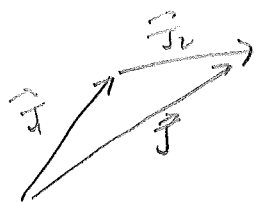


Addition of angular momentum

Classically, the angular momentum $\vec{J}$ of a particle is described by a vector.

The angular momentum of two particles is described by vector addition: $\vec{J} = \vec{J}_1 + \vec{J}_2$



$$|\vec{J}_1| - |\vec{J}_2| \leq |\vec{J}| \leq |\vec{J}_1| + |\vec{J}_2|$$

assume $|\vec{J}_2| < |\vec{J}_1|$



~~Quantum mechanics: $\vec{J} = \vec{J}_1 + \vec{J}_2$, $J_z = J_{1z} + J_{2z}$, $J^2 = J_1^2 + J_2^2$ when $\vec{J} = \vec{J}_1 + \vec{J}_2$, $\vec{J}_1 + \vec{J}_2 = \vec{J}$, $J_z = J_{1z} + J_{2z}$, $J^2 = J_1^2 + J_2^2$~~
~~misleading because what is J_z ? Do $(J_z)^2 = J_z^2$...~~

relevant ops: $J_1^2, J_{1z}, J_2^2, J_{2z}, J^2, J_z = J_{1z} + J_{2z}$

Quantum mechanically, the angular momentum of a particle is described by a state in a $(2j+1)$ dimensional space

$$\sum_{m=-j}^j a_m |j, m\rangle$$

basis of space

The angular momentum of two particles is described by a state in a $(2j_1+1) \times (2j_2+1)$ dim'l product space

$$\sum_{m_1=-j_1}^{j_1} \sum_{m_2=-j_2}^{j_2} a_{m_1 m_2} |j_1, m_1\rangle \otimes |j_2, m_2\rangle$$

basis of product space

This product space can be decomposed into a sum of spaces of the total angular momentum

$$(\text{spin } j_1) \otimes (\text{spin } j_2)$$

$$= \text{spin}(j_1+j_2) \oplus \text{spin}(j_1+j_2-1) \oplus \dots \oplus \text{spin}|j_1-j_2|$$

Do counting: $2 \oplus 2 = 2+1$

$2 \oplus 2 = 2+2$

$\vec{J}_1, \vec{J}_2, \vec{J}$ are operators that act on these spaces

$$\vec{J} = \vec{J}_1 + \vec{J}_2$$

$$\begin{aligned} J_z &= J_{1z} + J_{2z} \\ J_x &= J_{1x} + J_{2x} \Rightarrow J_- = J_{1-} + J_{2-} \\ J_y &= J_{1y} + J_{2y} \Rightarrow J_+ = J_{1+} + J_{2+} \end{aligned}$$

Relevant operators: $J_1^2, J_{1z}, J_2^2, J_{2z}, J^2, J_z$

Example: $j_1=1, j_2=2$ so $(\text{spin } 1) \otimes (\text{spin } 2) = (\text{spin } 3) \oplus (\text{spin } 2) \oplus (\text{spin } 1)$

(spin 1) space spanned by $|j_1, m_1\rangle = |1, 1\rangle, |1, 0\rangle, |1, -1\rangle$
 (spin 2) space spanned by $|j_2, m_2\rangle = |2, 2\rangle, |2, 1\rangle, |2, 0\rangle, |2, -1\rangle, |2, -2\rangle$

degenerate
 $J_1^2, J_{1z}, J_2^2, J_{2z}$

product space is 15-dimensional
 $|j_1, m_1\rangle |j_2, m_2\rangle$ or $|m_1, m_2\rangle$ for short
 $\Rightarrow 15$ basis states, labelled by $m = m_1 + m_2$
 (label rep by their multiplicity)
 $\tilde{3} \otimes \tilde{5} = \tilde{7} \oplus \tilde{5} \oplus \tilde{3} \leftarrow \text{notation}$
 label rep by their multiplicity
 eg J_1^2, J_2^2

		$ m_1, m_2\rangle$		$m = m_1 + m_2$	$ j, m\rangle$		
m_2	2	$ 1, 2\rangle$	$ 0, 2\rangle$	3	$ 3, 3\rangle$		
	1	$ 1, 1\rangle$	$ 0, 1\rangle$	2	$ 3, 2\rangle$	$ 2, 2\rangle$	
	0	$ 1, 0\rangle$	$ 0, 0\rangle$	1	$ 3, 1\rangle$	$ 2, 1\rangle$	$ 1, 1\rangle$
	-1	$ 1, -1\rangle$	$ 0, -1\rangle$	0	$ 3, 0\rangle$	$ 2, 0\rangle$	$ 1, 0\rangle$
	-2	$ 1, -2\rangle$	$ 0, -2\rangle$	-1	$ 3, -1\rangle$	$ 2, -1\rangle$	$ 1, -1\rangle$
m_1	1	$ 1, 1\rangle$	$ 0, 1\rangle$	-2	$ 3, -2\rangle$	$ 2, -2\rangle$	
	0	$ 1, 0\rangle$	$ 0, 0\rangle$	-3	$ 3, -3\rangle$		

The states in any row on the right are linear combinations of states in the same row on the left (and vice versa)

$$|j, m\rangle = \sum_{m_1, m_2} C_{m_1, m_2, m}^{j_1 j_2 j} |m_1, m_2\rangle$$

↑
Clebsch-Gordan coefficients

eg $|3, 1\rangle = C_{101}^{123} |1, 0\rangle + C_{011}^{123} |0, 1\rangle + C_{-1, 2, 1}^{123} |1, -2\rangle$

Simple example

$$g_1 = \frac{1}{2}, g_2 = \frac{1}{2}$$

(eg. 2 electron spin-1)
or electron + proton

$$(\text{spin } \frac{1}{2}) \otimes (\text{spin } \frac{1}{2}) = (\text{spin } 1) \oplus (\text{spin } 0)$$

$$\underline{\underline{2}} \otimes \underline{\underline{2}} = \underline{\underline{3}} \oplus \underline{\underline{1}}$$

$$\underline{\underline{m_1, m_2}}$$

$$|\frac{1}{2}, \frac{1}{2}\rangle$$

$$|+\frac{1}{2}, -\frac{1}{2}\rangle \quad |-\frac{1}{2}, +\frac{1}{2}\rangle$$

$$|-\frac{1}{2}, -\frac{1}{2}\rangle$$

$$m = m_1 + m_2$$

$$=$$

linear combinations $\rightarrow$

$$=$$

$$|j, m\rangle$$

$$|1, 1\rangle$$

$$|1, 0\rangle$$

$$|0, 0\rangle$$

$$|1, -1\rangle$$

There are eigenstates
of $J_1^2, J_{1z}, J_2^2, J_{2z}$

There are also eigenstates of J^2

$$\text{Because } J^2 = J_{1z} + J_{2z}$$

but they are not (necessarily)
eigenstates of J^2

$$(\text{because } [J^2, J_{1z}] \neq 0 \\ [J^2, J_{2z}] \neq 0)$$

There are eigenstates
of J^2 and J_z

They are not (necessarily)
eigenstates of J_{1z} and J_{2z}

$$(\text{because } [J^2, J_{1z}] \neq 0 \\ [J^2, J_{2z}] \neq 0)$$

Clearly $|1, 1\rangle = |\frac{1}{2}; \frac{1}{2}\rangle$

and $|1, -1\rangle = |-\frac{1}{2}; -\frac{1}{2}\rangle$

but what linear combinations comprise $|1, 0\rangle$ and $|0, 0\rangle$?

Recall $J_- |j, m\rangle = \sqrt{(j+m)(j-m+1)} |j, m-1\rangle$ (ignore $\hbar$)

so $J_- |1, 1\rangle = \sqrt{2} |1, 0\rangle$

Also $J_- |\frac{1}{2}; \frac{1}{2}\rangle = (J_{1-} + J_{2-}) |\frac{1}{2}; \frac{1}{2}\rangle$

$\uparrow$ acts on 1st label $\uparrow$ acts on 2nd

$= |-\frac{1}{2}; \frac{1}{2}\rangle + |\frac{1}{2}; -\frac{1}{2}\rangle$

$\therefore |1, 0\rangle = \frac{1}{\sqrt{2}} |-\frac{1}{2}; \frac{1}{2}\rangle + \frac{1}{\sqrt{2}} |\frac{1}{2}; -\frac{1}{2}\rangle$ (NB. automatically normalized)

$\uparrow$ CG coeffs $\nearrow$

$\rightarrow$ see chart.

Triplet "

$$\begin{aligned}
 |1, 1\rangle &= |\frac{1}{2}; \frac{1}{2}\rangle \\
 |1, 0\rangle &= \frac{1}{\sqrt{2}} (|\frac{1}{2}; -\frac{1}{2}\rangle + |-\frac{1}{2}; \frac{1}{2}\rangle) \\
 |1, -1\rangle &= |-\frac{1}{2}; -\frac{1}{2}\rangle
 \end{aligned}$$

all symmetric under $m_1 \leftrightarrow m_2$



[Since cannot describe ground state of helium
but can describe excited state $\Rightarrow$ orthohelium]

$\rightarrow$ describes p + e or

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What about $|0;0\rangle$? we know $J_+|0;0\rangle=0$

Consider $|\psi\rangle = \alpha|\frac{1}{2};-\frac{1}{2}\rangle + \beta|\frac{1}{2};\frac{1}{2}\rangle$

$$J_+|\psi\rangle = (J_{1+} + J_{2+})|\psi\rangle$$

$$= \alpha|\frac{1}{2};\frac{1}{2}\rangle + \beta|\frac{1}{2};\frac{1}{2}\rangle = 0 \Rightarrow \beta = -\alpha$$

Normalise $\Rightarrow \alpha = \frac{1}{\sqrt{2}}$

singlet $|0;0\rangle = \frac{1}{\sqrt{2}} \left(\underbrace{|\frac{1}{2};\frac{1}{2}\rangle - |\frac{1}{2};-\frac{1}{2}\rangle}_{\text{antisymmetric under } m_1 \leftrightarrow m_2} \right)$
[ground state of helium or excited state (para-helium)]

N.B. $|0;0\rangle$ is orthogonal to $|1;0\rangle$. Could have used this to determine $\alpha + \beta$.

Prob: find CG coeffs
for $3d^2 = 4G^2$

~~stop here~~
~~do it properly?~~

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$$J^2 = (\vec{J}_1 + \vec{J}_2) \cdot (\vec{J}_1 + \vec{J}_2) = J_1^2 + J_2^2 + 2\vec{J}_1 \cdot \vec{J}_2$$

$$= J_1^2 + J_2^2 + 2(J_{1z}J_{2z} + \underbrace{J_{1x}J_{2x} + J_{1y}J_{2y}}_{\left(\frac{J_{1+}+J_{1-}}{2}\right)\left(\frac{J_{2+}+J_{2-}}{2}\right) + \left(\frac{J_{1+}-J_{1-}}{2i}\right)\left(\frac{J_{2+}-J_{2-}}{2i}\right)})$$

$$= J_1^2 + J_2^2 + 2J_{1z}J_{2z} + J_{1+}J_{2-} + J_{1-}J_{2+} \quad \left(\begin{array}{l} \text{did up} \\ \text{to here} \\ \text{in 2013} \end{array} \right)$$

Apply this to the triplet states

$$\begin{aligned} J^2 \left| \frac{1}{2}; \frac{1}{2} \right\rangle &= (J_1^2 + J_2^2 + 2J_{1z}J_{2z} + J_{1+}J_{2-} + J_{1-}J_{2+}) \left| \frac{1}{2}; \frac{1}{2} \right\rangle \\ &= \left(\frac{3\hbar^2}{4} + \frac{3\hbar^2}{4} + 2\left(\frac{\hbar}{2}\right)\left(\frac{\hbar}{2}\right) + 0 + 0 \right) \left| \frac{1}{2}; \frac{1}{2} \right\rangle \\ &= 2\hbar^2 \left| \frac{1}{2}; \frac{1}{2} \right\rangle \end{aligned}$$

so indeed $\left| \frac{1}{2}; \frac{1}{2} \right\rangle = |1, 0\rangle$

$$J^2 \left| \frac{1}{2}; -\frac{1}{2} \right\rangle = \underbrace{\left(\frac{3\hbar^2}{4} + \frac{3\hbar^2}{4} + 2\left(\frac{\hbar}{2}\right)\left(-\frac{\hbar}{2}\right) \right)}_{\hbar^2} \left| \frac{1}{2}; -\frac{1}{2} \right\rangle + 0 + \hbar^2 \left| -\frac{1}{2}; \frac{1}{2} \right\rangle$$

$$J^2 \left| -\frac{1}{2}; \frac{1}{2} \right\rangle = \hbar^2 \left| -\frac{1}{2}; -\frac{1}{2} \right\rangle + \hbar^2 \left| \frac{1}{2}; -\frac{1}{2} \right\rangle + 0$$

so neither $\left| \frac{1}{2}; -\frac{1}{2} \right\rangle$ nor $\left| -\frac{1}{2}; \frac{1}{2} \right\rangle$ are eigenstates of J^2

$$J^2 \left(\left| \frac{1}{2}; -\frac{1}{2} \right\rangle + \left| -\frac{1}{2}; \frac{1}{2} \right\rangle \right) = 2\hbar^2 \left(\left| \frac{1}{2}; -\frac{1}{2} \right\rangle + \left| -\frac{1}{2}; \frac{1}{2} \right\rangle \right) \Rightarrow |1, 0\rangle$$

$$J^2 \left(\left| \frac{1}{2}; -\frac{1}{2} \right\rangle - \left| -\frac{1}{2}; \frac{1}{2} \right\rangle \right) = 0 \Rightarrow |0, 0\rangle$$

Problem

Calc $|\frac{3}{2}, m\rangle, |\frac{1}{2}, m\rangle$

in $(\text{spin } 1) \otimes (\text{spin } \frac{1}{2})$

$$J_- |j, m\rangle = \sqrt{j(j+1) - m(m-1)} |j, m-1\rangle$$

$$J_+ |j, m\rangle = \sqrt{j(j+1) - m(m+1)} |j, m+1\rangle$$

$$|\frac{3}{2}, \frac{3}{2}\rangle = |1; \frac{1}{2}\rangle$$

$$\sqrt{3} |\frac{3}{2}, \frac{1}{2}\rangle = \sqrt{2} |0; \frac{1}{2}\rangle + |1; -\frac{1}{2}\rangle$$

$$2\sqrt{3} |\frac{3}{2}, -\frac{1}{2}\rangle = 2|-1; \frac{1}{2}\rangle + \sqrt{2}|0; -\frac{1}{2}\rangle + \sqrt{2}|0; -\frac{1}{2}\rangle$$

$$6 |\frac{3}{2}, -\frac{3}{2}\rangle = 2|-1; -\frac{1}{2}\rangle + 2|-1; -\frac{1}{2}\rangle + 2|-1; -\frac{1}{2}\rangle$$

$$\begin{cases} |\frac{3}{2}, \frac{3}{2}\rangle = |1; \frac{1}{2}\rangle \\ |\frac{3}{2}, \frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |0; \frac{1}{2}\rangle + \sqrt{\frac{1}{3}} |1; -\frac{1}{2}\rangle \\ |\frac{3}{2}, -\frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |0; -\frac{1}{2}\rangle + \sqrt{\frac{1}{3}} |-1; \frac{1}{2}\rangle \\ |\frac{3}{2}, -\frac{3}{2}\rangle = |-1; -\frac{1}{2}\rangle \end{cases}$$

$$|\frac{1}{2}, \frac{1}{2}\rangle = \alpha |0; +\frac{1}{2}\rangle + \beta |1; -\frac{1}{2}\rangle$$

$$J_+ |\frac{1}{2}, \frac{1}{2}\rangle = \alpha\sqrt{2} |1; +\frac{1}{2}\rangle + \beta |1; \frac{1}{2}\rangle = 0 \Rightarrow \beta = -\sqrt{2}\alpha$$

$$\Rightarrow |\frac{1}{2}, \frac{1}{2}\rangle = \sqrt{\frac{1}{3}} |0; \frac{1}{2}\rangle - \sqrt{\frac{2}{3}} |1; -\frac{1}{2}\rangle$$

$$|\frac{1}{2}, -\frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |-1; \frac{1}{2}\rangle + \sqrt{\frac{1}{3}} |0; -\frac{1}{2}\rangle - \frac{2}{\sqrt{3}} |0; -\frac{1}{2}\rangle$$

$$= \sqrt{\frac{2}{3}} |-1; \frac{1}{2}\rangle - \frac{1}{\sqrt{3}} |0; -\frac{1}{2}\rangle$$

$$\begin{cases} |\frac{1}{2}, \frac{1}{2}\rangle = \sqrt{\frac{1}{3}} |0; +\frac{1}{2}\rangle - \sqrt{\frac{2}{3}} |1; -\frac{1}{2}\rangle \\ |\frac{1}{2}, -\frac{1}{2}\rangle = -\sqrt{\frac{1}{3}} |0; -\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |-1; \frac{1}{2}\rangle \end{cases}$$

$$J_- |1,1\rangle = \sqrt{2} |1,0\rangle$$

$$J_- |1,0\rangle = \sqrt{2} |1,-1\rangle$$

$$J^2 = J_1^2 + J_2^2 + 2 J_{1z} J_{2z} + J_{1+} J_{2-} + J_{1-} J_{2+}$$

$$J^2 |1, \frac{1}{2}\rangle = 2 + \frac{3}{4} + 2(1)(\frac{1}{2}) = \frac{15}{4} = \frac{3}{2}(\frac{3}{2} + 1) \quad \checkmark$$

$$J^2 |\frac{3}{2}, \frac{1}{2}\rangle = (\sqrt{\frac{2}{3}} |0, \frac{1}{2}\rangle + \frac{1}{\sqrt{3}} |1, -\frac{1}{2}\rangle)$$

$$= (2 + \frac{3}{4}) |\frac{3}{2}, \frac{1}{2}\rangle + \underbrace{\sqrt{\frac{2}{3}} (\sqrt{2} |1, -\frac{1}{2}\rangle) + \frac{1}{\sqrt{3}} (2(1)(-\frac{1}{2}) |1, -\frac{1}{2}\rangle + \sqrt{2} |0, \frac{1}{2}\rangle)}$$

$$\frac{1}{\sqrt{3}} |1, -\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |0, \frac{1}{2}\rangle$$

$$|\frac{3}{2}, \frac{1}{2}\rangle$$

$$= \frac{15}{4} |\frac{3}{2}, \frac{1}{2}\rangle \quad \checkmark$$

$$J^2 |\frac{1}{2}, \frac{1}{2}\rangle = J^2 (\sqrt{\frac{1}{3}} |0, \frac{1}{2}\rangle - \sqrt{\frac{2}{3}} |1, -\frac{1}{2}\rangle)$$

$$= (2 + \frac{3}{4}) |\frac{1}{2}, \frac{1}{2}\rangle + \underbrace{\sqrt{\frac{1}{3}} (\sqrt{2} |1, -\frac{1}{2}\rangle) - \sqrt{\frac{2}{3}} (2(1)(-\frac{1}{2}) |1, -\frac{1}{2}\rangle + \sqrt{2} |0, \frac{1}{2}\rangle)}$$

$$2\sqrt{\frac{2}{3}} |1, -\frac{1}{2}\rangle - 2\frac{1}{\sqrt{3}} |0, \frac{1}{2}\rangle$$

$$-2 |\frac{1}{2}, \frac{1}{2}\rangle$$

$$= \frac{3}{4} |\frac{1}{2}, \frac{1}{2}\rangle = \frac{1}{2}(1 + \frac{1}{2}) |\frac{1}{2}, \frac{1}{2}\rangle \quad \checkmark$$