{b^2}{4a}}$

$$= e^{\frac{b^2}{4a}} \int dq' e^{-aq'^2}$$

$$= \sqrt{\frac{\pi}{a}} e^{b^2/4a} \quad \checkmark$$

$$\psi(x) = \left(\frac{1}{2\pi\hbar}\right)^{\frac{1}{2}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \left(\frac{\pi}{\beta}\right)^{\frac{1}{2}} e^{\frac{1}{4\beta} \left(\frac{x}{\hbar}\right)^2} e^{i p_0 x / \hbar}$$

$$= \left(\frac{1}{2\pi\hbar^2\beta}\right)^{\frac{1}{4}} e^{-\frac{x^2}{4\beta\hbar^2}} e^{i p_0 x / \hbar}$$

$$= \left(\frac{2\alpha}{\pi}\right)^{\frac{1}{4}} e^{-\alpha x^2} e^{i p_0 x / \hbar} \quad \text{where } \alpha = \frac{1}{4\beta\hbar^2}$$

$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int dp e^{ipx/\hbar} \phi(p)$$

Alternative calculation,
but not as suitable for
generalization to time dependent
case $\rightarrow$ see
with
prob

What is the position space wavefunction of a
given momentum distribution?

$$\phi(p) = \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} e^{-\beta(p-p_0)^2}$$

$$\Delta p = \frac{1}{2\sqrt{\beta}}$$

$$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \int dp e^{-\beta p^2 + 2\beta p_0 p - \beta p_0^2 + \frac{ipx}{\hbar}}$$

Consider $\int_{-\infty}^{\infty} dp e^{-ap^2 + bp + c}$

$$= \int_{-\infty}^{\infty} dp e^{-a(p - \frac{b}{2a})^2 + \frac{b^2}{4a} + c}$$

$$= e^{\frac{b^2}{4a} + c} \int_{-\infty}^{\infty} dp' e^{-a(p')^2}$$

$$= e^{\frac{b^2}{4a} + c} \sqrt{\frac{\pi}{a}}$$

Then

$$\psi(x) = \left(\frac{1}{2\pi\hbar}\right)^{\frac{1}{2}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \left(\frac{\pi}{\beta}\right)^{\frac{1}{2}} e^{\frac{1}{4\beta} \left(\frac{ix}{\hbar} + 2\beta p_0\right)^2 - \beta p_0^2}$$

$$= \left(\frac{1}{2\pi\hbar^2\beta}\right)^{\frac{1}{4}} e^{-\frac{x^2}{4\beta\hbar^2} + \frac{ip_0 x}{\hbar}}$$

$$= \left(\frac{2\alpha}{\pi}\right)^{\frac{1}{4}} e^{-\alpha x^2} e^{\frac{ip_0 x}{\hbar}} \quad \text{where } \alpha \equiv \frac{1}{4\beta\hbar^2}$$

Trying to do the travelling Gaussian % doing the shift.
(this is more subtle + not suitable for this)

$$\psi(x,t) = \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \int dp e^{-\beta(p-p_0)^2 + \frac{ipx}{\hbar} - \frac{ip^2 t}{2m\hbar}}$$

$$= \frac{1}{\sqrt{2\pi\hbar}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \int dp e^{-\beta p^2} e^{-\left[\beta + \frac{it}{2m\hbar}\right] p^2 + \left[\frac{ix}{\hbar} + 2\beta p_0\right] p}$$

$$= \left(\frac{1}{2\pi\hbar}\right)^{\frac{1}{2}} \left(\frac{2\beta}{\pi}\right)^{\frac{1}{4}} \left(\frac{\pi}{\beta + \frac{it}{2m\hbar}}\right)^{\frac{1}{2}} e^{-\beta p_0^2 + \frac{\left(\frac{ix}{\hbar} + 2\beta p_0\right)^2}{4\left(\beta + \frac{it}{2m\hbar}\right)}}$$

this is not helpful

$$e^{\frac{1}{4\left(\beta + \frac{it}{2m\hbar}\right)} \left[-\frac{x^2}{\hbar^2} + \frac{4\beta p_0 i x}{\hbar} + 4\beta^2 p_0^2 - 4\beta p_0^2 \left(\beta + \frac{it}{2m\hbar}\right) - \frac{2\beta p_0^2 i t}{m\hbar} \right]}$$

$$e^{\frac{1}{4\left(\beta + \frac{it}{2m\hbar}\right)} \left[-\frac{x^2}{\hbar^2} + \frac{4\beta p_0 i}{\hbar} \left(x - \frac{p_0}{2m} t\right) \right]}$$

$$e^{-\beta p_0^2 + \frac{\left(\frac{ix}{\hbar} + 2\beta p_0\right)^2 \left(\beta - \frac{it}{2m\hbar}\right)}{4\left(\beta^2 + \frac{t^2}{4m^2\hbar^2}\right)}}$$

$$\left(\frac{ix}{\hbar} + 2\beta p_0\right)^2 \left(\beta - \frac{it}{2m\hbar}\right) = \beta \left(-\frac{x^2}{\hbar^2} + 4\beta^2 p_0^2\right) + \left(\frac{4\beta p_0 i x}{\hbar}\right) \left(\frac{t}{2m\hbar}\right) + (\text{imag part})$$

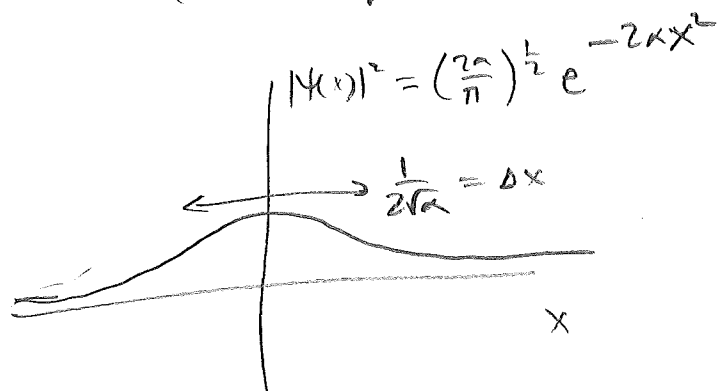
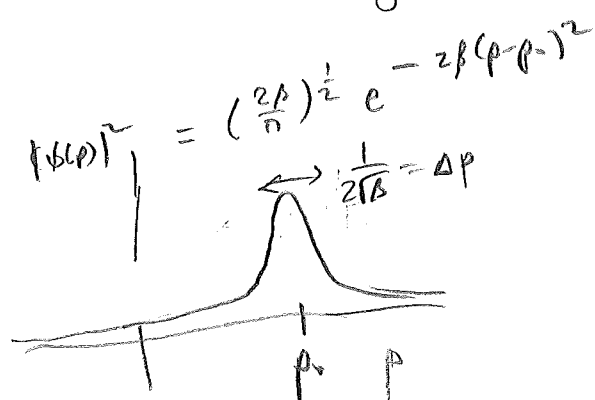
$$= -\frac{\beta}{\hbar^2} \left(x - \frac{p_0}{m} t\right)^2 + \frac{\beta p_0^2 t^2}{m^2 \hbar^2} + 4\beta^3 p_0^2$$

$$= -\frac{\beta \left(x - \frac{p_0}{m} t\right)^2}{4\hbar^2 \left(\beta^2 + \frac{t^2}{4m^2\hbar^2}\right)} + i(\text{stuff})$$

so e

Carols off first piece

Fourier of a gaussian is a gaussian (time a phase)

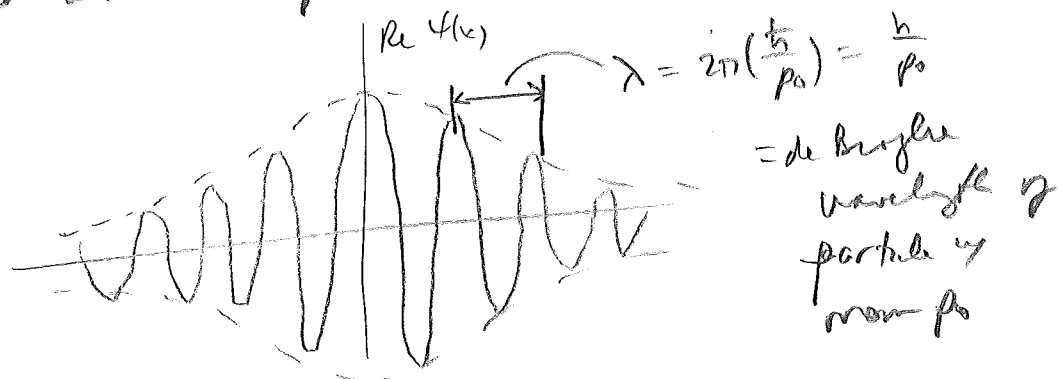


Since $\alpha\beta = \frac{1}{4\hbar^2}$

Products of widths $\left(\frac{1}{2\alpha}\right)\left(\frac{1}{2\beta}\right) = \frac{\hbar}{2}$ widths are inverse related

$\Delta x \Delta p = \frac{\hbar}{2}$ Heisenberg is saturated

What about the phase? $\text{Re } \psi(x) = \left(\frac{2\alpha}{\pi}\right)^{\frac{1}{4}} e^{-\alpha x^2} \cos\left(\frac{p_0 x}{\hbar}\right)$



[wavelength not perfectly well defined because momentum not "]