

Backgrd

21-0

$$e^{-\frac{i\epsilon p}{\hbar}} X e^{\frac{i\epsilon p}{\hbar}} = X + \frac{i\epsilon}{\hbar} [X, P] = X - \epsilon$$

$$\langle x_0 | e^{-\frac{i\epsilon p}{\hbar}} X e^{\frac{i\epsilon p}{\hbar}} | x_0 \rangle = x_0 - \epsilon$$

$$\Rightarrow e^{\frac{i\epsilon p}{\hbar}} | x_0 \rangle = | x_0 - \epsilon \rangle$$

$$\text{Now } |\psi\rangle = \int dx \psi(x) |x\rangle$$

$$e^{\frac{i\epsilon p}{\hbar}} |\psi\rangle = \int dx \psi(x) e^{\frac{i\epsilon p}{\hbar}} |x\rangle$$

$$= \int dx \psi(x) |x - \epsilon\rangle$$

$$= \int dx \psi(x + \epsilon) |x\rangle$$

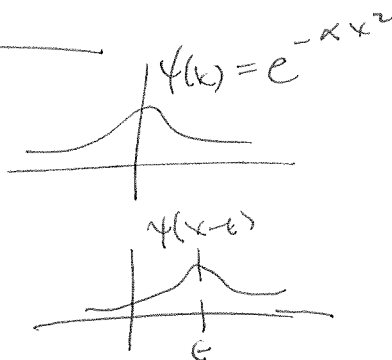
So acting on $\psi(x)$ we have

$$e^{\frac{i\epsilon p}{\hbar}} \psi(x) = \psi(x + \epsilon)$$

$$e^{\frac{i\epsilon}{\hbar} \frac{d}{dx}} \psi(x) = \psi(x + \epsilon)$$

$$e^{-\frac{i\epsilon p}{\hbar}} \psi(x) = \psi(x - \epsilon)$$

$$e^{-\frac{i\epsilon p}{\hbar}} |x\rangle = |x_0 + \epsilon\rangle$$



$U = e^{-\frac{i\epsilon p}{\hbar}}$ shifts state to the right by ϵ

Recall that a spatial rotation $\vec{\omega}$ is implemented by

unitary operator $U = e^{-\frac{i\vec{\omega} \cdot \vec{J}}{\hbar}}$ where $\vec{J}$ = angular momentum operator
 i.e. $\vec{J}$ generates spatial rotation

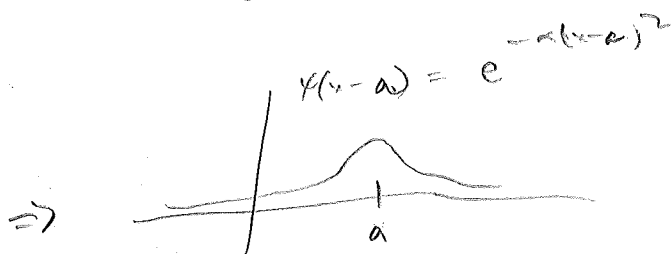
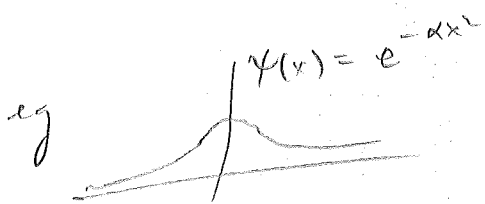
Recall that a time translation by t is implemented by

unitary operator $U = e^{-\frac{i\hat{H}t}{\hbar}}$ where $\hat{H}$ = Hamiltonian = energy operator
 i.e. $\hat{H}$ generates time translation

A spatial translation by a is implemented by

unitary operator $U = e^{-\frac{ia\hat{P}}{\hbar}}$ where $\hat{P}$ = momentum operator
 i.e. $\hat{P}$ generates spatial translation

U acts on a ^{position space} wavefunction $\psi(x)$ to give $\psi(x-a)$



$$U \psi(x) = \psi(x-a)$$

$$\begin{aligned} \text{But } \psi(x-a) &= \psi(x) - a \frac{d\psi}{dx} + \frac{1}{2!} a^2 \frac{d^2\psi}{dx^2} - \dots \\ &= e^{-a \frac{d}{dx}} \psi(x) \end{aligned}$$

$$\text{So } -a \frac{d}{dx} = -\frac{i}{\hbar} a \hat{P} \Rightarrow$$

$$\boxed{\hat{P} = \frac{\hbar}{i} \frac{d}{dx} \text{ on wavefunctions}}$$

$$\begin{array}{l}
 \hat{x} \text{ acts on } \psi(x) \text{ to give } x\psi(x) \\
 \hat{p} \text{ acts on } \psi(x) \text{ to give } \frac{\hbar}{i} \frac{d}{dx} \psi(x) \\
 \cancel{\hat{x} \text{ acts on } \psi(x) \text{ to give } x\psi(x)}
 \end{array}
 \left. \vphantom{\begin{array}{l} \hat{x} \text{ acts on } \psi(x) \text{ to give } x\psi(x) \\ \hat{p} \text{ acts on } \psi(x) \text{ to give } \frac{\hbar}{i} \frac{d}{dx} \psi(x) \end{array}} \right\} \text{ both are linear operations}$$

Consider $[\hat{x}, \hat{p}] \psi(x) = x p \psi - p x \psi$

$$\begin{aligned}
 &= x \frac{\hbar}{i} \frac{d\psi}{dx} - \frac{\hbar}{i} \frac{d}{dx} (x\psi) \\
 &\quad \underbrace{\psi + x \frac{d\psi}{dx}} \\
 &= i\hbar \psi(x)
 \end{aligned}$$

Since this is true for arbitrary $\psi(x)$, we find

$$[\hat{x}, \hat{p}] = i\hbar \mathbb{1}$$

Heisenberg commutation relations

x, p incompatible.

Recall generalized uncertainty relation

$$\Delta A \Delta B \geq \frac{1}{2} |\langle \psi | [A, B] | \psi \rangle|$$

$$\Rightarrow \Delta x \Delta p \geq \frac{\hbar}{2}$$

x, p incompatible $\Rightarrow$ eigenstates of x are not eigenstates of p .