

Rigid body

a system of particles whose relative positions are fixed
mediates an overall translation + rotation.

In particular, all $|\vec{r}_{ij}|$ are const.

[eg. point masses connected by weightless rods (chemistry models
of molecules)]

[atoms in a solid at their equilibrium positions]

[talk informally about d.o.f. $3N$ for N -pt mass
but only 6 for a rigid body $\Rightarrow$ Euler angles]

$$dE_{\text{mech}} = dW_{\text{ext}}$$

if external forces are conservative $dW_{\text{ext}} = -dU_{\text{ext}}$

$$\Rightarrow d(\underbrace{E_{\text{mech}} + U_{\text{ext}}}_E) = 0$$

Mechanical energy of a rigid body

$$E_{\text{mech}} = \sum T_i + \sum_{i < j} U_{ij}(\vec{r}_{ij}) + U_{\text{ext}}$$

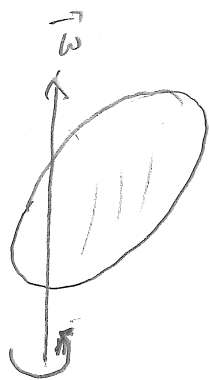
Internal constraint forces do no work $\therefore$
do not contribute to energy.

Particles at equilibrium positions $\Rightarrow \sum U_{ij} = \text{const}$
and can be ignored.

$$\Rightarrow \text{for a rigid body } E = T_{\text{sys}} + U_{\text{ext}}$$

Rotating rigid body

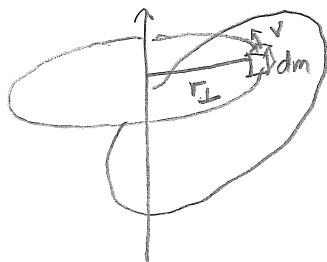
Assume rigid body is rotating about some fixed axis
of angular velocity $\vec{\omega}$



$\hat{\omega}$ = axis of rotation
(direction given by r.h. rule)

$\omega = |\vec{\omega}|$ = angular speed
(need not be constant)

Break into infinitesimal elements



element dm has speed $|\vec{v}| = \omega r_{\perp}$
where $r_{\perp}$ = perpendicular distance from axis

$$T_{\text{sys}} = \int \frac{1}{2} (dm) v^2 = \frac{1}{2} \left(\int dm r_{\perp}^2 \right) \omega^2$$

Define moment of inertia $I = \int dm r_{\perp}^2 = \int d^3r \overset{\text{density}}{\rho} r_{\perp}^2$

$$T_{\text{sys}} = \frac{1}{2} I \omega^2 \quad \text{NB.}$$

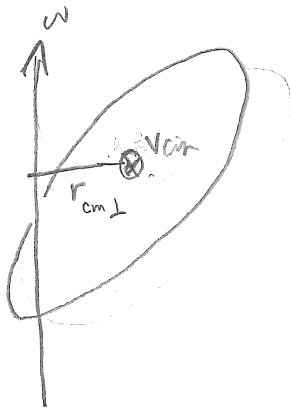
$$T^{sys} = \frac{1}{2} I \omega^2$$

L4

Earlier we showed

$$T^{sys} = \frac{1}{2} M V_{cm}^2 + T'^{sys}$$

↑
kinetic energy in cm frame



$$|\vec{V}_{cm}| = r_{cm\perp} \omega$$

$$T^{sys} = \frac{1}{2} M r_{cm\perp}^2 \omega^2 + \frac{1}{2} I' \omega^2$$

↑
mom of inertia
wrt axis through cm

$$= \frac{1}{2} (M r_{cm\perp}^2 + I') \omega^2$$

$$\Rightarrow I = M r_{cm\perp}^2 + I' \quad (\text{parallel axis theorem})$$

[cm frame is moving + accelerating but not rotating]

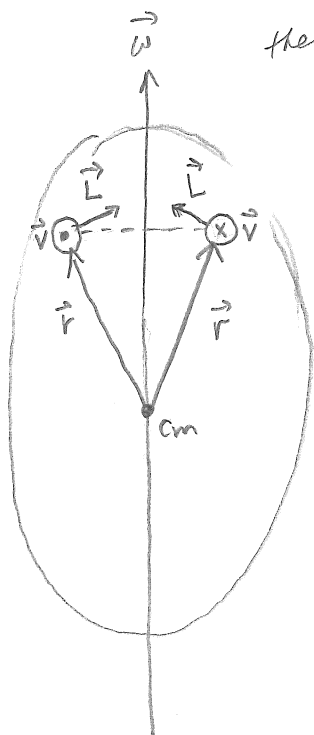
$\therefore$ object is rotating at same speed in cm as in rest frame]

Earlier we showed $\vec{L} = \vec{r}_{cm} \times M \vec{v}_{cm} + \vec{L}'$
 $\uparrow$
 any mass w.r.t. - cm

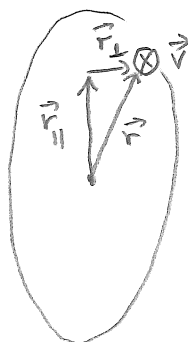
m1

Balanced objects

Consider $\vec{L}'$ for a rotating balanced body
 • for every element on one side of the axis
 there is an equal element on the opposite side



Radial components of $\vec{L}$ cancel
 leaving only components along $\vec{\omega}$



$$\vec{r} \times \vec{v} = (\vec{r}_{\parallel} + \vec{r}_{\perp}) \times \vec{v}$$

$\vec{r}_{\parallel} \times \vec{v}$ are radial

$\vec{r}_{\perp} \times \vec{v}$ are parallel to $\vec{\omega}$

$$\Rightarrow \vec{r}_{\perp} \times \vec{v} = r_{\perp} v \hat{\omega}$$

$$= r_{\perp} (r_{\perp} \omega) \hat{\omega}$$

$$= r_{\perp}^2 \vec{\omega}$$

$$\vec{L}' = \int dm \vec{r} \times \vec{v} = \left[\int dm r_{\perp}^2 \right] \vec{\omega} = I' \vec{\omega}$$

[Slater + Frank, p. 99 Mechanics = on dynamical vs static balance]

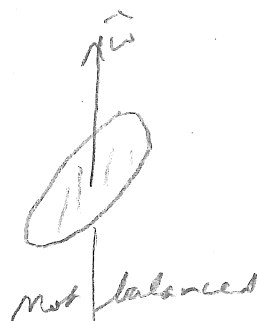
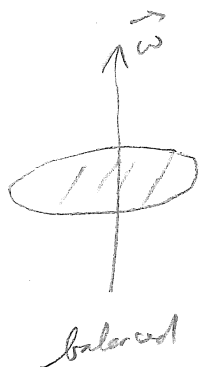
m2

more generally, we call an object dynamically balanced w.r.t. axis $\hat{\omega}$ if

$\vec{L}'$ is parallel to $\vec{\omega}$.



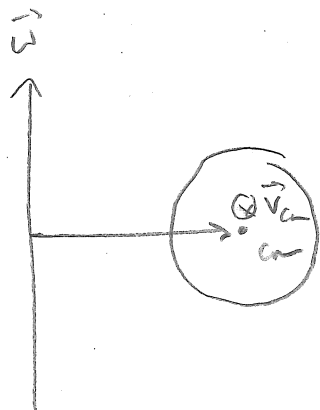
Being balanced depends not only on object but also on the axis,



We will later show that for any object (however misshapen) there are at least 3 (perpendicular) axes about which it is balanced (called principle axes)

For a balanced object

$$\vec{L} = \vec{r}_{cm} \times M \vec{v}_{cm} + I' \vec{\omega}$$



Choose origin (for $\vec{L}$) to be the point on the axis closest to the cm

$$M \vec{r}_{cm} \times \vec{v}_{cm} = M r_{cm} v_{cm} \hat{\omega} = M r_{cm} (r_{cm\perp} \omega) \hat{\omega}$$

$\downarrow$
 $= r_{cm\perp}$

$$= M r_{cm\perp}^2 \vec{\omega}$$

$$\Rightarrow \vec{L} = (M r_{cm\perp}^2 + I') \vec{\omega}$$

$$= I \quad (\text{parallel axis theorem})$$

$\vec{L} = I \vec{\omega}$ valid for any fixed axis. $\forall$ possibly changing angular speed
(provided origin is at point on axis closest to cm)

For an object balanced w/ axis $\hat{\omega}$

$$\vec{L} = I \vec{\omega}$$

$$\frac{d\vec{L}}{dt} = I \frac{d\vec{\omega}}{dt} = I \vec{\alpha}$$

$\vec{\alpha}$ = angular acceleration

Recall $\frac{d\vec{L}}{dt} = \vec{\tau}$

valid in any IRF
and also in CM frame even if accelerating

$$\vec{\tau} = I \vec{\alpha}$$

analog of

$$\vec{F} = m \vec{a}$$

Perpendicular axis theorem

- applies only to flat objects
- assume object is in the xy plane
- assume object is balanced wrt. x, y and z axes.

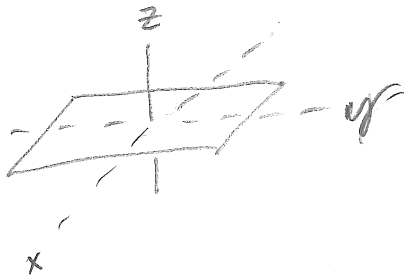
Then $I_{zz} = I_{xx} + I_{yy}$

where I_{zz} = moment of inertia wrt. z -axis, etc.

[explain double index later]

[not necessarily wrt. cm]

Proof:



$$I_{zz} = \int dm r_{\perp}^2$$

$$= \int dm (x^2 + y^2)$$

As usual
 $dm = \rho d^2r$

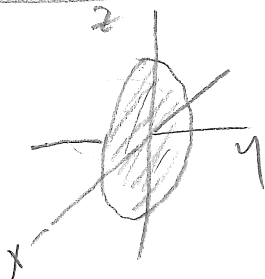
$$I_{xx} = \int dm y^2$$

$$I_{yy} = \int dm x^2$$

Q.E.D.

Note: a solid object can often be built from a stack of flat objects

Eg.



Moment of inertia of disk wrt axis
in the plane of the disk I_{zz}

$$I_{xx} = I_{yy}$$

$$I_{zz} + I_{yy} = I_{xx} = \frac{1}{2}MR^2$$

$$\Rightarrow I_{zz} = \frac{1}{4}MR^2$$