

Position observable X

Review of what has gone before:

observable $A \rightarrow$ operator \hat{A}

$A|a_n\rangle = a_n|a_n\rangle$ eigenstates = states of definite A

$$\sum_n |a_n\rangle \langle a_n| = \mathbb{1} \quad \text{completeness}$$

$$|\psi\rangle = \sum |a_n\rangle \langle a_n|\psi\rangle =$$

define $\psi_n = \langle a_n|\psi\rangle = \text{prob. amplitude}$
 $|a_n\rangle = \psi_n |a_n\rangle \rightarrow \text{indeterminate w.r.t. } A$

$$|\psi\rangle = \sum \psi_n |a_n\rangle \quad \sum \psi_n^* \psi_n = \sum |\psi_n|^2 = 1$$

$$\langle \psi | \psi \rangle = \sum \psi_n \langle a_n | \psi \rangle = \sum \psi_n^* \psi_n = 1$$

$|\psi_n|^2 = \text{probability of measuring } a_n$

$$\langle A \rangle = \langle \psi | \hat{A} | \psi \rangle = \sum_n \underbrace{\langle \psi | \hat{A} | a_n \rangle}_{a_n | a_n \rangle} \frac{\langle a_n | \psi \rangle}{\psi_n} = \sum_n a_n \psi_n \checkmark$$

$$\langle A \rangle = \langle \psi | \hat{A} | \psi \rangle = \sum_n \underbrace{\langle \psi | \hat{A} | a_n \rangle}_{a_n | a_n \rangle} \frac{\langle a_n | \psi \rangle}{\psi_n} = \sum_n |\psi_n|^2 a_n \checkmark$$

operator \hat{x}

$$\hat{x}|x\rangle = x|x\rangle$$

operator x = possible result of a measurement of position
= any real number

$|x\rangle$ = state of perfectly well-defined position (idealized)

Completeness (since all values of x are allowed, $\Sigma \rightarrow []$)

$$\int dx |x\rangle \langle x| = \mathbb{I}$$

$$|\psi\rangle = \int dx |x\rangle \langle x|\psi\rangle$$

Define $\psi(x) = \langle x|\psi\rangle = \text{a number that depends on } x \text{ and } \psi$
or a function x !

$$|\psi\rangle = \int dx \psi(x) |x\rangle = \text{state of indeterminate position}$$

$\psi(x)$ = probability amplitude, or wavefunction

Suppose $|f\rangle$ is normalized

$$\begin{aligned} 1 = \langle f | f \rangle &= \int dx \langle f(x) \rangle \langle x | f \rangle \\ &= \int dx f^*(x) f(x) \\ &= \int dx |f(x)|^2 \end{aligned}$$

$|f(x)|^2$ = probability $\underline{\text{of}} \rightarrow$ obtaining x

A function obeying $\int dx |f(x)|^2$ is called
"square integrable" or "normalizable"

$$f(x) \xrightarrow[x \rightarrow \pm\infty]{} 0 \quad \text{for the} \quad \frac{1}{\sqrt{|x|}}$$

How does \hat{x} act on a wavefunction $\psi(x)$?

$$\text{Let } |\psi\rangle = \int dx \psi(x) |x\rangle$$

$$\hat{x}|\psi\rangle = \int dx \psi(x) \hat{x}|x\rangle = \int dx \psi(x) x |x\rangle$$

so if $|\psi\rangle \rightarrow \psi(x)$ then $\hat{x}|\psi\rangle \rightarrow x\psi(x)$

i.e. \hat{x} multiplies $\psi(x)$ by x

$$\langle x \rangle = \langle \psi | \hat{x} | \psi \rangle = \int dx \underbrace{\langle \psi | \hat{x} | x \rangle}_{\langle \psi | x \rangle} \underbrace{\langle x | \psi \rangle}_{\psi^*(x)}$$

$$= \int dx \psi^*(x) x \psi(x)$$

[as in 2140]

For seq: $\hat{x} = \int dx \int dx' |x\rangle \underbrace{\langle x | x' | x' \rangle}_{x f(x-x')} \langle x |$

$$= \int dx |x\rangle x \langle x|$$

i.e. diagonal

What about orthonormality of $\langle \cdot | \cdot \rangle$?

Define $\langle x' | x \rangle = \delta(x', x)$, a func of x, x'

If $x \neq x'$, then orthogonal $\Rightarrow \delta(x', x) = 0$

what is $\delta(x, x) = ?$

Consider

$$\psi(x') = \langle x' | \int dx \psi(x) | x \rangle = \int_{-\infty}^{\infty} dx \psi(x) \delta(x', x)$$

- integral vanishes if $x \neq x'$.
- nonzero on set of measure zero.
- $\delta(x', x)$ must be as at $x=x'$.

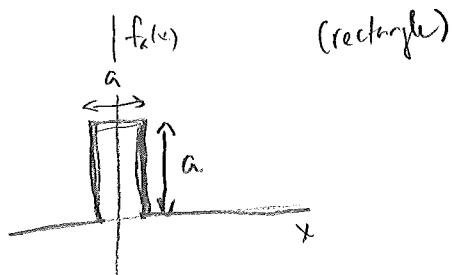
Define Dirac delta function $\delta(x)$ [really a distribution]

$$\text{and that } \begin{cases} \textcircled{1} & \delta(x) = 0 \quad \text{if } x \neq 0 \\ \textcircled{2} & \int_{-\infty}^{\infty} \delta(x) dx = 1 \end{cases}$$

$\delta(x)$ is usually represented as a limit of a family of functions

i.e.

$$\delta(x) = \lim_{a \rightarrow 0} R_a(x), \quad R_a(x) = \begin{cases} 0, & |x| > \frac{1}{2}a \\ \frac{1}{a}, & |x| < \frac{1}{2}a \end{cases}$$



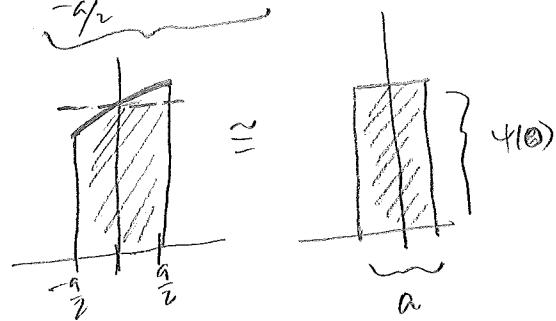
① Let $x_0 \neq 0$, then $\lim_{a \rightarrow 0} R_a(x_0) = 0$ [choose $a < 2|x_0|$]

② $\int_{-\infty}^{\infty} R_a(x) dx = 1$ $\therefore \int_{-\infty}^{\infty} \lim_{a \rightarrow 0} R_a(x) dx = 1$ ✓

Let $\psi(x)$ be an arbitrary function

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$$\text{Consider } \int_{-\infty}^{\infty} R_a(x) \psi(x) dx = \frac{1}{a} \int_{-\frac{a}{2}}^{\frac{a}{2}} \psi(x) dx \stackrel{?}{=} \psi(0)$$



$$\lim_{a \rightarrow 0} \int_{-\infty}^{\infty} R_a(x) \psi(x) dx = \psi(0)$$

$$\int_{-\infty}^{\infty} \delta(x) \psi(x) dx = \psi(0)$$

more generally

$$\int_{-\infty}^{\infty} \delta(x-x') \psi(x) dx = \psi(x')$$

[This is exactly what we want]

$$\langle x' | x \rangle = \delta(x-x')$$

$$\langle x' | x \rangle = \delta(x-x') = \delta(x'-x)$$

"Dirac delta for normalization"

[use for
since $\langle x' | x \rangle = \langle x | x' \rangle$]

(NB \hat{x} eigenstates not normalizable $\langle x | x \rangle \neq 1$)

Claim: position operator hermitian. How prove?

Hermite conjugation

[or std Hermitian problem?]

$$\text{Recall } \langle \psi | A^\dagger | \phi \rangle = \langle \phi | A | \psi \rangle^*$$

In posit space representation

$$\int d\omega \psi^*(\omega) (\hat{A}^\dagger)_{\text{pos.}} \phi(\omega) = \left[\int d\omega \psi^*(\omega) (\hat{A})_{\text{pos.}} \psi(\omega) \right]^*$$

Now

$$\int d\omega \psi^*(\omega) (\hat{x}^\dagger)_{\text{pos.}} \phi(\omega) = \left(\int d\omega \psi^*(\omega) \times \psi(\omega) \right)^*$$

$$= \int d\omega \psi^*(\omega) \times^* \phi(\omega)$$

$$\Rightarrow (\hat{x}^\dagger)_{\text{pos.}} = \times^* = (\hat{x})_{\text{pos.}}$$

$$\boxed{\hat{x}^\dagger = \hat{x}}$$

\hat{x} expectation in position space

$$\langle x_0 \rangle = \int dx \langle x | x_0 \rangle |x\rangle = \int dx \delta(x-x_0) |x\rangle$$

not normalizable

$$\int dx |\psi|^2 = \int dx \delta(x-x_0) \delta(x-x_0) = \delta(0)$$

Alt,

$$\hat{x} |x_0\rangle = x_0 |x_0\rangle$$

$$\Rightarrow \langle x | \hat{x} | x_0 \rangle \langle x | x_0 \rangle = x_0 \langle x | x_0 \rangle$$

$$\text{or } x \langle x | x_0 \rangle = x_0 \langle x | x_0 \rangle$$

$$(x-x_0) \langle x | x_0 \rangle = 0 \quad \Rightarrow \quad \langle x | x_0 \rangle = \delta(x-x_0)$$

$$\hat{x} = \int_{\text{all } x} dx' |x\rangle \underbrace{\langle x | \hat{x} | x' \rangle}_{x' \delta(x-x')} \langle x' | = \int_{\text{all } v} dv |x\rangle \times \langle x |$$

as diagonal operator

$$\langle \psi | \psi \rangle = 1 \Rightarrow |\psi\rangle \text{ unitless}$$

$|\psi\rangle$ has units $\frac{1}{\sqrt{\text{mass}}}$

$$\int |\psi|^2 dx = 1 \Rightarrow \psi \text{ has units } \frac{1}{\sqrt{L}}$$

$$\langle x | \psi \rangle = \frac{1}{\sqrt{\text{mass}}} e^{i k x}$$

$$\int \delta(x-x') dx = 1 \Rightarrow \delta(x) \text{ has units } \frac{1}{L}$$

has units $\frac{1}{\sqrt{L \cdot \text{mass}}}$

$$\langle x | x' \rangle = \delta(x-x') \Rightarrow |x\rangle \text{ has units } \frac{1}{\sqrt{L}}$$

x has units $L \cdot \text{mass}$

$$\psi(x) = \langle x | \psi \rangle \text{ has units } \frac{1}{\sqrt{L}}$$

$$|\psi\rangle = \underbrace{\int_{-\infty}^{\infty} \frac{1}{\sqrt{L}} \frac{1}{\sqrt{L}}}_{\text{L}} \psi(x) |x\rangle \sim \text{unitless}$$

$$(\hat{A})_{pos} = \langle x | \hat{A} | x' \rangle = A_{xx'}$$

$$\hat{A} = \oint dx dx' |x\rangle A_{xx'} \langle x'|$$

$$\hat{x}_{xx'} = x \delta(x - x')$$

$$\hat{p}_{xx'} = \langle x | \hat{p} | x' \rangle = \int dp dp' \langle x | p \rangle p \frac{\delta(p - p')}{\hbar} \langle p | x' \rangle$$

$$= \frac{1}{2\pi\hbar} \int dp p e^{\frac{i(x-x')p}{\hbar}}$$

$$= \frac{\hbar}{c} \frac{d}{dx} \left[\frac{1}{2\pi\hbar} \int dp e^{\frac{i(x-x')p}{\hbar}} \right]$$

$$= \frac{\hbar}{c} \delta'(x - x')$$

$$\hat{p} = \int dx dx' |x\rangle \frac{\hbar}{c} \delta'(x - x') \langle x' |$$

$$\hat{A} = \sum |E_m\rangle A_{mn} \langle E_n|$$

$$A_{mn} = \langle E_m | \hat{A} | E_n \rangle$$

✓

Eigenfunction of \hat{X} in momentum space:

Solve
for
exam

$$\psi_{x_0}(x) = \delta(x - x_0)$$

$$\phi_{x_0}(p) = \frac{1}{\sqrt{2\pi\hbar}} \int_{-\infty}^{\infty} dx \delta(x - x_0) e^{-\frac{ipx}{\hbar}} = \frac{1}{\sqrt{2\pi\hbar}} e^{-\frac{ipx_0}{\hbar}}$$

①

$$\hat{X} \phi_{x_0}(p) = i\hbar \frac{d\phi}{dp} = p_0 \phi$$

$$\Rightarrow \phi(p) \sim e^{-\frac{ipx_0}{\hbar}}$$