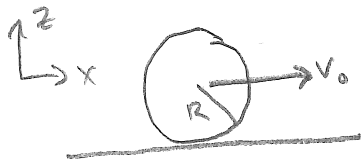


Sliding bowling ball



Hits ground w/ speed v_0 and no rotation.

When does it begin to roll?

What is its speed?

How far has it slid?



$$F_f = \mu N = \mu mg$$

μ = coeff of kinetic friction

$$\text{Net force } \vec{F} = -\mu mg \hat{i}$$

$$\vec{F} = m\vec{a}_{cm} \quad a_x = \frac{F_x}{m} = -\mu g$$

$$v_x = v_0 - \mu g t$$

$$x = x_0 + v_0 t - \frac{1}{2} \mu g t^2$$

$$\text{Torque about cm } \vec{\tau} = R F_f \hat{j} = \mu mg R \hat{j}$$

$$\vec{\tau} = I \vec{\alpha} \quad \alpha_y = \frac{\tau_y}{I} = \frac{R \mu mg}{I} = \frac{5}{2} \frac{\mu g}{R}$$

$$\omega_y = \frac{5}{2} \frac{\mu g}{R} t$$

Ball starts to roll when $v_x = R \omega_y$ (point of contact not moving)

$$v_0 - \mu g t = \frac{5}{2} \mu g t$$

$$t_{\text{roll}} = \frac{2}{7} \frac{v_0}{\mu g}$$

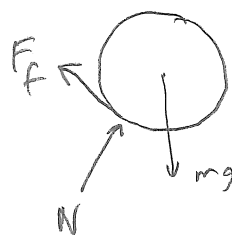
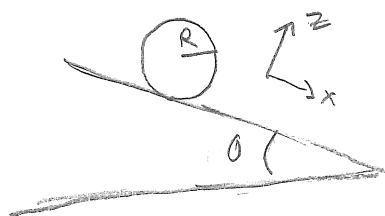
$$\text{Speed of ball } v_x = R \left(\frac{5}{2} \frac{\mu g}{R} \right) \left(\frac{2}{7} \frac{v_0}{\mu g} \right) = \frac{5}{7} v_0$$

$$\text{It has slid } v_0 t - \frac{1}{2} \mu g t^2 = \frac{2}{7} \frac{v_0^2}{\mu g} - \frac{2}{49} \frac{v_0^2}{\mu g} = \frac{12}{49} \frac{v_0^2}{\mu g}$$

What happens after roll?

$$[F_f = 0, \quad \vec{v} + \vec{\omega} \text{ are constant}]$$

Cylinder rolling down incline



Net force

$$F_x = mg \sin \theta - F_f = m a_x$$

[What is F_f ? $\mu mg \cos \theta$? No!] F_f is what it needs to be to keep ball rolling provided $F_f \leq \mu mg \cos \theta$.

Torque about cm [accelerating]

$$\tau_y = R F_f = I \alpha_y$$

Assume rolling w/o sliding:

$$v_x = R \omega_y$$

$$a_x = R \alpha_y$$

$$\Rightarrow mg \sin \theta - F_f = m R \alpha_y = \frac{m R^2 F_f}{I}$$

$$F_f = \frac{mg \sin \theta}{(1 + \frac{m R^2}{I})}$$

$$a_x = g \sin \theta - \frac{F_f}{m} = \frac{g \sin \theta}{(\frac{I}{m R^2} + 1)}$$

eg. $I = \frac{1}{2} m R^2$
 $\Rightarrow a = \frac{2}{3} g \sin \theta$

check whether cylinder will roll or slide

$$F_f \leq \mu mg \cos \theta$$

~~$$\sin \theta \leq (1 + \frac{m R^2}{I}) \mu \cos \theta$$~~

$$\tan \theta \leq \mu (1 + \frac{m R^2}{I})$$

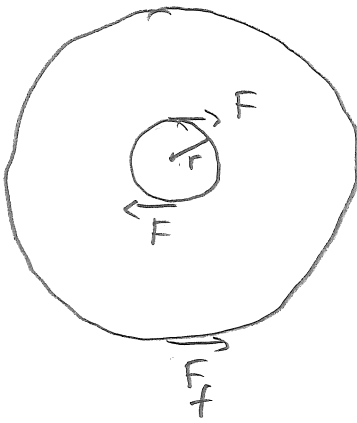
eg. disk $I = \frac{1}{2} m R^2$
 $\tan \theta \leq 3\mu$

otherwise will start sliding and also begin to roll.

(Jan 2018)

Work done to accelerate a wheel

03



$$\bullet F_f = m a_{cm} \quad \text{net force}$$

$$\bullet 2rF - RF_f = I\alpha \quad \text{net torque}$$

$$\bullet a = R\alpha \quad \text{rolling w/o slipping}$$

$\Downarrow$

$$\frac{F_f}{m} = R \frac{(2rF - RF_f)}{I}$$

$$2mRrF = (I + mR^2)F_f$$

Frictional force does no work because point of contact moves perpendicular to force. All work is done by the torque applied to the axle

$$\Delta W = \tau \Delta \phi$$

$$W = 2F r \Delta \phi = 2rF \omega \Delta t = \left(\frac{I + mR^2}{mR} \right) F_f \omega \Delta t$$

Change in kinetic energy

$$K = \frac{1}{2}mv^2 + \frac{1}{2}I\omega^2 = \frac{1}{2}(I + mR^2)\omega^2$$

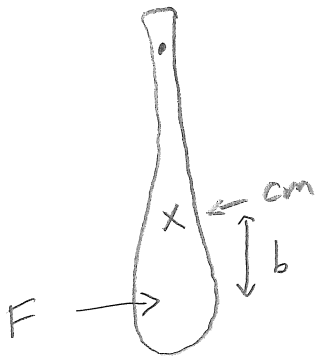
$\uparrow$
 $v = \omega R$

$$\Delta K = (I + mR^2)\omega \Delta \omega$$

$$= (I + mR^2)\omega \alpha \Delta t \quad \text{but } \alpha = \frac{a}{R} = \frac{F_f}{mR}$$

$$= \left(\frac{I + mR^2}{mR} \right) \omega F_f \Delta t, \text{ agrees with above}$$

Center of percussion



A force F is applied for a short time interval Δt at a distance b below the cm. of some object.

$b = \text{impact parameter.}$

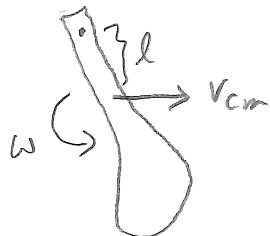
$$F = m \frac{\Delta v_{cm}}{\Delta t}$$

If object initially at rest, after the impulse $v_{cm} = \frac{F \Delta t}{m}$

The force results in a torque about cm $\tau = bF$

$$\tau = I' \frac{\Delta \omega}{\Delta t}$$

After the impulse, object rotates $\omega = \frac{\tau \Delta t}{I'} = \frac{b F \Delta t}{I'}$



Speed of a point a distance l above cm is $v_{cm} - l\omega$

This point remains stationary if

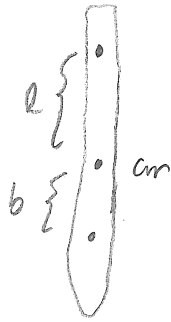
$$l\omega = v_{cm}$$

it acts as a pivot

$$l \frac{b F \Delta t}{I'} = \frac{F \Delta t}{m}$$

$$b = \frac{I'}{ml}$$

Point of impact called
Center of percussion



Distance of impact point from pivot

$$b + l = \frac{(I' + ml^2)}{ml} = \frac{I}{ml}$$

where I = moment of inertia about pivot

This is same as radius of gyration, k_{eff} .

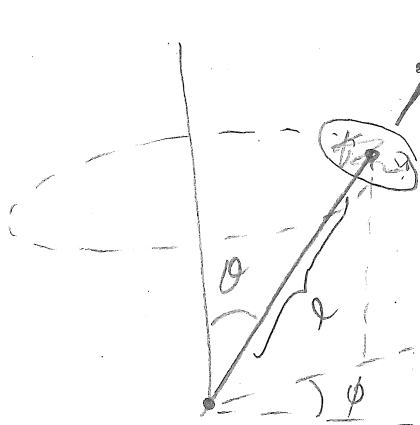
For a rod of length $L = 2l$, $I = \frac{1}{3}ML^2$

$$k_{\text{eff}} = \frac{2}{3}L$$

(For a bat, a bit longer because heavier at the end)
beyond the cm

Simplified treatment of precessing top

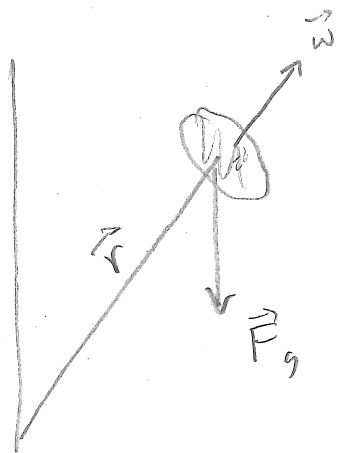
Consider a top with its tip fixed
spinning w/ angular speed $\vec{\omega}$



The external gravitational torque
will cause it to precess
around a vertical line
w/ precession frequency $\dot{\phi}$

given approximately by $\frac{mgl}{I\omega}$

We'll derive this now in a simplified treatment
and later more rigorously



$$\vec{L} = I\vec{\omega}$$

$$\vec{\tau}_{\text{ext}} = \vec{r} \times \vec{F}_g$$

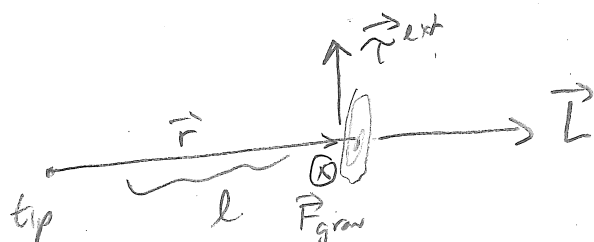
($\vec{L}$ & $\vec{r}$ both defined
wrt. fixed tip,

force on tip exerts no torque.)

$$\frac{d\vec{L}}{dt} = \vec{\tau}_{\text{ext}}$$

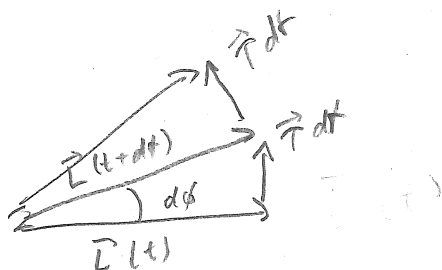
First consider $\theta = \frac{\pi}{2}$

View from overhead



$$\frac{d\vec{L}}{dt} = \vec{\tau}_{\text{ext}} \Rightarrow \vec{L}(t+dt) = \vec{L}(t) + \vec{\tau} dt$$

torque causes $\vec{L}$ to change direction ($\tau \therefore$ top itself)

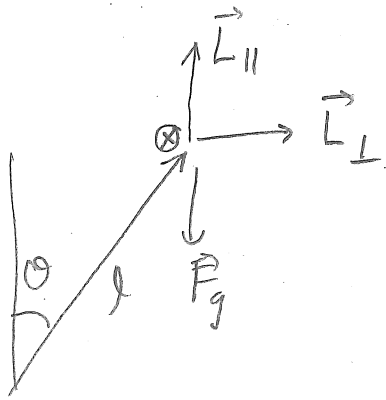


direction of top changes by $d\phi = \frac{\tau dt}{L}$

$$= \frac{Lmg}{I\omega} dt$$

$$\dot{\phi} = \frac{mgl}{I\omega}$$

Now consider arbitrary θ

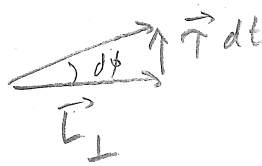


$$|\vec{\tau}| = lmg \sin \theta$$

$$L_{\perp} = I\omega \sin \theta$$

$$L_{\parallel} = I\omega \cos \theta$$

Torque causes $\vec{L}_{\perp}$ to change direction; $\vec{L}_{\parallel} = \text{const}$

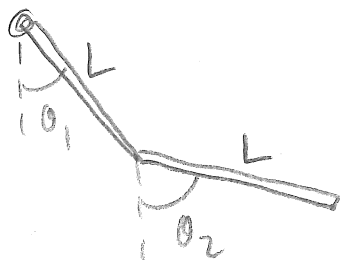


$$d\phi = \frac{mgl \sin \theta}{I\omega \sin \theta} dt$$

$$\dot{\phi} = \frac{mgl}{I\omega}$$

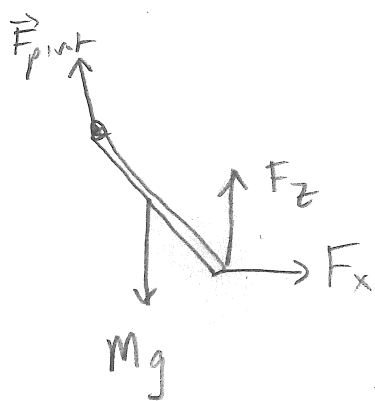
precession rate independent of
angle of tilt.

Double pendulum



Treat each limb as a system.

Assume each limb has length L
and mass M



Let F_x, F_z be components of force
exerted on upper by lower limb

Calculate torque about pivot so we can
ignore $\vec{F}_{pivot}$

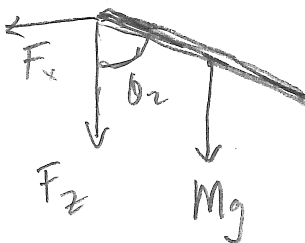
$$\begin{aligned}\tau_y &= \frac{L}{2} Mg \sin\theta - L F_x \cos\theta_1 - L F_z \sin\theta_1 \\ &= \frac{dL_y}{dt} = -I\ddot{\theta}_1 \quad \text{where } I = \frac{1}{3} ML^2\end{aligned}$$

$$\Rightarrow \ddot{\theta}_1 = -\frac{3}{2} \frac{g}{L} \sin\theta_1 + \frac{3F_x}{mL} \cos\theta_1 + \frac{3F_z}{mL} \sin\theta_1$$

Define $\omega_0^2 \equiv \frac{g}{L}$

$$\ddot{\theta}_1 = -\frac{3}{2} \omega_0^2 \sin\theta_1 + 3 \frac{F_x}{mL} \cos\theta_1 + 3 \frac{F_z}{mL} \sin\theta_1$$

Force in lower limb are equal & opposite



Calc torque about cm (accelerating)
 → can't calc about joint because it is not inertial.

$$\tau_y = -\frac{L}{2} F_x \cos \theta_2 - \frac{L}{2} F_z \sin \theta_2$$

$$= \frac{dL_y}{dt} = -I' \ddot{\theta}_2 \quad \text{where } I' = \frac{1}{12} ML^2$$

$$\Rightarrow \ddot{\theta}_2 = \frac{6F_x}{ML} \cos \theta_2 + \frac{6F_z}{ML} \sin \theta_2$$

2 eqns, 4 unknowns $\theta_1, \theta_2, F_x, F_z$

Use $\vec{F}_{\text{ext}} = m\vec{a}_{\text{cm}}$

let x, z be cm of lower link

$$-F_x = M\ddot{x}$$

$$-F_z - Mg = M\ddot{z}$$

$$x = L \sin \theta_1 + \frac{L}{2} (1 - \cos \theta_2)$$

$$z = -L \cos \theta_1 - \frac{L}{2} \cos \theta_2$$

Rest of calculation for problem set

- Take derivatives
- eliminate F_x, F_z from the 4 eqns
- write eqns for $\ddot{\theta}_1, \ddot{\theta}_2$