

OPERATORS and EIGENVALUES

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An operator $\hat{A}$ acts on a ket to give another ket

$$\hat{A}|\psi\rangle = |\chi\rangle$$

An operator is linear if

$$\hat{A}[|\psi\rangle + |\phi\rangle] = \hat{A}|\psi\rangle + \hat{A}|\phi\rangle$$

$$\text{and } \hat{A}[\alpha|\psi\rangle] = \alpha \hat{A}|\psi\rangle$$

Eigensate of an operator (math def.)

If $\hat{A}|\psi\rangle$ produces a ket proportional to $|\psi\rangle$

$$\text{i.e. } \hat{A}|\psi\rangle = \lambda|\psi\rangle$$

then $|\psi\rangle$ is called an eigensate / eigenket / eigenstate
of $\hat{A}$ and λ is its associated eigenvalue

- If $|\psi\rangle$ is an eigensate, so is $\alpha|\psi\rangle$ for $\alpha \in \mathbb{C}$ [due to linearity]
- If ~~more than~~ ^{only} one ^{independent} eigensate corresponds to ~~the same~~ ^{a given} eigenvalue λ
then that eigensate is called "nondegenerate"
- If several indep eigensats corresp to same eigenvalue λ
degenerate: # of indep eigensats = degree of degeneracy

Observables $\leftrightarrow$ operators

Associated with each observable A

is a linear, Hermitian operator $\hat{A}$

whose eigenvalues are the possible results of measuring A ,
 and whose eigenstates are state vectors of definite A .
 $\hat{A}|a_n\rangle = a_n|a_n\rangle$ [definition of Hermitian later]

(see phys. def. = math def.)

Since $|+\rangle$ and $|-\rangle$ are states of definite S_z ,

$$\hat{S}_z |+\rangle = \frac{\hbar}{2} |+\rangle$$

$$\hat{S}_z |-\rangle = -\frac{\hbar}{2} |-\rangle$$

Note: states of indefinite S_z are not eigenstates of $\hat{S}_z$

$$\begin{aligned} \hat{S}_z |\psi\rangle &= \hat{S}_z (\alpha|+\rangle + \beta|-\rangle) \\ &= \alpha \hat{S}_z |+\rangle + \beta \hat{S}_z |-\rangle \quad (\text{by linearity}) \\ &= \alpha \frac{\hbar}{2} |+\rangle + \beta \left(-\frac{\hbar}{2}\right) |-\rangle \quad (\text{eigenvalues}) \\ &= \frac{\hbar}{2} (\alpha|+\rangle - \beta|-\rangle) \end{aligned}$$

Q: Is $|\psi\rangle$ an eigenstate of $\hat{S}_z$?

iff $\alpha=0$ or $\beta=0$

Inner product [There is a "dual vector space" isomorphic to state space] $\rightarrow$ not necessary if on dual 5-30

Associated to each ket $|\psi\rangle$ in the vector space of states is a "bra" $\langle\psi|$ in the dual vector space (left half of bracket)

[eg $z \rightarrow z^*$]

Combine $\langle\psi|$ with $|\phi\rangle$ to get bracket $\langle\psi|\phi\rangle$ which is a complex number

[like a dot product]

"inner product of $|\psi\rangle$ & $|\phi\rangle$ "

Postulate

① Inner product is linear

$$\langle\psi|[c|\phi\rangle] = c \langle\psi|\phi\rangle$$

$$\langle\psi|[\phi_1\rangle + \phi_2\rangle] = \langle\psi|\phi_1\rangle + \langle\psi|\phi_2\rangle$$

② Self-adjoint: $\langle\psi|\phi\rangle = \langle\phi|\psi\rangle^*$

[almost sym'm like dot product but not quite]

(Note $\langle\psi|\psi\rangle = \langle\psi|\psi\rangle^*$ so $\langle\psi|\psi\rangle$ is real.)

③ non-negative: $\langle\psi|\psi\rangle \geq 0$ for $|\psi\rangle \neq 0$

If $|\psi\rangle = \alpha|+\rangle + \beta|-\rangle$ then $\langle\psi| = \alpha^*\langle+| + \beta^*\langle-|$

Proof: for any $\langle\phi|$, $\langle\phi|\psi\rangle = \alpha\langle\phi|+\rangle + \beta\langle\phi|-\rangle$ (linearity i.p.)
 $\langle\phi|\psi\rangle^* = \alpha^*\langle\phi|+\rangle^* + \beta^*\langle\phi|-\rangle^*$ (c.c. both sides)
 $\langle\psi|\phi\rangle = \alpha^*\langle+|\phi\rangle + \beta^*\langle-|\phi\rangle$ (postulate ②)
 $\langle\psi| = \alpha^*\langle+| + \beta^*\langle-|$ (because valid for any $|\phi\rangle$)

Normalization

$\sqrt{\langle \psi | \psi \rangle}$ is called the norm of $|\psi\rangle$

If $\langle \psi | \psi \rangle = 1$, then $|\psi\rangle$ is normalized

If $\langle \psi | \psi \rangle = N$ then define $|\psi'\rangle = \frac{1}{\sqrt{N}} |\psi\rangle$

$$\text{thus } \langle \psi' | \psi' \rangle = \frac{1}{N} \langle \psi | \psi \rangle = 1$$

This is called "normalizing $|\psi\rangle$ ".

Orthogonality

If $\langle \phi | \psi \rangle = 0$ then $|\psi\rangle$ and $|\phi\rangle$ are orthogonal.

$\langle \phi | \psi \rangle$ is a bracket (= complex #)

$|\psi\rangle\langle\phi|$ is an operator?

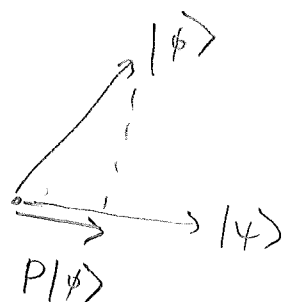
it acts on $|\chi\rangle$ to give $\langle\phi|\chi\rangle |\psi\rangle$

$$(|\psi\rangle\langle\phi|) |\chi\rangle = |\psi\rangle \langle\phi|\chi\rangle = \langle\phi|\chi\rangle |\psi\rangle$$

If $|\psi\rangle$ is normalized then

$|\psi\rangle\langle\psi|$ is a projection operator P

(which projects any ~~ket~~ $|\phi\rangle$ onto $|\psi\rangle$
and $|\psi\rangle$ onto itself)



[if $|\phi\rangle \perp |\psi\rangle$, then $P|\phi\rangle = 0$]

$$\begin{aligned} P^2 &= (|\psi\rangle\langle\psi|)(|\psi\rangle\langle\psi|) \\ &= \underbrace{\langle\psi|\psi\rangle}_1 \underbrace{|\psi\rangle\langle\psi|}_P \end{aligned}$$

$$P^2 = P$$

Assume that S_z eigenstates $|+\rangle$ and $|-\rangle$
are normalized

$$\langle + | + \rangle = 1$$

$$\langle - | - \rangle = 1$$

[if not, normalize them]

We'll prove later that they are orthogonal

$$\langle + | - \rangle = 0$$

Since we assumed that eigenstates are completely rep. the space
~~the~~ $|+\rangle$ and $|-\rangle$ form an ORTHONORMAL BASIS
for the space of states.

$$|\psi\rangle = \alpha|+\rangle + \beta|-\rangle$$

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The coefficients of a state in an orthonormal basis are given by inner product of the state w/ the corresponding eigenvector

1. p. $\alpha = \langle +|\psi\rangle$
 $\beta = \langle -|\psi\rangle$

Proof $\langle +|\psi\rangle = \langle +|(\alpha|+\rangle + \beta|-\rangle)$
 $= \alpha\langle +|+\rangle + \beta\langle +|-\rangle$
 $= \alpha$

linearity
orthonormality

$$|\psi\rangle = |+\rangle\alpha + |-\rangle\beta$$

$$= |+\rangle\langle +|\psi\rangle + |-\rangle\langle -|\psi\rangle$$

$$= |+\rangle\langle +|\psi\rangle + |-\rangle\langle -|\psi\rangle$$

$$\mathbb{I}|\psi\rangle = [|+\rangle\langle +| + |-\rangle\langle -|]|\psi\rangle$$

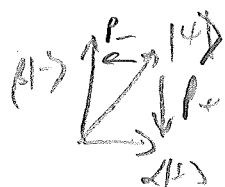
$$\mathbb{I} = |+\rangle\langle +| + |-\rangle\langle -| \quad (\text{completeness relation})$$

$$P_+ = |+\rangle\langle +| \Rightarrow P_+ + P_- = \mathbb{I}$$

$$P_- = |-\rangle\langle -|$$

In particular $P_+|\psi\rangle = |+\rangle\langle +|\psi\rangle = \alpha|+\rangle$

so P_+ is the projection operator onto $|+\rangle$ (measurement)



Also, $\langle \psi|P_+|\psi\rangle = \alpha\langle \psi|+\rangle = \alpha\langle +|\psi\rangle = \alpha\alpha^* = |\alpha|^2$
 $\langle \psi|P_-|\psi\rangle = |\beta|^2$