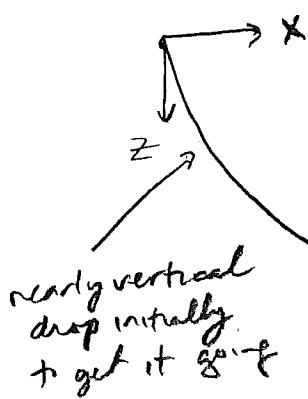


Brachistochrone (shortest time)

Johann Bernoulli posed
this in 1696 as a challenge
 - Newton solved it in a day
 - Leibniz & Huygen
also solved it



A particle starting at rest at $(0,0)$
slides down a frictionless wire
to point (x_f, z_f) . What shape
minimizes the time it takes?

$$\frac{ds}{dt} = v \quad \text{where } ds = \text{arc length along wire}$$

$$dt = \frac{ds}{v} = \frac{\sqrt{dx^2 + dz^2}}{v}$$

$$\text{Energy conserved} \Rightarrow E = \frac{1}{2}mv^2 - mgz = \text{const}$$

$$\text{initially } v=0 \text{ and } z=0 \Rightarrow E=0$$

$$\Rightarrow v = \sqrt{2gz}$$

$$T = \int dt = \int \sqrt{\frac{dx^2 + dz^2}{2gz}}$$

can ignore $\frac{1}{\sqrt{2g}}$

Rev 2024

C2

Choose either x or z as independent variable

$$T = \int dx \sqrt{\frac{1 + (\frac{dz}{dx})^2}{z}}$$

$$\text{or} \quad \int dz \sqrt{\frac{(\frac{dx}{dz})^2 + 1}{z}}$$

$$f(z, \frac{dx}{dz}, x) = \sqrt{\frac{1 + z'^2}{z}}$$

Indep of x so use 2nd form writing

↓ skip

$$f(x, \frac{dx}{dz}, z) = \sqrt{\frac{x'^2 + 1}{z}}$$

Index of x in use $\stackrel{?}{=}$ from condition

$$\frac{\partial f}{\partial x'} = \frac{x'}{\sqrt{z(x'^2 + 1)}} = k, \text{ indep of } z$$

$$x'^2 = k^2 z (x'^2 + 1)$$

$$(1 - k^2 z) x'^2 = k^2 z$$

$$x'^2 = \frac{k^2 z}{1 - k^2 z}$$

$$\frac{dx}{dz} = \sqrt{\frac{z}{(\frac{1}{k^2}) - z}}$$

$$\frac{1}{k^2} = 2R$$

$$\frac{dx}{dz} = z' = \sqrt{\frac{1}{C^2 z} - 1} = \sqrt{\frac{1 - C^2 z}{C^2 z}}$$

$$\frac{dx}{dz} = \sqrt{\frac{z}{(\frac{1}{C^2}) - z}}$$

$$\text{let } \frac{1}{C^2} = 2R$$

$$\frac{dx}{dz} = \sqrt{\frac{z}{2R - z}}$$

C3

$$\frac{dx}{dz} = \sqrt{\frac{z}{2R-z}}$$

$$x = \int dz \sqrt{\frac{z}{2R-z}}$$

[or TB: let $z = R(1 - \cos \theta)$]

$$\text{Let } z = 2R \sin^2 \alpha$$

$$dz = 4R \sin \alpha \cos \alpha d\alpha$$

$$x = 4R \int \sin \alpha \cos \alpha d\alpha \sqrt{\frac{2R \sin^2 \alpha}{2R(1-\sin^2 \alpha)}}$$

$$= 4R \int \sin^2 \alpha d\alpha$$

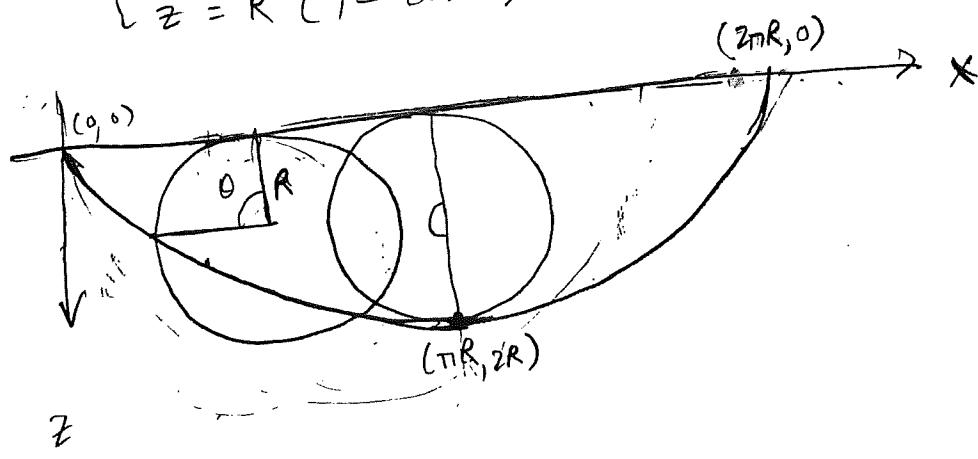
$$= 2R \int (1 - \cos 2\alpha) d\alpha$$

$$= 2R \left(\alpha - \frac{\sin 2\alpha}{2} \right)$$

$$\text{Let } \theta = 2\alpha$$

$$\begin{cases} x = R(\theta - \sin \theta) \\ z = R(1 - \cos \theta) \end{cases}$$

eqn of
cyclone



$$x_f = R (\theta_f - \sin \theta_f)$$

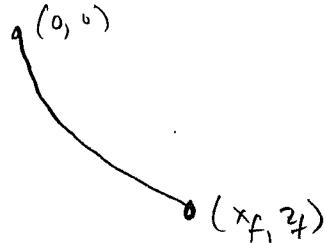
$$z_f = R (1 - \cos \theta_f)$$

$$\Rightarrow \frac{z_f}{x_f} = \frac{1 - \cos \theta_f}{\theta_f - \sin \theta_f}$$

Given the slope between $(0, 0)$ and (x_f, z_f) $0 \leq \theta_f \leq 2\pi$
 numerically solve this transcendental eqn for θ_f

The solve $R = \frac{x_f}{\theta_f - \sin \theta_f}$ w/ initial x_f .

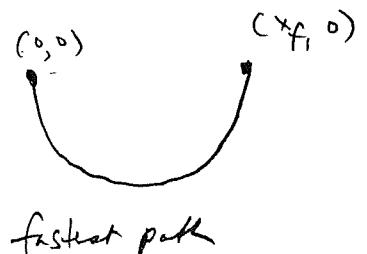
If $\frac{z_f}{x_f} > \frac{2}{\pi}$



If $\frac{z_f}{x_f} < \frac{2}{\pi}$



If $z_f = 0 \Rightarrow \theta_f = 2\pi \Rightarrow R = \frac{x_f}{2\pi}$



Let's calculate the time

$$T = \int \sqrt{\frac{dx^2 + dz^2}{2gz}}$$

$$dx = R(1 - \cos\theta) d\theta$$

$$dz = R \sin\theta d\theta$$

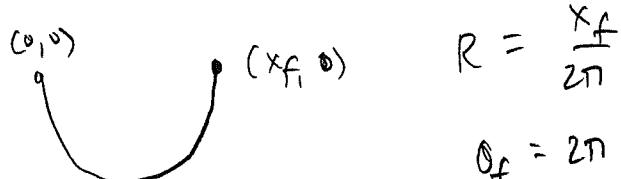
$$dx^2 + dz^2 = R^2 [(1 - 2\cos\theta + \cos^2\theta) + \sin^2\theta] d\theta^2$$

$$= R^2 (2 - 2\cos\theta) d\theta^2$$

$$2gz = 2gR(1 - \cos\theta)$$

$$T = \int \sqrt{\frac{dx^2 + dz^2}{2gz}} = \frac{R}{g} d\theta^2$$

$$T = \sqrt{\frac{R}{g}} \int d\theta = \sqrt{\frac{R}{g}} \theta_f =$$



$$R = \frac{x_f}{2\pi}$$

$$\theta_f = 2\pi$$

$$T = \sqrt{2\pi} \frac{x_f}{g} \approx 2.517 \sqrt{\frac{x_f}{g}}$$

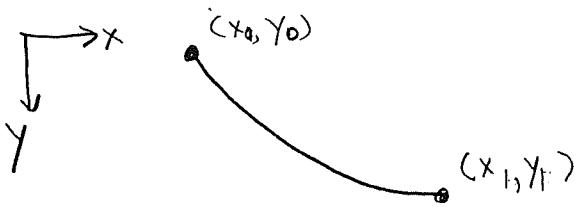
$$x_f = 1.5 \text{ m} \Rightarrow T \approx 1 \text{ sec}$$

$$\left[\text{For a circle } T = \underbrace{2.62 \sqrt{\frac{R}{g}}}_{\sqrt{2} K(\frac{1}{2})}, \text{ pendulum} \right]$$

June 2006 at Sing's Conf.

- Johann Bernoulli presented it as a challenge 1696
- Newton solved this in 1 day in 1696
- Leibniz, L'Hopital & Bernoulli also found it.
- Huygen \rightarrow (1673) refined that ~~cycloid~~ is a tautochrone.

Brachistochrone



$$\left[\frac{1}{2}mv^2 + mgy = \frac{1}{2}mv_0^2 - mgy_0 \right]$$

$$v = \sqrt{v_0^2 + 2g(y-y_0)}$$

Let $v_0=0$ and $y_0=0 \Rightarrow v=\sqrt{2gy}$

Time of trajectory $T[y(x)] = \int \frac{ds}{v} = \int \frac{\sqrt{1+y'^2} dx}{\sqrt{2gy}}$

$$\Rightarrow L = \sqrt{1+y'^2} \Rightarrow \frac{\partial L}{\partial y} = -\frac{1}{2}\sqrt{\frac{1+y'^2}{y^3}} \Rightarrow \frac{\partial L}{\partial y'} = \frac{y'}{\sqrt{y(1+y'^2)}}$$

$$\begin{aligned} \frac{d}{dx}\left(\frac{\partial L}{\partial y'}\right) &= \frac{y''}{\sqrt{y(1+y'^2)}} - \frac{1}{2}\frac{y'^2}{\sqrt{y^3(1+y'^2)}} - \frac{y'^2 y''}{\sqrt{y(1+y'^2)^3}} \\ &= \frac{y(1+y'^2)y'' - \frac{1}{2}y'^2(1+y'^2) - yy'^2y''}{\sqrt{y^3(1+y'^2)^3}} \end{aligned}$$

$$= \frac{yy'' - \frac{1}{2}y'^2(1+y'^2)}{\sqrt{y^3(1+y'^2)^3}} = -\frac{1}{2}\frac{(1+y'^2)^2}{\sqrt{y^3(1+y'^2)}}$$

$$2yy'' - y'^2(1+y'^2) = - (1+y'^2)^2$$

$$\Rightarrow \boxed{2yy'' + 1+y'^2 = 0}$$

Better to parametrize curve: $x(p), y(p)$

$$\text{Then } y' = \frac{\dot{y}}{\dot{x}} \Rightarrow y'' = \frac{\ddot{y}}{\dot{x}^2} - \frac{\dot{y}\ddot{x}}{\dot{x}^3} = \frac{\ddot{x}\dot{y} - \dot{y}\ddot{x}}{\dot{x}^3}$$

$$\Rightarrow \boxed{2y(\ddot{x}\dot{y} - \dot{y}\ddot{x}) + \dot{x}^3 + \dot{x}\dot{y}^2 = 0}$$

since $\frac{\partial L}{\partial x} = 0$
one can use
Beltrami identity

$$xL - y \frac{\partial L}{\partial y} = C$$

$$\sqrt{\frac{1+y'^2}{y}} - \frac{y'^2}{\sqrt{y(1+y'^2)}} = C$$

$$\frac{1}{\sqrt{y(1+y'^2)}} = C$$

$$y(1+y'^2) = k^2$$

$$\Rightarrow \begin{cases} x = \frac{k^2}{2}(p - \sin p) \\ y = \frac{k^2}{2}(1 - \cos p) \end{cases}$$

Brachistochrone (Cont.)

$$T = \int \frac{\sqrt{\dot{x}^2 + \dot{y}^2}}{\sqrt{2gy}} dt$$

$$L = \sqrt{\frac{\dot{x}^2 + \dot{y}^2}{y}}, \quad \frac{\partial L}{\partial y} = -\frac{1}{2} \sqrt{\frac{\dot{x}^2 + \dot{y}^2}{y^3}}, \quad \frac{\partial L}{\partial \dot{y}} = \frac{\dot{y}}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}}$$

$$\frac{\partial L}{\partial x} = 0, \quad \frac{\partial L}{\partial \dot{x}} = \frac{\dot{x}}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}}$$

$$\frac{d}{dt} \left(\frac{\dot{x}}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}} \right) = 0 \Rightarrow \boxed{\dot{x}^2 = \frac{1}{2} K y (\dot{x}^2 + \dot{y}^2)}$$

$$\hookrightarrow \frac{\ddot{x}}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}} + \dot{x} \frac{d}{dt} \left(\frac{1}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}} \right) = 0$$

$$\frac{d}{dt} \left(\frac{\dot{x}}{\dot{y}} \right) = \frac{\ddot{y}}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}} + \dot{y} \underbrace{\frac{d}{dt} \left(\frac{1}{\sqrt{y(\dot{x}^2 + \dot{y}^2)}} \right)}_{-\frac{\dot{x}}{\dot{x} \sqrt{y(\dot{x}^2 + \dot{y}^2)}}} = -\frac{1}{2} \frac{\dot{x}^2 + \dot{y}^2}{\sqrt{y^3(\dot{x}^2 + \dot{y}^2)}}$$

$$\ddot{y} - \frac{\dot{y}}{\dot{x}} \ddot{x} = -\frac{1}{2} \frac{\dot{x}^2 + \dot{y}^2}{y}$$

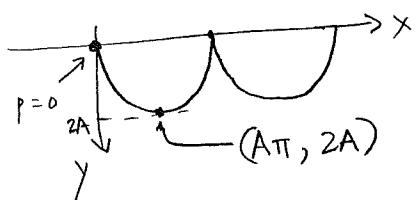
$$\boxed{2y(\dot{x}\ddot{y} - \dot{y}\ddot{x}) + \dot{x}(\dot{x}^2 + \dot{y}^2) = 0} \text{ as before}$$

Brachistochrone (cont)

To show: cycloid = brachistochrone

$$\begin{aligned} \dot{x}^2 + \frac{1}{2}K\gamma(\dot{x}^2 + \dot{y}^2) \\ 2\gamma(\ddot{x}\dot{y} - \dot{y}\ddot{x}) + \dot{x}(\dot{x}^2 + \dot{y}^2) = 0 \end{aligned} \quad \left. \begin{array}{l} \dot{x}^2 + \frac{1}{2}K\gamma(\dot{x}^2 + \dot{y}^2) \\ 2\gamma(\ddot{x}\dot{y} - \dot{y}\ddot{x}) + \dot{x}(\dot{x}^2 + \dot{y}^2) = 0 \end{array} \right\} \quad \ddot{y}^2 - K\gamma^2(\ddot{x}\dot{y} - \dot{y}\ddot{x}) + \dot{x}^3 = 0$$

$$\text{let } \begin{cases} x = A(p - \sin p) \\ y = A(1 - \cos p) \end{cases} \Rightarrow \begin{cases} \dot{x} = y \\ \dot{y} = A \sin p \end{cases} \Rightarrow \dot{x}^2 + \dot{y}^2 = 2A^2(1 - \cos p) = 2Ay$$



$$\text{check } \dot{x}^2 = \frac{1}{2}K\gamma(\dot{x}^2 + \dot{y}^2) \\ y^2 = \frac{1}{2}K\gamma(2Ay) \Rightarrow K = \frac{1}{A}$$

$$\text{check } K\gamma^2(\ddot{x}\dot{y} - \dot{y}\ddot{x}) + \dot{x}^3 = 0$$

$$\frac{1}{A}y^2(\underbrace{A(1-\cos p)[A\sin p] - [A\sin p]^2}_{A^2\cos p - A^2}) + \dot{x}^3 = 0 \quad \checkmark$$

$$-Ay$$

$$\left. \begin{array}{l} \text{Hence } \frac{y_1}{x_1} = \left(\frac{1 - \cos p_1}{p_1 - \sin p_1} \right) \text{ determines } p_1 \\ \text{then } y_1 = A(1 - \cos p_1) \text{ determines } A \end{array} \right\} \rightarrow \text{so if } \frac{y_1}{x_1} \geq \frac{2}{A} \text{ then } p_1 \leq \pi$$

ie vs

$$\text{Then } T(x_1, y_1) = \int_{0}^{p_1} \sqrt{\frac{\dot{x}^2 + \dot{y}^2}{2g}} dp = \int_{0}^{p_1} \sqrt{\frac{A}{g}} dp = \sqrt{\frac{A}{g}} p_1$$

along brachistochrone

$$(eg \text{ if } x_1 \rightarrow 0, \text{ then } p_1 \rightarrow 0 \text{ so } y_1 \rightarrow A \frac{1}{2}p_1^2 \text{ so } T(0, y_1) = \sqrt{\frac{2y_1}{g}} \quad \checkmark)$$

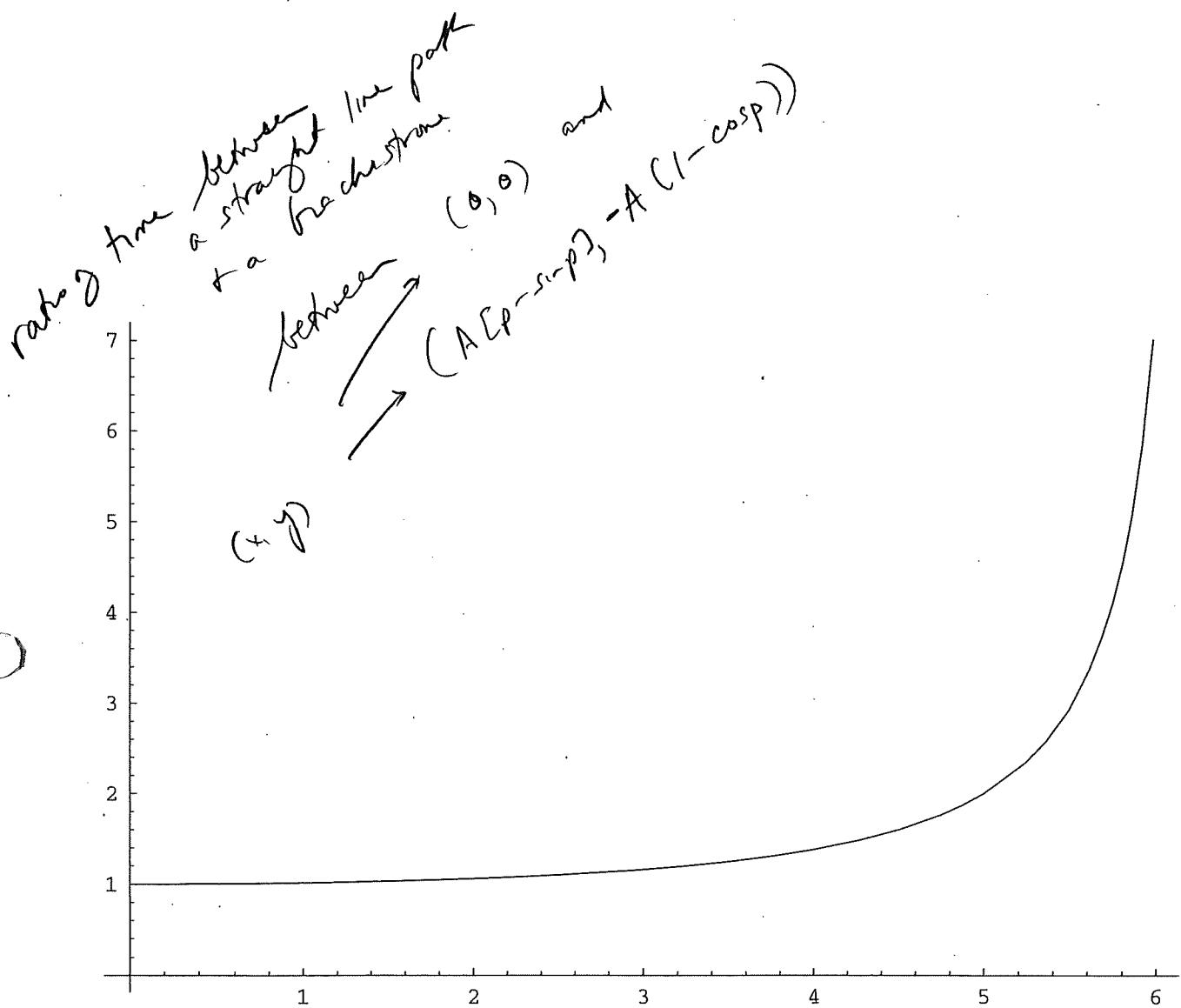
$$\text{Consider straight line path } \Rightarrow \begin{cases} x = x_1 p \\ y = y_1 p \end{cases} \Rightarrow T(x_1, y_1) = \sqrt{\frac{x_1^2 + y_1^2}{2g y_1}} \int_0^1 \sqrt{p} dp = \sqrt{\frac{2(x_1^2 + y_1^2)}{g y_1}}$$



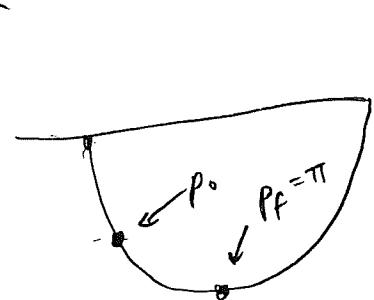
$$\text{Also: } T = \sqrt{\frac{2S}{a}} = \sqrt{\frac{2S}{g \sin \theta}} = \sqrt{\frac{2 \cdot x_1^2 + y_1^2}{g(y_1 / \sqrt{x_1^2 + y_1^2})}} = \sqrt{\frac{2(x_1^2 + y_1^2)}{g y_1}} \quad \checkmark$$

plot this.

$$\text{Compare times: } \frac{T_{\text{straight}}}{T_{\text{cyc}}} = \frac{\sqrt{\frac{2(x_1^2 + y_1^2)}{g y_1}}}{\sqrt{\frac{A}{g} p_1^2}} = \sqrt{\frac{2[(p - \sin p)^2 + (1 - \cos p)^2]}{p_1^2[1 - \cos p]}} = \boxed{\frac{2(p_1^2 - 2p_1 \sin p_1 + 2(1 - \cos p_1))}{p_1^2(1 - \cos p_1)}} \quad \text{if } p_1 \rightarrow 0 \Rightarrow \text{ratio} \rightarrow 1$$



Huygen noticed (1673) that a cycloid is ~~a clock~~ a isochrone



$$x = A(p - \sin p)$$

$$y = A(1 - \cos p)$$

Start at p_0 , end at $p_f = \pi$

$$T = \int_{p_0}^{p_f} \sqrt{\frac{x^2 + y^2}{2g(y - y_0)}} = \sqrt{\frac{A}{g}} \int_{p_0}^{p_f} \sqrt{\frac{1 - \cos p}{\cos p_0 - \cos p}} dp$$

As calc'd before, $p_0 = 0 \Rightarrow T = \sqrt{\frac{A}{g}} \pi$

But now (help from Mathworld)

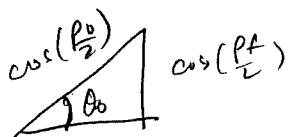
$$\begin{aligned} \cos p &= 2\cos^2\left(\frac{p}{2}\right) - 1 \\ 1 - \cos p &= 2\sin^2\left(\frac{p}{2}\right) \end{aligned} \quad \Rightarrow \int_{p_0}^{p_f} \sqrt{\frac{1 - \cos p}{\cos p_0 - \cos p}} dp = \int_{p_0}^{p_f} \frac{\sin\left(\frac{p}{2}\right)}{\sqrt{\cos^2\left(\frac{p}{2}\right) - \sin^2\left(\frac{p}{2}\right)}} dp$$

$$= \int_{p_0}^{p_f} \frac{\sin\left(\frac{p}{2}\right)}{\sqrt{1 - \frac{\cos^2\left(\frac{p}{2}\right)}{\cos^2\left(\frac{p_0}{2}\right)}}} dp = \int_1^{\frac{\cos\left(\frac{p_f}{2}\right)}{\cos\left(\frac{p_0}{2}\right)}} \frac{-2du}{\sqrt{1-u^2}}$$

$$u = \frac{\cos\left(\frac{p}{2}\right)}{\cos\left(\frac{p_0}{2}\right)}$$

$$du = -\frac{1}{2} \frac{\sin\left(\frac{p}{2}\right)}{\cos^2\left(\frac{p}{2}\right)} dp$$

$$= 2 \arcsin u \Big|_1^{\frac{\cos\left(\frac{p_f}{2}\right)}{\cos\left(\frac{p_0}{2}\right)}} = 2 \arcsin \frac{\cos\left(\frac{p_f}{2}\right)}{\cos\left(\frac{p_0}{2}\right)}$$



$$= \pi - \theta_0 \xrightarrow{\substack{p_f \rightarrow \pi \\ \cos(p_0/2) \rightarrow 0 \\ \theta_0 \rightarrow 0}} \pi, \text{ indy } p_0$$

Thus time from p_0 to π is indy of starting pt.

~~is perpendic to circle at H~~

\therefore an isochrone-indy pendulum

(M/18)



Cycloid



$$x = R(\theta - \sin \theta)$$

$$z = -R(1 - \cos \theta)$$

$$\text{Let } \theta = \pi + \phi$$

$$x = R(\pi + \phi - \sin(\pi + \phi))$$

$$= R\pi + R(\phi + \sin \phi)$$

$$z = -R(1 - \cos(\pi + \phi))$$

$$= -R - R \cos \phi$$

$$= -2R + R(1 - \cos \phi)$$

Shift origin by $R\pi, -2R$

parametric path

$$\begin{cases} x = R(\phi + \sin \phi) \\ z = R(1 - \cos \phi) \end{cases} \approx \begin{cases} 2R\phi \\ \frac{1}{2}R\phi^2 \end{cases}$$

$$\Rightarrow z \approx \frac{1}{8R}x^2 \Rightarrow U = mgz = \frac{mg}{8R}x^2$$

$$\dot{x} = R(1 + \cos \phi) \dot{\phi}$$

$$\dot{z} = R \sin \phi \dot{\phi}$$

$$L = \frac{1}{2}mR^2 \left[\underbrace{(1 + \cos \phi)^2 + (\sin \phi)^2}_{2 + 2 \cos \phi} \right] \dot{\phi}^2 - mg R(1 - \cos \phi)$$

~~$$\frac{\partial L}{\partial \dot{\phi}} = 2mR^2(1 + \cos \phi) \dot{\phi}$$~~
~~$$\frac{\partial L}{\partial \phi} = 2L - 2mR^2 \sin \phi \dot{\phi}^2 - mg R \sin \phi$$~~

$$\dot{\phi} \frac{\partial L}{\partial \dot{\phi}} - L = mR^2(1 + \cos \phi) \dot{\phi}^2 + mg R(1 - \cos \phi) = E$$

$$dt = \sqrt{\frac{(1 + \cos \phi)}{E - \frac{g}{R}(1 - \cos \phi)}} d\phi = \int_0^{2\pi} \frac{d\phi}{\sqrt{\frac{g}{R}(1 + \cos \phi)(\cos \phi - \cos \phi_0)}}$$