

In QM, every observable A (something that can be measured) is associated with a linear operator $\hat{A}$

which acts on wavefunctions ψ to give another function

eg. energy $E \rightarrow \hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$

position $X \rightarrow \hat{X} = x$ [multiply by x]

momentum $P \rightarrow \hat{P} = ?$ [T.B.D.]

- An eigenfunction of $\hat{A}$ is a wavefunction ψ_a that satisfies the eigenvalue equation

$$\hat{A} \psi_a = a \psi_a$$

- ψ_a represents a state w/ well-defined values of observable A .
ie if we measure A of state ψ_a , result will always be the eigenvalue a

- State ψ_a has zero uncertainty for the observable A

ie $\Delta A = 0$

proof: $\langle A \rangle = \int dx \psi_a^* \hat{A} \psi_a = a \int dx \psi_a^* \psi_a = a$

$$\langle A^2 \rangle = \int dx \psi_a^* \hat{A}^2 \psi_a = a^2$$

$\underbrace{\hat{A} \psi_a}_{a \psi_a}$
 $\underbrace{\psi_a^* a \psi_a}_{a^2 \psi_a^* \psi_a}$

$$\Delta A = \sqrt{\langle A^2 \rangle - \langle A \rangle^2} = \sqrt{a^2 - a^2} = 0$$

[can talk about compatibility, incompatibility, & non commuting operators]

$$\hat{p} = ?$$

Eigenfunctions of $\hat{p}$ are states of well-defined momentum

$$\hat{p} \psi_p = p \psi_p$$

For a free particle, a state of well-defined momentum & energy is

$$\psi = e^{i \frac{px}{\hbar}} e^{-i \frac{Et}{\hbar}}$$

$$i \hbar \frac{\partial \psi}{\partial t} = \hat{H} \psi$$

$$\underbrace{-i \frac{E}{\hbar} \psi}_{E \psi}$$

$$\hat{p} e^{i \frac{px}{\hbar}} = p e^{i \frac{px}{\hbar}}$$

$$\hat{p} = B \frac{d}{dx} \Rightarrow \hat{p} e^{i \frac{px}{\hbar}} = B \frac{ip}{\hbar} e^{i \frac{px}{\hbar}} = p e^{i \frac{px}{\hbar}}, \quad B = \frac{\hbar}{i}$$

$$\boxed{\hat{p} = \frac{\hbar}{i} \frac{d}{dx}}$$

Even though we "derived" $\hat{p}$ using a free particle, the result is valid for any particle

$$\langle p \rangle = \int dx \psi^* \hat{p} \psi = \frac{\hbar}{i} \int dx \psi^* \frac{d\psi}{dx}$$

[exercise: calc $\langle p \rangle$, Δp for particle in a box]

$$\langle p \rangle = 0, \quad \Delta p = \frac{\hbar \pi n}{L}$$

[After exercise presented]

Energy eigenfunctions for particle in a box

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

$$E_n = \frac{\frac{1}{2} n^2 \hbar^2}{2mL^2}$$

$$\Delta E_n = 0$$

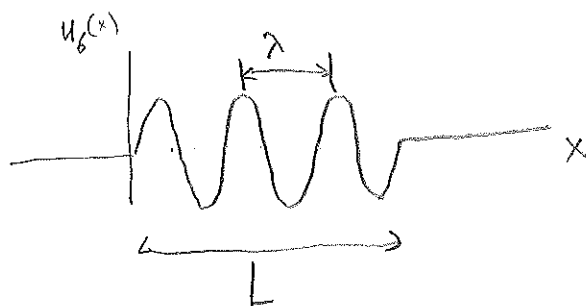
$$\langle x \rangle = \frac{L}{2}$$

$$\Delta x = L \sqrt{\frac{1}{12} - \frac{1}{2n^2\pi^2}} \quad [\text{from problem}]$$

$$\langle p \rangle = 0$$

$$\Delta p = \frac{\hbar n \pi}{L} \quad [\text{from exercise}]$$

[neither position nor momentum eigenstate]



$$\lambda = \frac{2L}{n} \quad \text{from inspection}$$

$$p = \frac{h}{\lambda} = \frac{hn}{2L} = \frac{\pi \hbar n}{L}$$

would "expect" momentum measurement to give $\approx \frac{\pi \hbar n}{L}$

Why $\langle p \rangle = 0$?

$$p = \pm \frac{\pi \hbar n}{L}$$

[particle can move either direction]

A measurement of p gives (see plot)

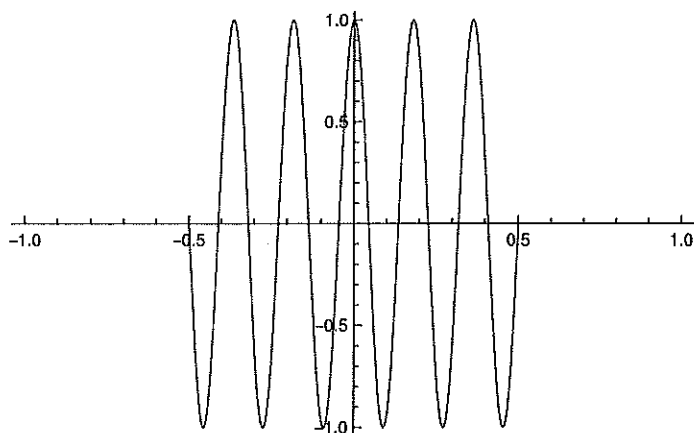
```
f[x_] := If[Abs[x] < 1/2, Cos[n Pi x], 0]
```

```
nn = 11;
```

```
FourierTransform[f[x], x, p]
```

$$\frac{2 p \cos\left[\frac{n \pi}{2}\right] \sin\left[\frac{p}{2}\right] - 2 n \pi \cos\left[\frac{p}{2}\right] \sin\left[\frac{n \pi}{2}\right]}{\sqrt{2 \pi} \left(p^2 - n^2 \pi^2\right)}$$

```
Plot[f[x] /. n -> nn, {x, -1, 1}]
```



$$n = 11$$

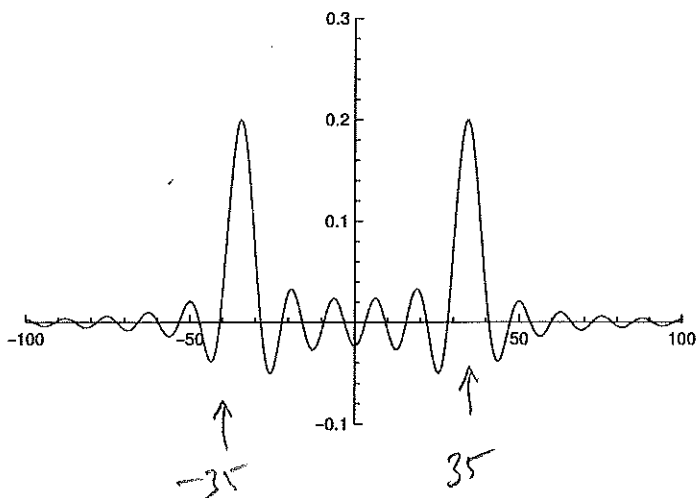
$$\lambda = \frac{2}{11} L$$

```
FourierTransform[f[x] /. n -> nn, x, p]
```

$$\frac{11 \sqrt{2 \pi} \cos\left[\frac{p}{2}\right]}{p^2 - 121 \pi^2}$$

$$p = \frac{h}{\lambda} = \frac{11 \pi h}{L}$$

```
Plot[%, {p, -100, 100}, PlotRange -> {{-100, 100}, {-0.1, 0.3}}]
```



$$\text{let } h = L = 1$$

$$p = 11 \pi \approx 35$$

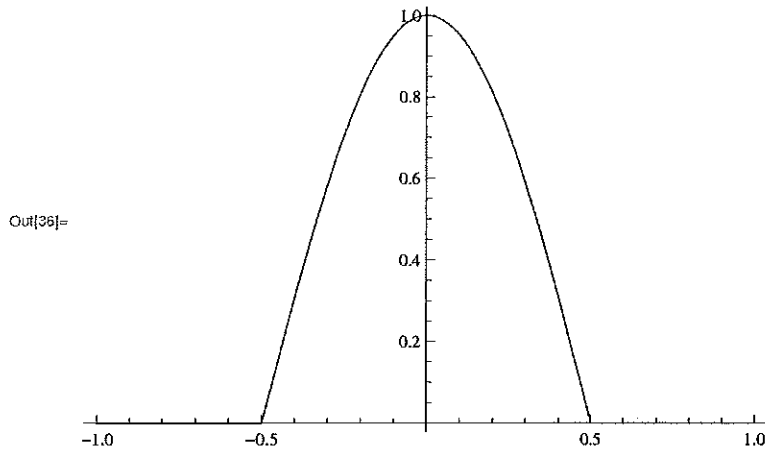
rest of these plots
just for my reference

In[34]:= **f[x_] := If[Abs[x] < 1/2, Cos[n Pi x], 0]**

In[35]:= **FourierTransform[f[x], x, p]**

$$\text{Out[35]} = \frac{2 p \cos\left[\frac{n \pi}{2}\right] \sin\left[\frac{p}{2}\right] - 2 n \pi \cos\left[\frac{p}{2}\right] \sin\left[\frac{n \pi}{2}\right]}{\sqrt{2} \pi (p^2 - n^2 \pi^2)}$$

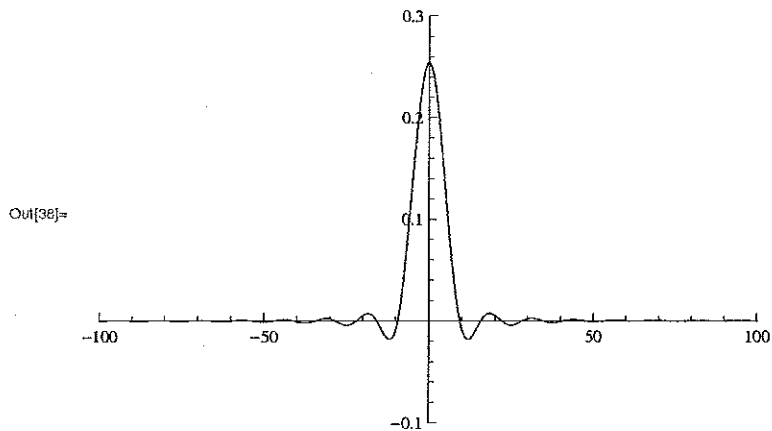
In[36]:= **Plot[f[x] /. n -> 1, {x, -1, 1}]**



In[37]:= **FourierTransform[f[x] /. n -> 1, x, p]**

$$\text{Out[37]} = \frac{\sqrt{2} \pi \cos\left[\frac{p}{2}\right]}{-p^2 + \pi^2}$$

In[38]:= **Plot[%, {p, -100, 100}, PlotRange -> {{-100, 100}, {- .1, .3}}]**



In[41]:= **f[x_] := If[Abs[x] < 1/2, Cos[n Pi x], 0]**

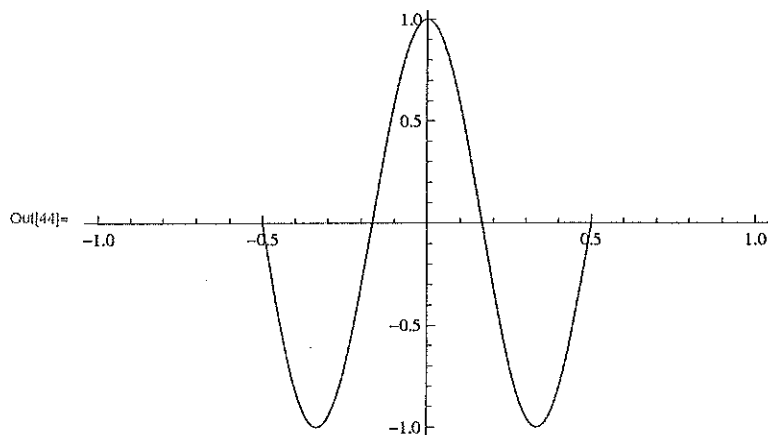
In[42]:= **nn = 3**

Out[42]:= 3

In[43]:= **FourierTransform[f[x], x, p]**

Out[43]:=
$$\frac{2 p \cos\left[\frac{n \pi}{2}\right] \sin\left[\frac{p}{2}\right] - 2 n \pi \cos\left[\frac{p}{2}\right] \sin\left[\frac{n \pi}{2}\right]}{\sqrt{2} \pi (p^2 - n^2 \pi^2)}$$

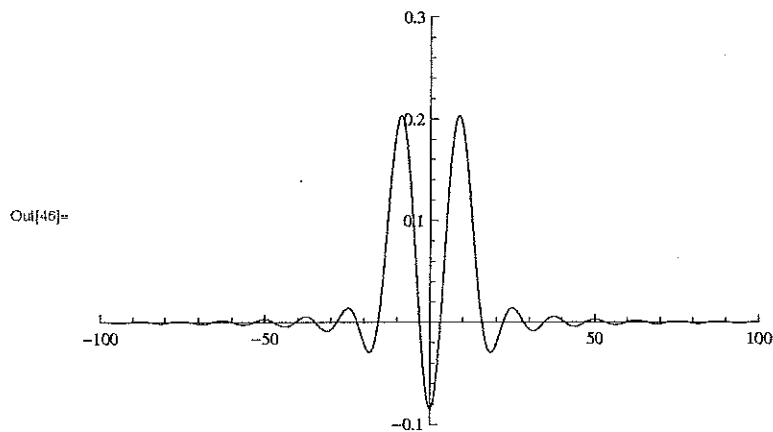
In[44]:= **Plot[f[x] /. n -> nn, {x, -1, 1}]**



In[45]:= **FourierTransform[f[x] /. n -> nn, x, p]**

Out[45]=
$$\frac{3 \sqrt{2} \pi \cos\left[\frac{p}{2}\right]}{p^2 - 9 \pi^2}$$

In[46]:= **Plot[%, {p, -100, 100}, PlotRange -> {{-100, 100}, {- .1, .3}}]**



In[47]:= **f[x_] := If[Abs[x] < 1/2, Cos[n Pi x], 0]**

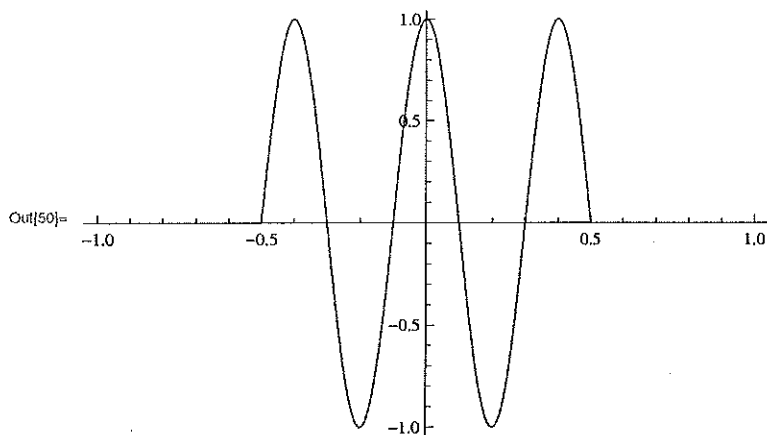
In[48]:= **nn = 5**

Out[48]= 5

In[49]:= **FourierTransform[f[x], x, p]**

Out[49]=
$$\frac{2 p \cos\left[\frac{n \pi}{2}\right] \sin\left[\frac{p}{2}\right] - 2 n \pi \cos\left[\frac{p}{2}\right] \sin\left[\frac{n \pi}{2}\right]}{\sqrt{2 \pi} \left(p^2 - n^2 \pi^2\right)}$$

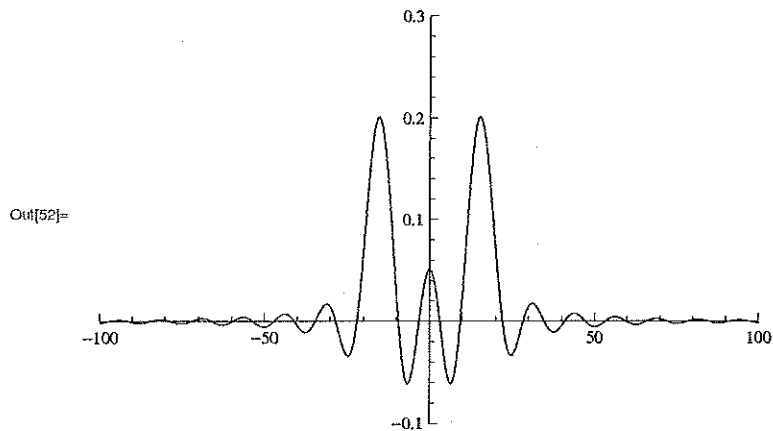
In[50]:= **Plot[f[x] /. n -> nn, {x, -1, 1}]**



In[51]:= **FourierTransform[f[x] /. n -> nn, x, p]**

Out[51]=
$$-\frac{5 \sqrt{2 \pi} \cos\left[\frac{p}{2}\right]}{p^2 - 25 \pi^2}$$

In[52]:= **Plot[%, {p, -100, 100}, PlotRange -> {{-100, 100}, {- .1, .3}}]**



```
In[1]:= f[x_] := If[Abs[x] < 1/2, Cos[n Pi x], 0]
```

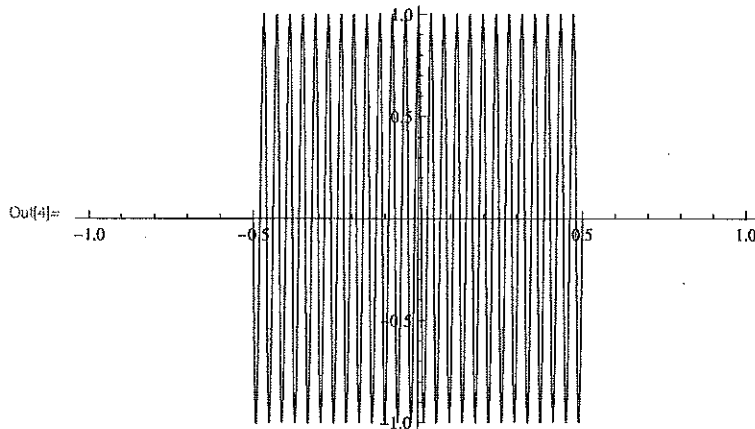
```
In[2]:= nn = 51
```

```
Out[2]:= 51
```

```
In[3]:= FourierTransform[f[x], x, p]
```

$$\text{Out[3]} = \frac{2 p \cos\left[\frac{n \pi}{2}\right] \sin\left[\frac{p}{2}\right] - 2 n \pi \cos\left[\frac{p}{2}\right] \sin\left[\frac{n \pi}{2}\right]}{\sqrt{2 \pi} (p^2 - n^2 \pi^2)}$$

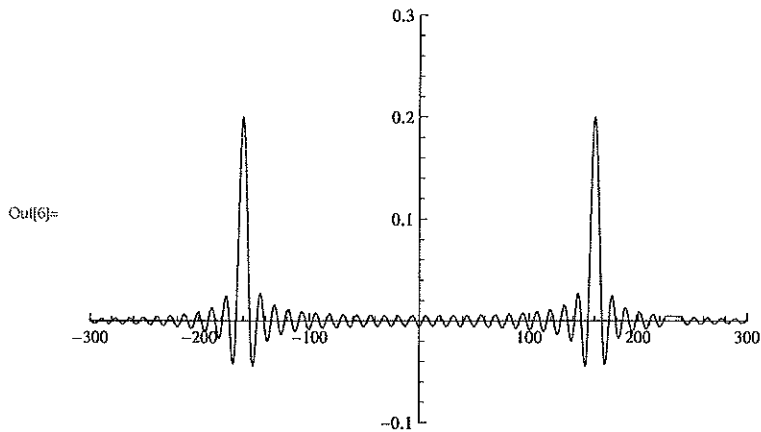
```
In[4]:= Plot[f[x] /. n -> nn, {x, -1, 1}]
```



```
In[5]:= FourierTransform[f[x] /. n -> nn, x, p]
```

$$\text{Out[5]} = \frac{51 \sqrt{2 \pi} \cos\left[\frac{p}{2}\right]}{p^2 - 2601 \pi^2}$$

```
In[6]:= Plot[%, {p, -300, 300}, PlotRange -> {{-300, 300}, {-0.1, 0.3}}]
```



Note $\Delta x \sim L$ $\Delta p \sim \frac{1}{L}$

$$\Delta x \Delta p = \hbar \pi n \sqrt{\frac{1}{12} - \frac{1}{2\pi^2 n^2}}$$

$$= \hbar \sqrt{\frac{\pi^2 n^2}{12} - \frac{1}{2}} \quad n=1, 2, \dots$$

Minimum value for $n=1$

$$(\Delta x \Delta p)_{\min} = \hbar \sqrt{\frac{\pi^2}{12} - \frac{1}{2}} \approx 0.57 \hbar$$

In general one can prove

$$\Delta x \Delta p \geq \frac{\hbar}{2} \quad (\text{Heisenberg})$$

If $\Delta x \Delta p = \frac{\hbar}{2}$ we say inequality is saturated

Consider a free particle ($V(x) = 0$)

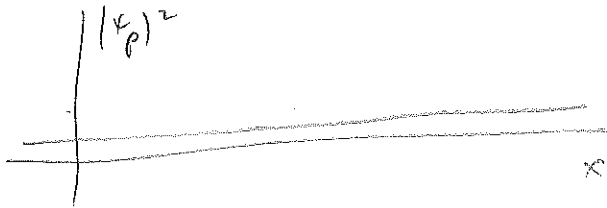
described by a momentum & energy eigenstate

$$\psi_p(x, t) = A(p) e^{i \frac{p x}{\hbar}} e^{-i \frac{E_p t}{\hbar}} \quad E_p = \frac{p^2}{2m}$$

If measure its momentum or energy, the result will be p or E_p respectively w/ 100% probability.

What if measure its position?

Probability density $|\psi_p|^2 = |A(p)|^2$ (indep. of x and t)



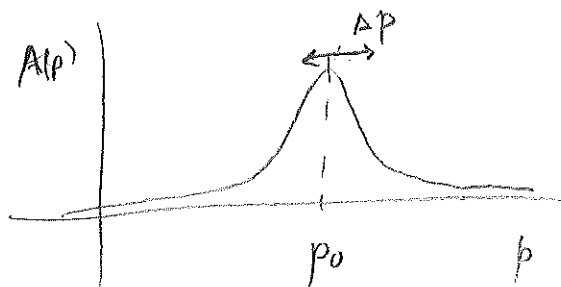
Problems w/ this description

- ① stationary state $\Rightarrow$ does not move
- ② not localized $\Rightarrow$ equally likely to be found anywhere
- ③ not normalizable $\int_{-\infty}^{\infty} |A(p)|^2 dx = \infty$

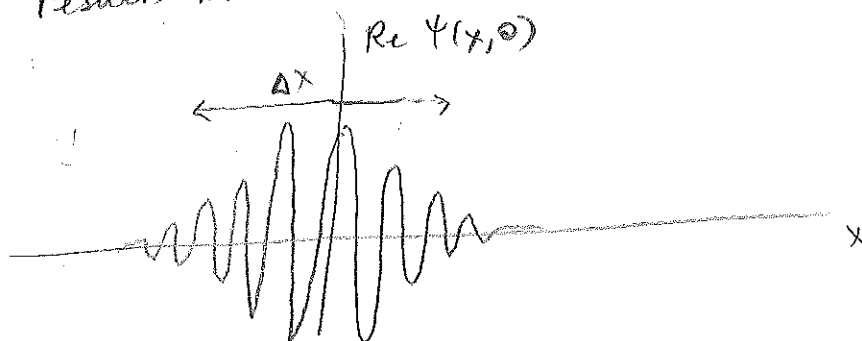
All 3 problems can be solved by forming
a wave packet (linear combination of eigenstates)

$$\begin{aligned}\psi(x, t) &= \int dp \psi_p(x, t) \\ &= \int dp A(p) e^{i \frac{px}{\hbar}} e^{-i \frac{p^2 t}{2m\hbar}}\end{aligned}$$

where $A(p)$ is a function peaked near p_0 w/ width $\sim \Delta p$



Constructive interference between different components,
results in



The wavepacket

① moves (w/ speed $\frac{p_0}{m}$)

② is localized to within $\Delta x \sim \frac{\hbar}{\Delta p}$

③ can be normalized.

