

Derivation of general representation of angular momentum algebra

Let $\vec{J} = (J_x, J_y, J_z)$ be a set of Hermitian operators obeying

$$[J_x, J_y] = i\hbar J_z$$

$$[J_y, J_z] = i\hbar J_x$$

$$[J_z, J_x] = i\hbar J_y$$

Define $J^2 = J_x^2 + J_y^2 + J_z^2$

[earlier we saw that for $sp = \frac{1}{2}$, S^2 and S_z are compatible because $|\pm\rangle$ are eigenstates of both.]

HW: show $[J_z, J^2] = 0$

$\therefore J^2$ and J_z are compatible and have a common set of mutual eigenstates $|\alpha, m\rangle$

$$J^2 |\alpha, m\rangle = \alpha\hbar^2 |\alpha, m\rangle$$

$$J_z |\alpha, m\rangle = m\hbar |\alpha, m\rangle$$

[NB we do not assume α, m are integers]

Assume $|\alpha, m\rangle$ form an orthonormal basis

Define $J_{\pm} = J_x \pm iJ_y$

As before $[J_z, J_{\pm}] = \pm \hbar J_{\pm}$

so $J_z (J_{\pm} |\alpha, m\rangle) = (m \pm 1)\hbar (J_{\pm} |\alpha, m\rangle)$

i.e. $J_{\pm}$ raises/lowers the J_z eigenvalue by 1

$$[J^2, J_x] = 0$$

$$[J^2, J_y] = 0$$

$$\Rightarrow [J^2, J_{\pm}] = 0$$

$$J^2 (J_{\pm} |\alpha, m\rangle) = \alpha \hbar^2 (J_{\pm} |\alpha, m\rangle)$$

i.e. $J_{\pm}$ does not change the J^2 eigenvalue

$$J_{\pm} |\alpha, m\rangle \sim |\alpha, m \pm 1\rangle$$

To prove: $\alpha \geq m^2$

$$J^2 = J_z^2 + J_y^2 + J_x^2$$

$$\langle \alpha, m | J^2 | \alpha, m \rangle = \langle \alpha, m | J_z^2 | \alpha, m \rangle + \langle \alpha, m | J_y^2 | \alpha, m \rangle + \langle \alpha, m | J_x^2 | \alpha, m \rangle$$

$$\alpha \hbar^2 = m^2 \hbar^2 + \quad ? \quad + \quad ??$$

Recall $\langle X | X \rangle \geq 0$ for any $|X\rangle$

Let $|X\rangle = A|Y\rangle \Rightarrow \langle X| = \langle Y|A^\dagger$ then $\langle Y|A^\dagger A|Y\rangle \geq 0$

If $A = A^\dagger$ then $\langle Y|A^2|Y\rangle \geq 0$

J_x and J_y are Hermitian so ? and ?? are ≥ 0

$$\Rightarrow \alpha \geq m^2$$

Hence J_+, J_- cannot operate indefinitely

$\exists$ a maximal value of m , call it $j \Rightarrow J_+ |\alpha, j\rangle = 0$

and a minimum value, call it $\bar{j} \Rightarrow J_- |\alpha, \bar{j}\rangle = 0$

of states = $j - \bar{j} + 1$

$$\begin{array}{c}
 J_+ \\
 \uparrow \\
 \left\{ \begin{array}{l} |\alpha, j\rangle \\ |\alpha, j-1\rangle \end{array} \right\} \\
 J_- \\
 \downarrow
 \end{array}$$

$$\begin{array}{c}
 J_+ \\
 \uparrow \\
 \left\{ \begin{array}{l} |\alpha, \bar{j}\rangle \end{array} \right\} \\
 J_- \\
 \downarrow
 \end{array}$$

This finite set of states is a representation of angular momentum algebra

(m 's not necessarily integers)

We previously showed $J^2 = J_z^2 + \hbar J_z + J_- J_+$

Act on top state

$$J^2 |\alpha, j\rangle = (J_z^2 + \hbar J_z + J_- J_+) |\alpha, j\rangle$$

$$\alpha \hbar^2 |\alpha, j\rangle = (j^2 \hbar^2 + \hbar(j\hbar) + 0) |\alpha, j\rangle$$

$$\Rightarrow \alpha = j(j+1)$$

We also showed $J^2 = J_z^2 - \hbar J_z + J_+ J_-$

Act on lowest state

$$\alpha \hbar^2 |\alpha, \bar{j}\rangle = ((\bar{j}\hbar)^2 - \hbar(+\bar{j}\hbar) + 0) |\alpha, \bar{j}\rangle$$

$$\alpha = \bar{j}(\bar{j}-1)$$

set $\bar{j}(\bar{j}-1) = j(j+1)$ $\Rightarrow \bar{j} = \begin{cases} -j \\ j+1 \end{cases}$ because $\bar{j} < j$

" $\bar{j}^2 - \bar{j} - j(j+1) = 0$

$$\left. \begin{array}{l} |\alpha, j\rangle \\ \vdots \\ |\alpha, \bar{j}\rangle \end{array} \right\}$$

$$\# \text{ of states} = j - \bar{j} + 1 = 2j + 1 \in \mathbb{Z}$$

j can be either
integer (orbital angular momentum for spin-0 boson)

or
half-integer (spin-1/2 fermion)

"Spin j representation"

~~Most general state in the rep~~

Change in notation:

instead of $|\alpha, m\rangle$ w $\alpha = j(j+1)$

we write $|j, m\rangle$ where $j = m_{\max}$

$$J^2 |j, m\rangle = j(j+1) \hbar^2 |j, m\rangle$$

$$J_z |j, m\rangle = m \hbar |j, m\rangle$$

Assuming that $|j, m\rangle$ is an orthonormal basis,

$$J_+ |j, m\rangle = c_+(j, m) |j, m+1\rangle$$

$$J_- |j, m\rangle = c_-(j, m) |j, m-1\rangle$$

(Problem: calculate $c_+(j, m)$ and $c_-(j, m)$)

$$\left[\begin{array}{l} \text{Ans: } c_+ = \sqrt{j(j+1) - m(m+1)} \quad \hbar = \sqrt{(j-m)(j+m+1)} \hbar \\ c_- = \sqrt{j(j+1) - m(m-1)} \quad \hbar = \sqrt{(j+m)(j-m+1)} \hbar \end{array} \right]$$

most general state of spin j is $|\psi\rangle = \sum_{m=-j}^j a_m |j, m\rangle$
has well-defined J^2 but not J_z