

Radioactivity

emission of particles by unstable nuclei

[1896 Becquerel discovers "uranic rays" from uranium]

[1898 Marie Curie identifies new elements (Po , Ra)
by their distinctive half-lives]

[1899 Rutherford classifies radioactivity as α, β, γ
by the curvature of their tracks in a magnetic field]

(cloud chamber: superheated vapor condenses on ions)

[see illustration]

γ = neutral = high energy photons ($E > 1 \text{ keV}$)

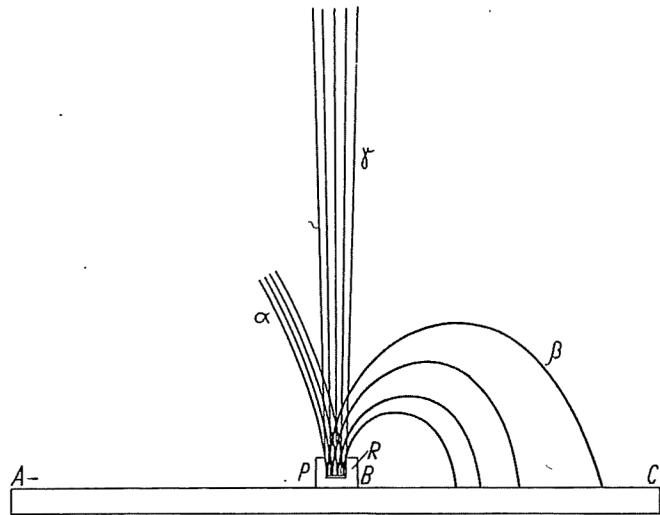
β = negatively charged = electron

α = positively charged = ${}^4\text{He}$ nuclei

[1909 Boyds, Rutherford stored this in Manchester]

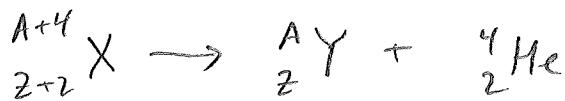
[α particles from radioactive nuclei can't accumulate in
natural gas deposits, from which it can be extracted.
Once in atmosphere, it escapes.]

Figure 1-2 Deflection of alpha, beta, and gamma rays in a magnetic field. The nomenclature is due to Rutherford (1899). [Mme Curie, Thesis, 1904.]



Segré, p. 3

α -2



$$Q = m_{nuc}(X) - m_{nuc}(Y) - m_{nuc}({}_2^4He)$$

$$= [(A+4)m_n + \Delta(X) + (Z+2)m_e] - [Am + \Delta(Y) + Zm_e] - [4m + \Delta({}_2^4He) - 2m_e]$$

$$= A(X) - \Delta(Y) - \Delta({}_2^4He) \quad \rightarrow \text{see next}$$

[In other words, since $\frac{B}{A} \uparrow$ as $A \downarrow$, and ${}_2^4He$ has large binding energy]

If $Q > 0$, α -decay can occur

Typically $Q \approx 5 \text{ MeV}$ for heavy unstable nuclei (e.g. U)

$$Q = T_Y + T_{He} \quad [\text{fin } X \text{ at rest}]$$

$$\text{Because } m_{He} \ll m_Y \Rightarrow T_{He} \gg T_Y$$

so α particle gets most of the released energy.

In air, an α -particle travels a few cm.

\Rightarrow [picture]

$$\Delta(^{12}\text{C}) = 0 \quad \text{typical}$$

$$\Delta(^{56}_{26}\text{Fe}) = -60.6 \text{ MeV} \quad \text{tightly bound}$$

$$\Delta(^{235}_{92}\text{U}) = 40.9 \text{ MeV} \quad \text{less bound}$$

$$\Delta(^{230}_{91}\text{Pa}) = 40.3 \text{ MeV} \rightarrow \text{differen: } 0.6 \text{ MeV, not because } \Delta(\rho) = 7.3$$

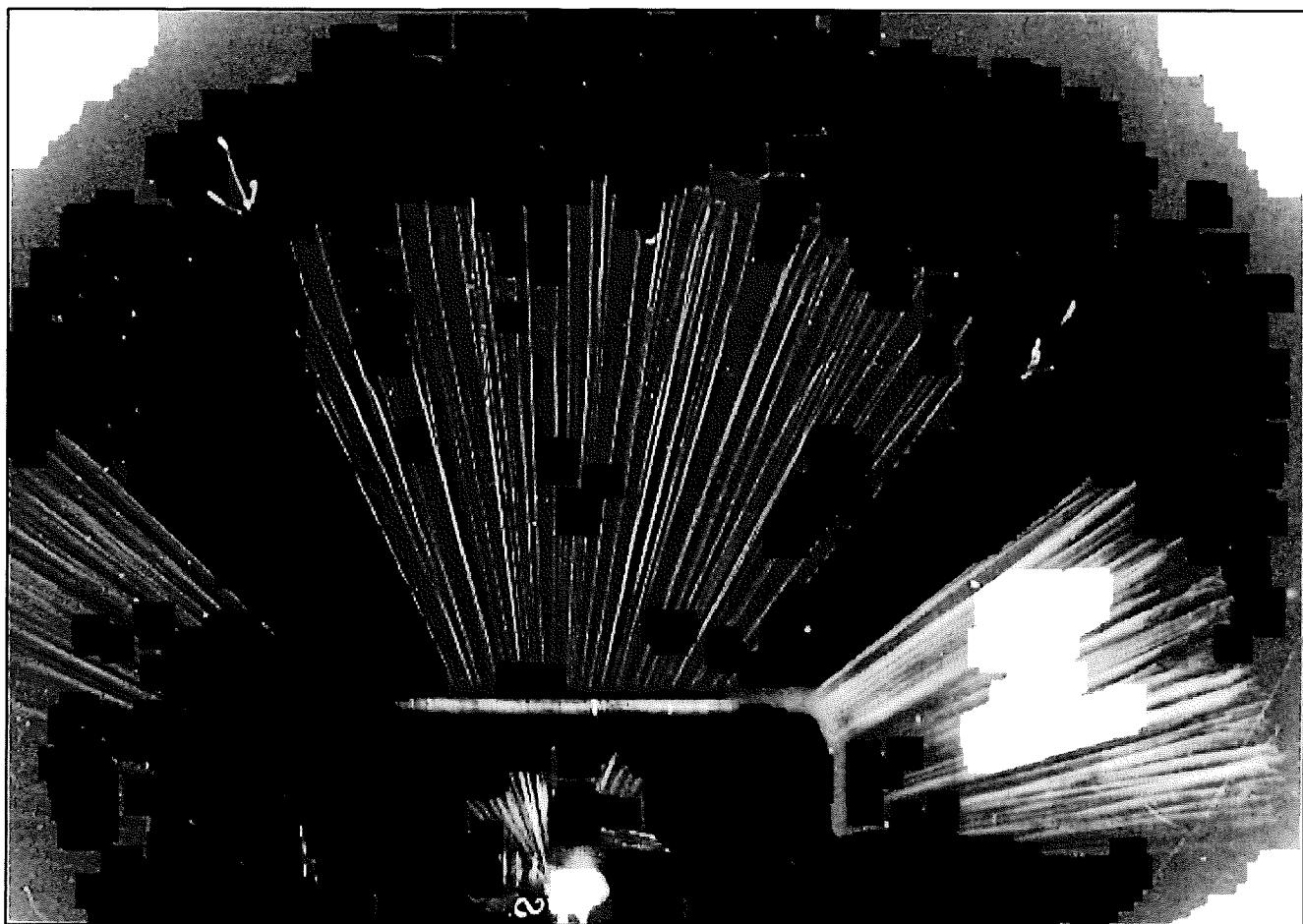
$$\Delta(^{231}_{90}\text{Th}) = 33.8 \text{ MeV} \rightarrow \text{differen: } 7 \text{ MeV, ok, because } \Delta(\alpha) = 1.425$$



$$\Delta(\rho) = 7.3$$

$$\Delta(n) = 8.1$$

$$\Delta(^4\text{He}) = 2.425$$



Because all α particles get same kinetic energy

they all travel the same distance in air, a few cm

Assume neutrino		but different nuclei have diff ranges (see picture →)
Bragg 1904		
^{232}Th	2.8 cm	
^{226}Ra	3.3 cm	
^{232}Po	8.6 cm	

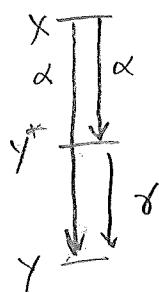
Not strictly true e.g. see ^{227}Th decay

]} maximum $K_\alpha = 6.04 \text{ MeV}$

but also smaller K_α , due to decay into an excited state of ^{223}Ra .

These are followed by series of γ 's

$$Q = m_{\text{nuc}}(X) - m_{\text{nuc}}(Y^*) - m_{\text{nuc}}(\alpha)$$



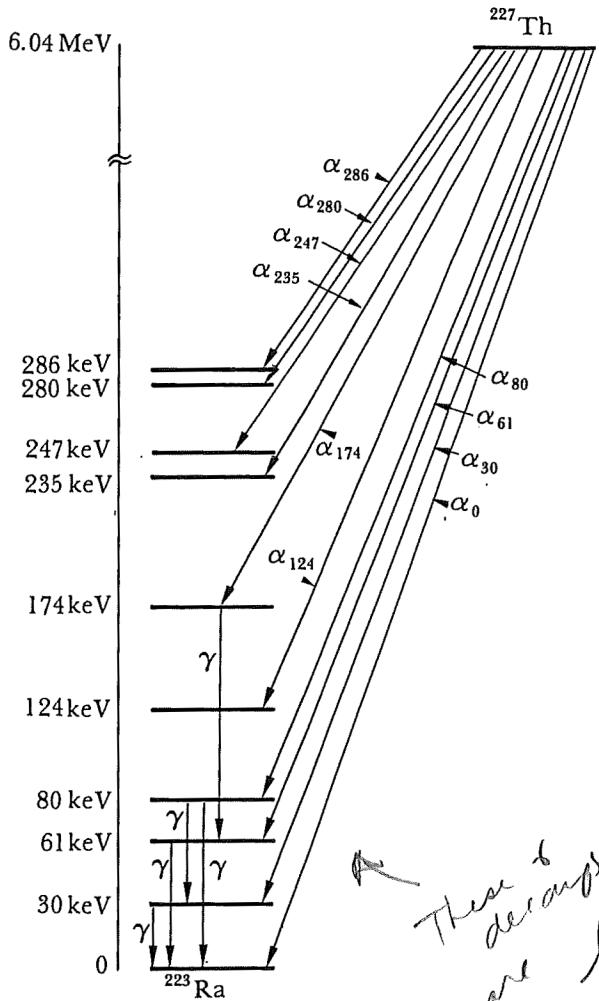
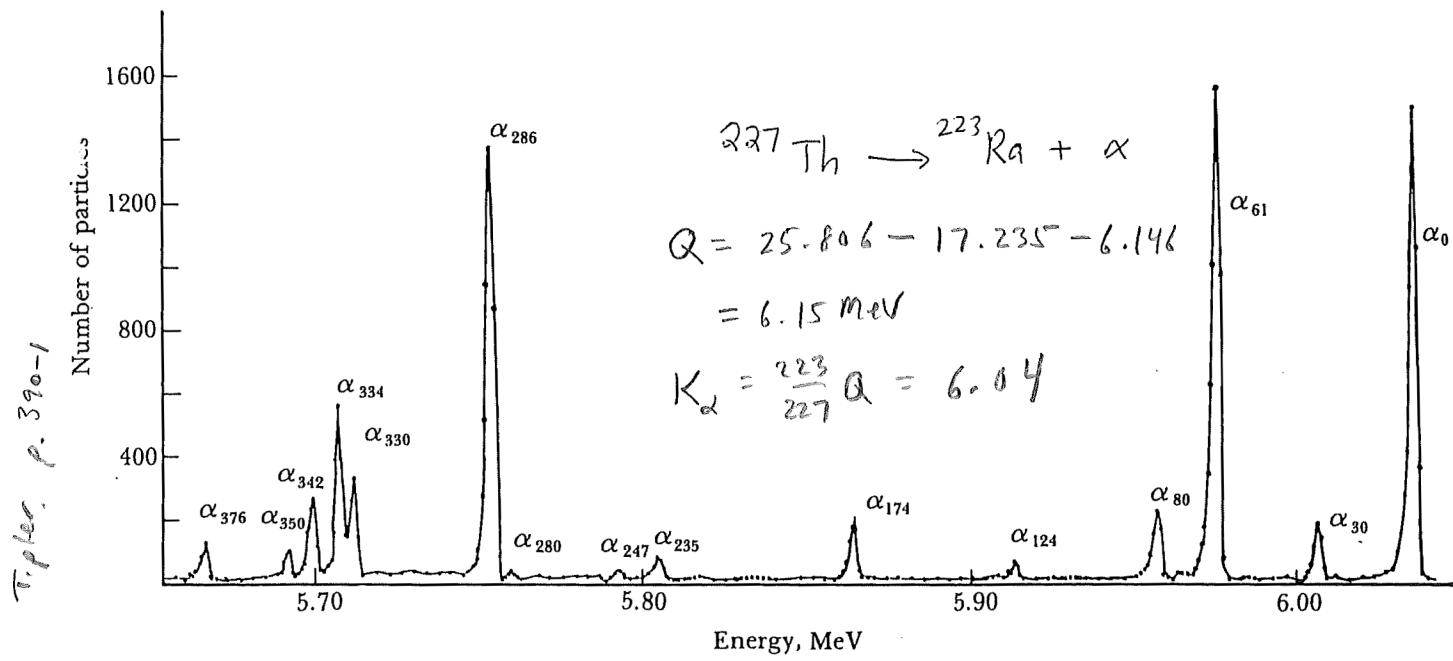
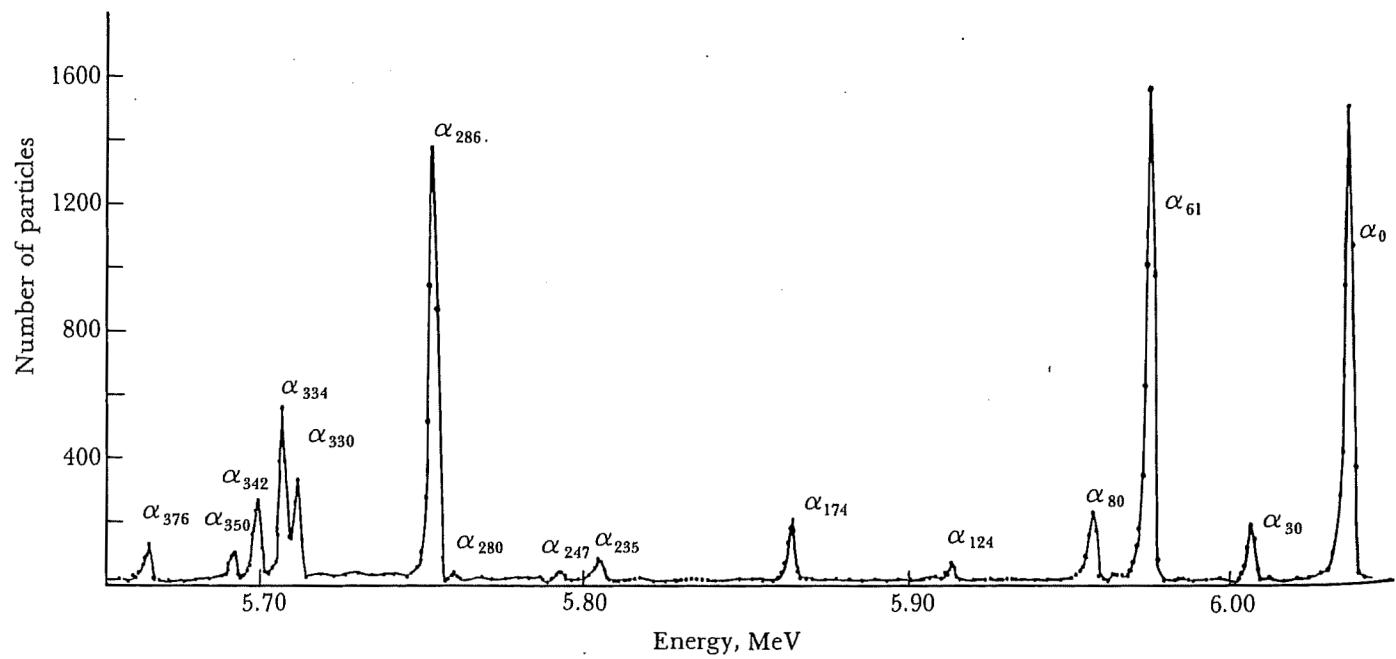
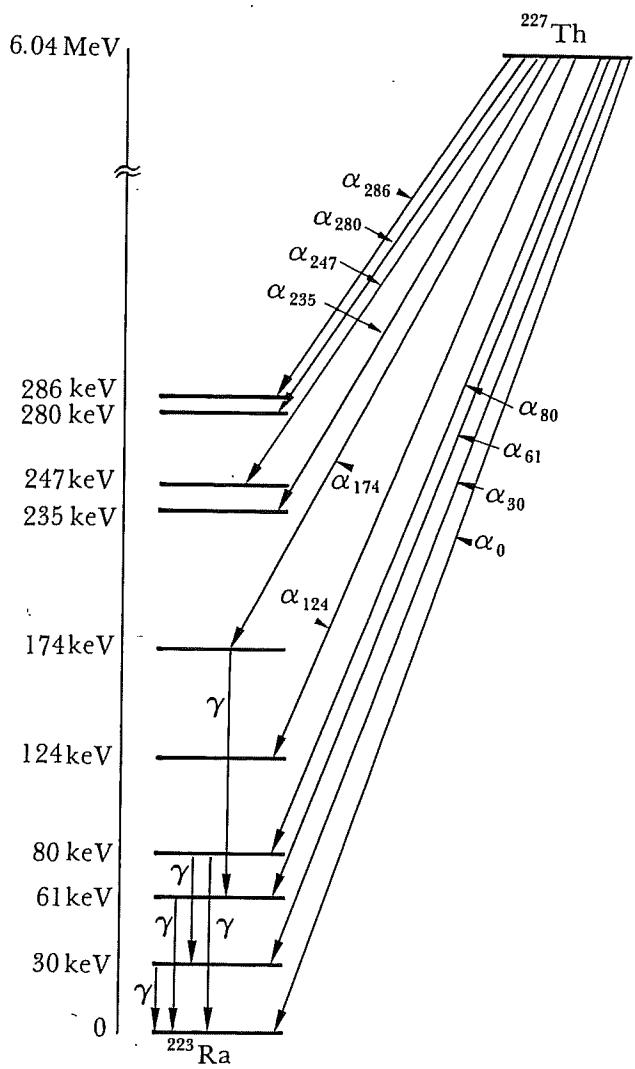


Figure 11-13
Energy levels of ^{223}Ra determined by measurement of α -particle energies from ^{227}Th , as shown in Figure 11-12. Only the lowest-lying levels and some of the γ -ray transitions are shown.

widths due to experimental resolution, not to finite width due to finite lifetime
 These levels are believed to be in <
 $b_{1/2}$ and $1/n$ (see Table 11-18)





Half-life of α -emitters depends strongly on Q

	<u>Q</u>	<u>$T_{1/2}$</u>
^{227}U	7.2 MeV	1 minute
^{229}U	6.5 MeV	58 minutes
^{230}U	6.0 MeV	21 days
^{232}U	5.4 MeV	69 yrs
^{233}U	4.9 MeV	1.6×10^5 yrs
^{235}U	4.7 MeV	7.0×10^8 years
^{238}U	4.3 MeV	4.5×10^9 yrs
		99.3%

As $Q \downarrow$, $T_{1/2} \uparrow$

[see plot]

1911 Geiger-Nuttall rule : $\ln T_{1/2} = 9 \frac{Z}{\sqrt{Q}} + c_2$

(c_2 is constant)

Recall $T_{1/2} = (\ln 2) R = \frac{\ln 2}{R}$

(from NW) $R = \frac{\text{prob. of decay}}{\text{time}}$ ← inherently quantum mechanical
(what is the process?)

How do we calculate the probability?



Imagine a uranium nucleus in an α particle trap and emit a thorium nucleus

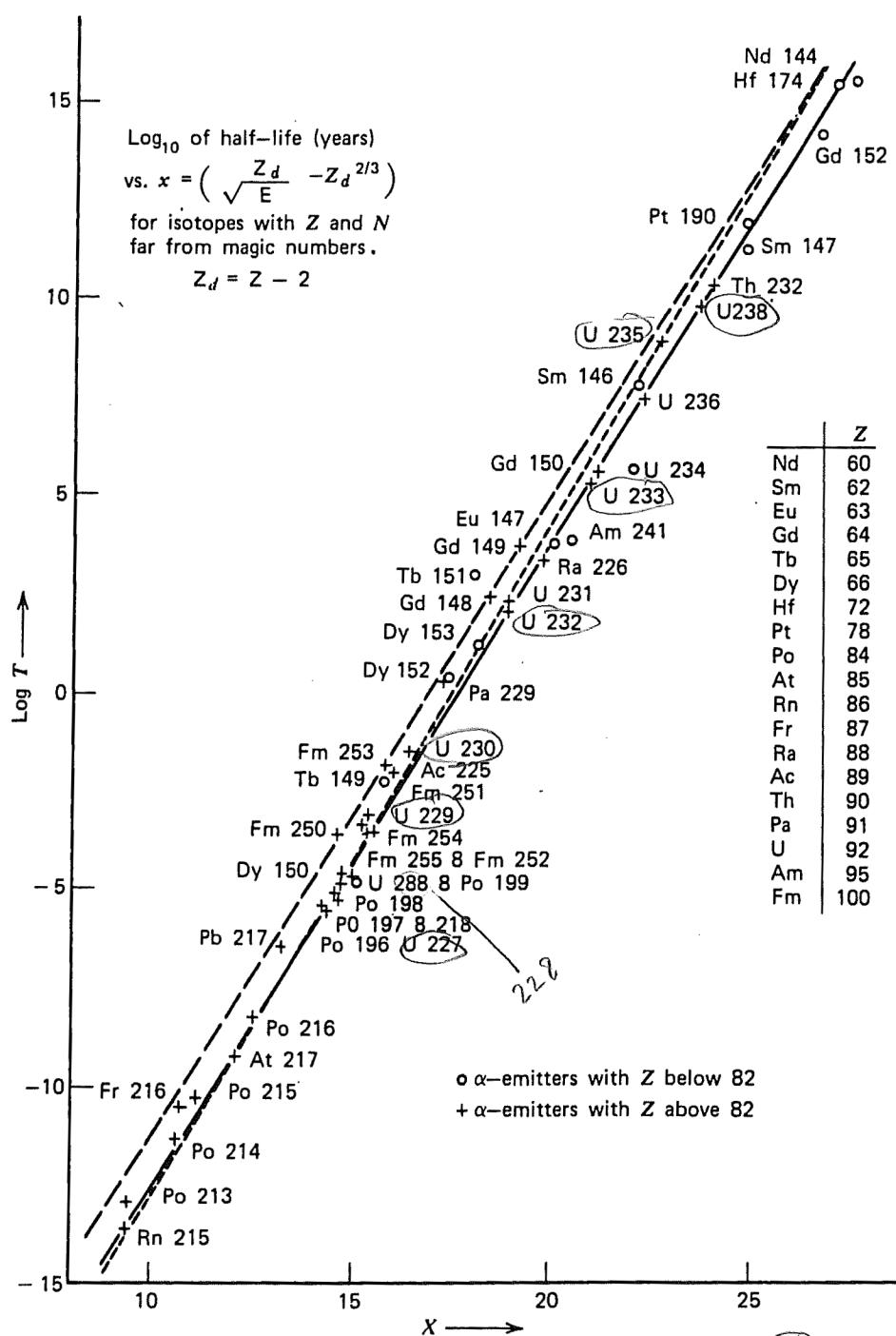


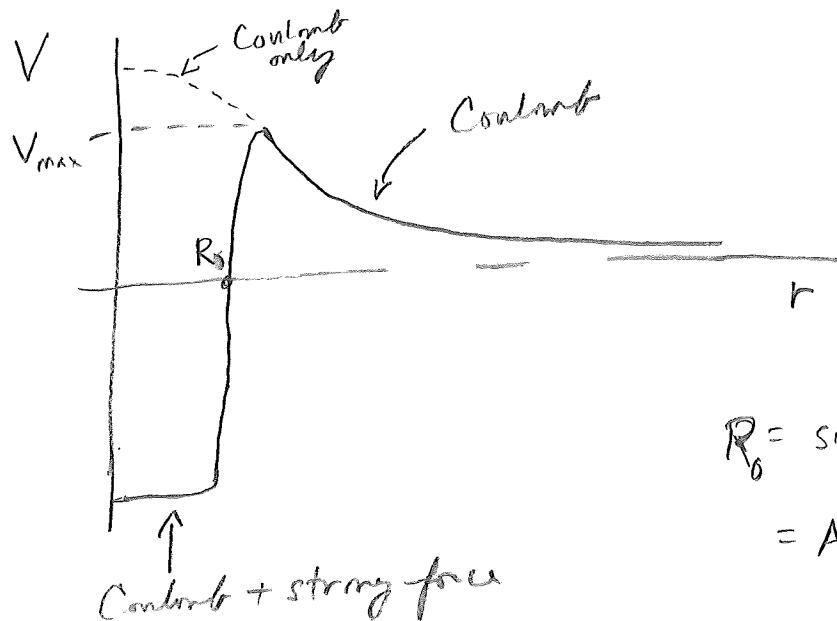
Figure 5-9. Plot of $\log_{10} 1/\tau$ versus $C_2 - C_1 Z_1 / \sqrt{E}$ with $C_1 = 1.61$ and a slowly varying $C_2 = 28.9 + 1.6Z_1^{2/3}$. (From E. K. Hyde, I. Perlman, and G. T. Seaborg, *The Nuclear Properties of the Heavy Elements*, Vol. 1, Prentice-Hall, Englewood Cliffs, N.J. (1964), reprinted by permission.)

$$\ln \frac{1}{\tau} \sim -3.71$$



α -5

Consider the potential energy between Y and α as a function of separation r



$$R_0 = \text{size of } Y \text{ nucleus} \\ = A^{1/3} r_0$$

Rough estimate of height of barrier

$$V_{\max} \sim \frac{K q_1 q_2}{R} = \frac{K(Ze)(2e)}{A^{1/3} r_0} = 2 \left(\frac{Ke^2}{r_0} \right) \frac{Z}{A^{1/3}}$$

$$\sim 2 \frac{(1.44 \text{ MeV fm})}{(1.2 \text{ fm})} \frac{Z}{A^{1/3}}$$

$$\therefore \sim (2.4 \text{ MeV}) \frac{Z}{A^{1/3}}$$

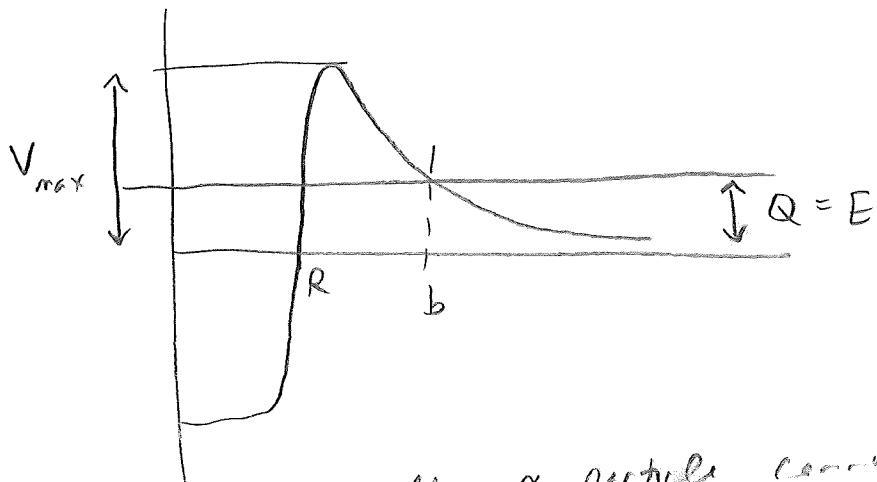
for ${}^{238}\text{U}$ decay, $Z = 90$, $A = 238 \rightarrow V_{\max} \sim 35 \text{ MeV}$

[Actual $V_{\max} \sim 25 \text{ MeV}$]

q-6

What is energy of α particle?

$$\text{After it escapes, } E = T + V \\ \approx Q + 0$$



Classically, α particle cannot escape
Quantum mechanically, it can "tunnel" through barrier

[Gamow (1928)
Gurney-Condon]

Tunnelling rate

$$R = (\text{const}) e^{-\frac{Z}{\hbar} \int_R^b dr \sqrt{2m(V(r)-E)}}$$

$$R = (\text{const}) e^{-C_1 \frac{Z}{\sqrt{Q}}}$$

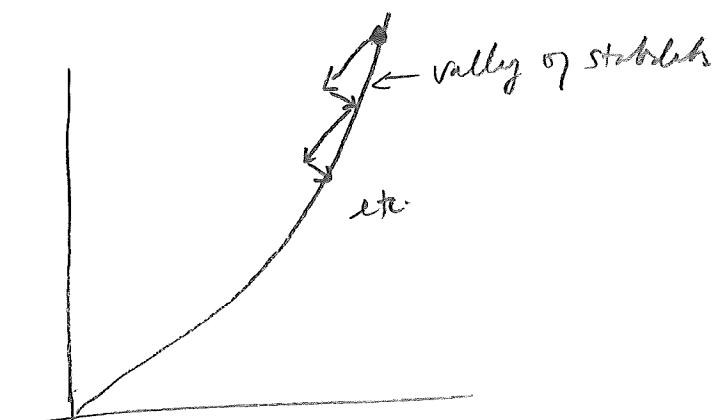
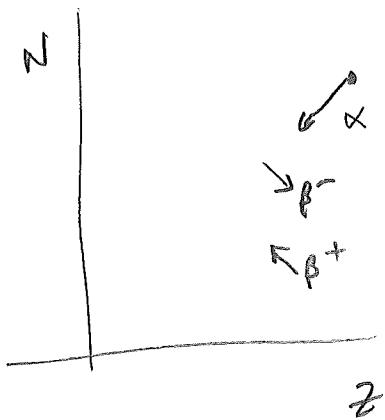
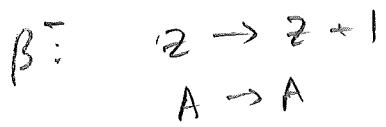
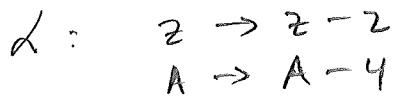
[HW] calculate this to find $R \approx e^{-C_1 \frac{Z}{\sqrt{Q}}}$

[HW: const = 3.96
expt. 3.68 (Gamow's)]

$$T_{1/2} \sim \frac{\ln 2}{R}$$

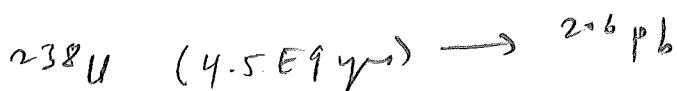
$$\ln T_{1/2} \sim -\ln R + \text{const}$$

$$\sim C_1 \frac{Z}{\sqrt{Q}} + \text{const} \quad (\text{Gurney-Nuttall})$$



Since α, β conserve $A \bmod 4$,
there are 4 distinct radioactive series

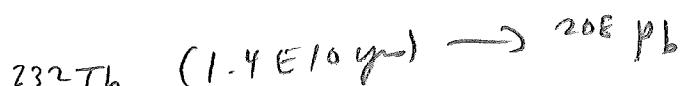
$$A = 4n + 2$$



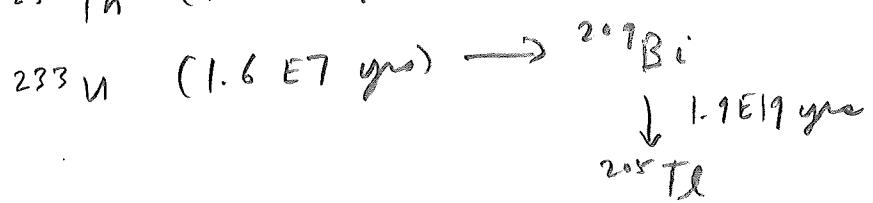
$$A = 4n + 3$$



$$A = 4n$$



$$A = 4n + 1$$



hyperphysics.phy-astr.
gsu.edu/hbase/nuclear

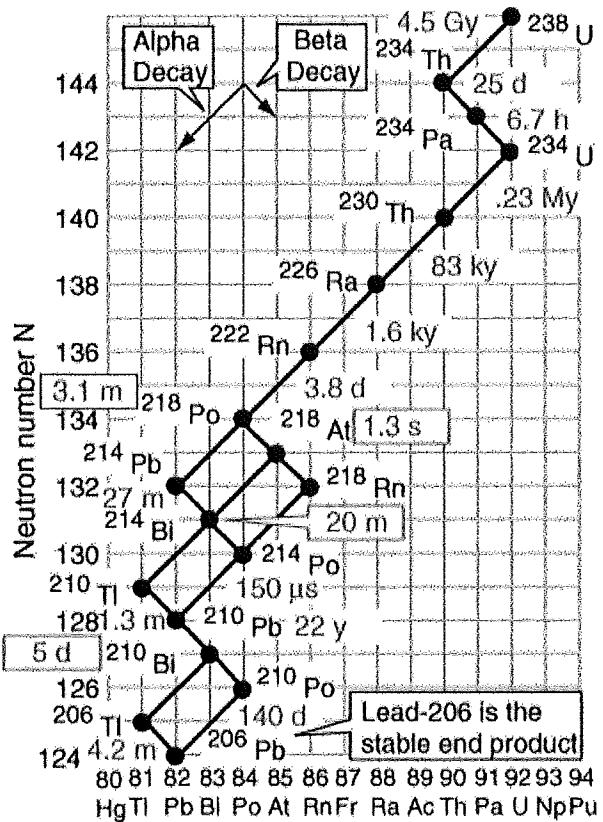
(georgia state u.)

The Uranium-238 Decay Series

- ^{235}U Series
- ^{232}Th Series
- ^{238}U Series
- ^{237}Np Series

The four natural radioactive series

Boxed values
for half-life are
for multiple
decay paths



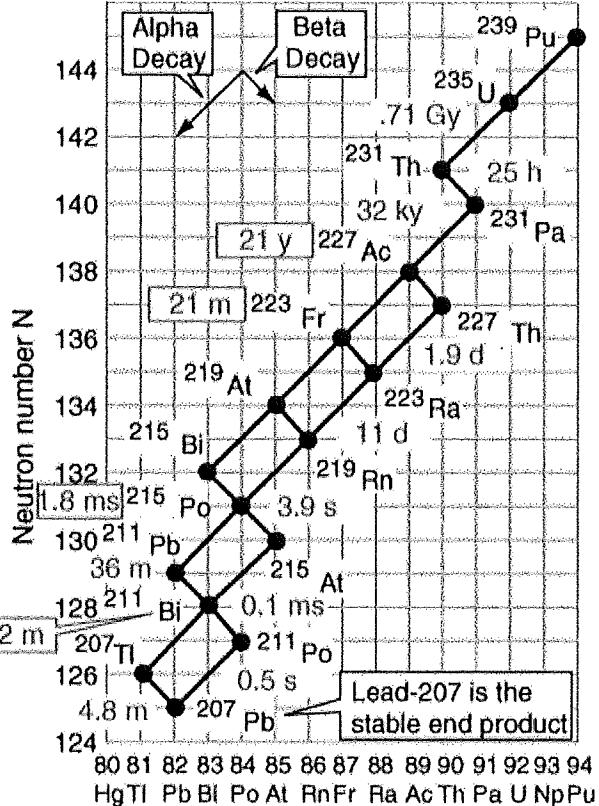
The Uranium-235 Decay Series

- ^{235}U Series
- ^{232}Th Series
- ^{238}U Series
- ^{237}Np Series

The four natural radioactive series

This series is traditionally called the Actinium series.

Boxed values for half-life are for multiple decay paths

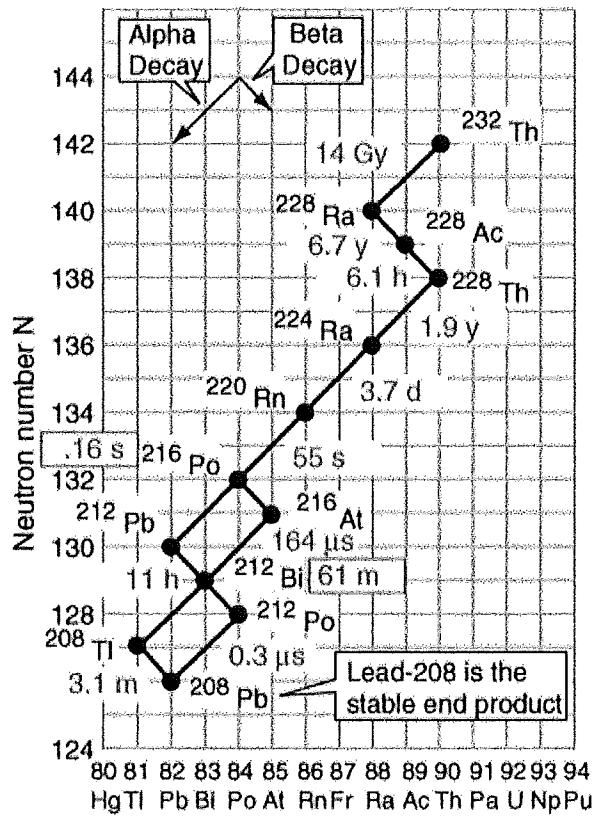


The Thorium-232 Decay Series

- ^{235}U Series
- ^{232}Th Series
- ^{238}U Series
- ^{237}Np Series

The four natural radioactive series

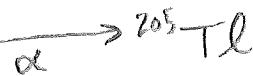
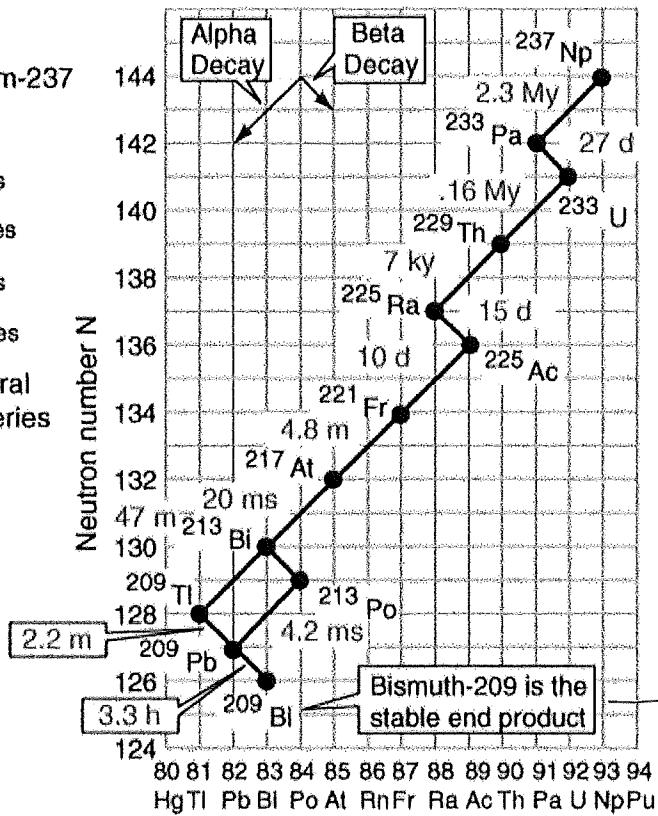
Boxed values for half-life are for multiple decay paths



The Neptunium-237 Decay Series

- ^{235}U Series
- ^{232}Th Series
- ^{238}U Series
- ^{237}Np Series

The four natural radioactive series



$$\tau_{1/2} = 1.9 \times 10^{11} \text{ years}$$

$$[2^{103}]$$

Try to estimate energy released in α -decay using semi-empirical
 $238\text{U} \rightarrow 234\text{Th} + \alpha$ (accuracy to about 15%)

$$\Delta = 47.304 \quad \Delta = 40.604 \quad \Delta = 2.425 \Rightarrow Q = 4.270 \text{ MeV}$$

6.695 MeV ← actual difference

use binding energy formula

$$B(238\text{U}) = 1778.63 \Rightarrow \Delta = 71.05$$

$$B(234\text{Th}) = 1756.03 \Rightarrow \Delta = 62.93$$

22.60

8.12 ← not too bad,
but not very good

{ not too close}

Now do Taylor expansion

$$B = c_1 A - c_2 A^{2/3} - c_3 \frac{Z^2}{A^{1/3}} - c_4 \frac{(A-2Z)^2}{A}$$

$$B\left(\frac{A}{Z}X\right) - B\left(\frac{A-Z}{Z}Y\right) = 4c_1 - \frac{8c_2}{3A^{1/3}} - c_3 \frac{Z^2}{A^{1/3}} \left[\frac{4}{Z} - \frac{4}{3A} \right]$$

$$+ 4c_4 \frac{(A-2Z)^2}{A^2}$$

$$= 62 - \frac{44.8}{A^{1/3}} - \underbrace{\frac{2.88 Z^2}{A^{1/3}}}_{0.464} \left(\frac{1}{Z} - \frac{1}{3A} \right) + \frac{92(A-2Z)^2}{A^2}$$

$$238\text{U}_{92} \rightarrow 62 - 7.23 - 37.25 + 4.74 = 22.3$$

$$\text{Since } \Delta = Zm_H + (A-Z)m_n - A(931.5) - B$$

$$\Delta(X) - \Delta(Y) = \underbrace{Zm_H + Zm_n - 4(931.5)}_{30.73} - \Delta B$$

$$= -31.3 - \frac{44.8}{A^{1/3}} - 2.08 \frac{Z^2}{A^{1/3}} \left(\frac{1}{Z} - \frac{1}{3A} \right) + 92 \frac{(A-2Z)^2}{A^2}$$

$$= -31.3 + 7.2 + 37.2 - 4.7 = 8.4$$