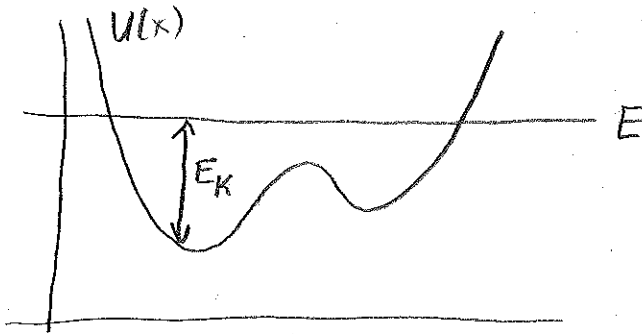


Schrödinger equation

An electron experiencing a force
(e.g. due to the proton in a hydrogen atom,
or due to the conductor in which it is confined)
has potential energy U as well as kinetic energy E_K



Total mechanical energy is $E = E_K + U = \text{constant}$

The matter wave ψ describing the electron
obeys the Schrödinger equation

$$\left[-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + U(x)\psi = E\psi \right]$$

valid for an electron with well-defined energy

Let's check that free electron obeys Schrodinger eqn.

Free electron is described by a travelling matter wave

$$\psi = A \cos(kx - \omega t) + i A \sin(kx - \omega t) \quad \text{where } k = \frac{2\pi}{\lambda}$$

$$\text{and } \lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE_k}} \Rightarrow E_k = \frac{h^2}{2m\lambda^2}$$

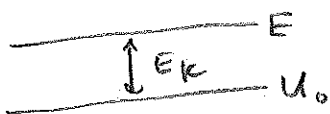
Take derivatives

$$\frac{d\psi}{dx} = -Ak \sin(kx - \omega t) + iAk \cos(kx - \omega t)$$

$$\frac{d^2\psi}{dx^2} = -Ak^2 \cos(kx - \omega t) - iAk^2 \sin(kx - \omega t)$$

$$= -k^2 \psi$$

A free electron experiences no force, so $U = U_0 = \text{const}$



L.h.s of Schrodinger eqn

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + U\psi = \frac{\hbar^2 k^2}{2m} \psi + U_0 \psi = (E_k + U_0) \psi = E \psi$$

$$\text{But } \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2}{2m} \left(\frac{2\pi}{\lambda}\right)^2 = \frac{(2\pi\hbar)^2}{2m\lambda^2} = \frac{h^2}{2m\lambda^2} = E_k$$

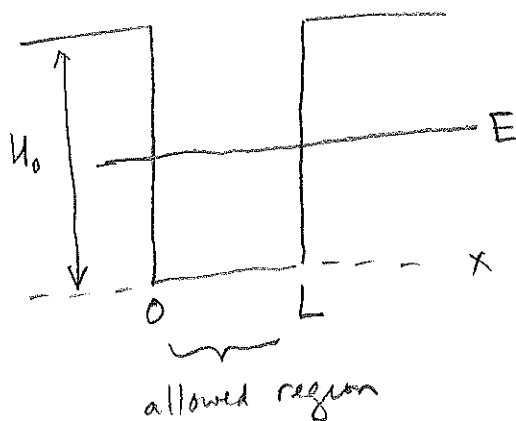
r.h.s of Schrodinger eqn

Particle in a box

Electron is confined to a region $0 < x < L$
by a strong force (eg. conduction electrons in a metal)

$$\text{Potential energy } U(x) = \begin{cases} 0 & 0 < x < L \\ U_0 & \text{outside box} \end{cases}$$

↑ very large



Schrodinger eqn inside box ($U=0$) is

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi$$

Verify that it has solution (similar to free electron)

$$\psi = N \sin(kx) \quad \leftarrow \text{purely real}$$

$$\frac{d\psi}{dx} = Nk \cos(kx)$$

$$\frac{d^2\psi}{dx^2} = -Nk^2 \sin(kx) = -k^2\psi$$

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = \frac{\hbar^2 k^2}{2m} \psi = \frac{\hbar^2}{2m\lambda^2} \psi = E\psi$$

As before $E = \frac{\hbar^2}{2m\lambda^2}$

but what is λ ?
Determined by boundary conditions.

If confining force is very strong, $U_0 \rightarrow \infty$ outside box
and there is zero probability to find electron there

$$\Rightarrow \psi = 0 \text{ for } x < 0 \text{ and } x > L$$

Thus

$$\psi = \begin{cases} 0 & x < 0 \\ N \sin kx & 0 < x < L \\ 0 & x > L \end{cases}$$

We require that ψ be continuous at the boundaries

ψ is automatically continuous at $x=0$ since $\sin(0)=0$.

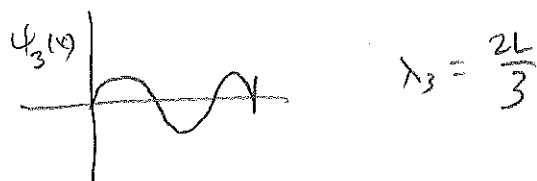
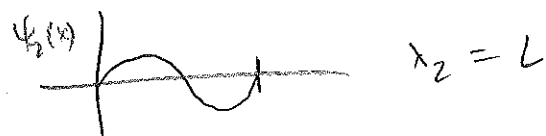
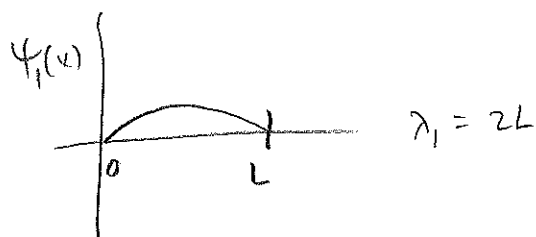
ψ is continuous at $x=L$ if $N \sin(kL) = 0$

only possible if $kL = n\pi$ for some integer n

Define $k_n = \frac{n\pi}{L} \Rightarrow \lambda_n = \frac{2\pi}{k_n} = \frac{2L}{n}$

De Broglie wavelength is quantized inside the box

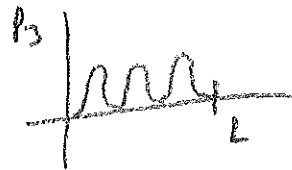
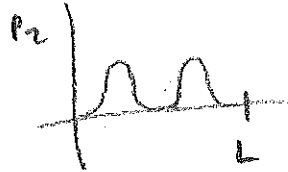
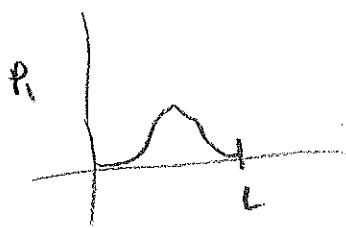
$$\psi_n(x) = N \sin(k_n x) = N \sin\left(\frac{n\pi x}{L}\right)$$



Probability density $P(x) = |\psi(x)|^2 = \psi_R^2 + \psi_I^2$

For particle in a box, ψ is purely real so $P(x) = \psi_R^2$

$$P_n(x) = \begin{cases} N^2 \sin^2\left(\frac{n\pi x}{L}\right), & \text{inside box} \\ 0 & \text{outside} \end{cases}$$



$P(x)dx = \text{prob. of finding particle between } x \text{ and } x+dx$

$\int_a^b P(x)dx = \text{prob. of finding particle between } x=a \text{ and } x=b$

$\int_{-\infty}^{\infty} P(x)dx = \text{prob. of finding particle anywhere} = 1$

$$\Rightarrow \boxed{\int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1} \quad \text{normalization condition}$$

For particle in a box

$$\int_{-\infty}^{\infty} |\psi|^2 dx = \int_0^L N^2 \sin^2\left(\frac{n\pi x}{L}\right) dx = 1$$

use this to find the value of N

[problem]

$$E = \frac{\hbar^2}{2m\lambda^2} = \frac{\hbar^2}{2m} \left(\frac{n}{2L}\right)^2 = \left(\frac{\hbar^2}{8mL^2}\right) n^2$$

X6

Energy of electron increases $\propto n$

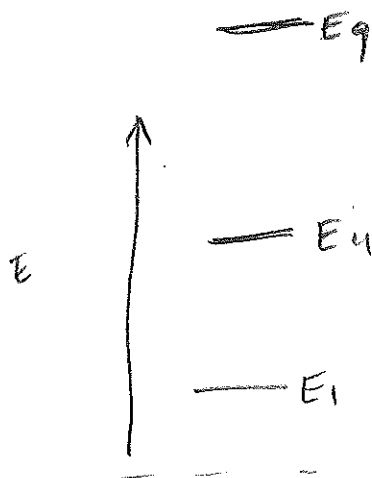
$$E_n = \frac{\hbar^2 k_n^2}{2m} = \left(\frac{\hbar^2 \pi^2}{2mL^2}\right) n^2$$

$$n=1 \Rightarrow E_1 = \frac{\hbar^2 \pi^2}{2mL^2} \quad (\text{ground state})$$

$$E_n = E_1 n^2$$

$$n=2 \quad E_2 = 4E_1 \quad (1^{\text{st}} \text{ excited state})$$

$$n=3 \quad E_3 = 9E_1 \quad (2^{\text{nd}} \text{ excited state})$$

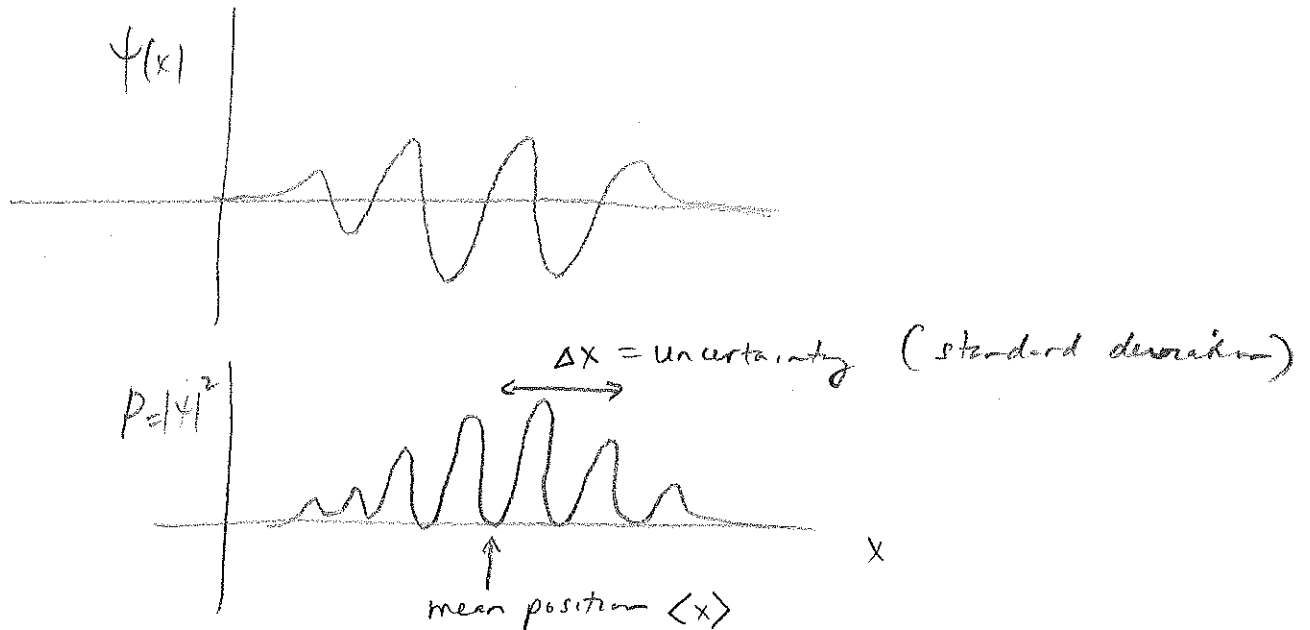


analogous to hydrogen atom;
also electron in a metal

$n=0$ not possible because $\psi=0 \Rightarrow$ no particle
 $n<0$ do not give separate solutions

When electron transitions from one energy level to another, it emits a photon whose energy is the difference between the energy levels

The position of particle is not necessarily well-defined



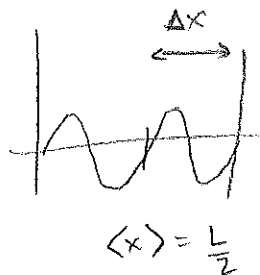
The momentum of a particle is also not necessarily well-defined (if not monochromatic)

$\Delta p = \text{uncertainty in momentum}$

Heisenberg established the uncertainty principle which holds for all particles

$$\Delta p \Delta x \geq \frac{\hbar}{2}$$

[Let's see how this works for particle in a box]



$$\Delta x \sim \frac{L}{2}$$

1

$$p = \frac{h}{\lambda} = \frac{hn}{2L} = \frac{\pi \hbar n}{L}$$

But can move left or right $\Rightarrow p_x = \pm p$

$$\Delta p \sim \frac{p - (-p)}{2} = p = \frac{\pi \hbar n}{L}$$

As $L \downarrow$, $\Delta x \downarrow$ but $\Delta p \uparrow$ so

product remains const

$$\Delta p \Delta x \sim \frac{\pi \hbar n}{2}$$

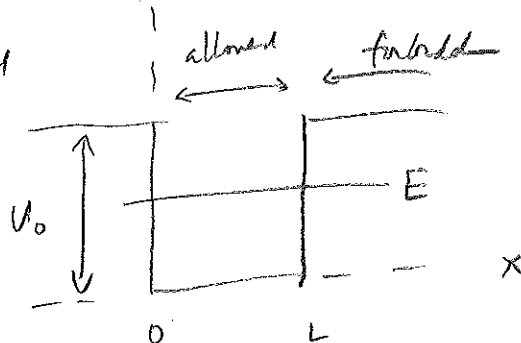
uncertainty relation
satisfied

$$\text{Since } n\pi > 1, \quad \Delta p \Delta x \geq \frac{\hbar}{2} \quad \nearrow$$

for $n \geq 1$

Consider a finite potential well

$$U(x) = \begin{cases} U_0 & x < 0 \\ 0 & 0 < x < L \\ U_0 & x > L \end{cases}$$



Schrödinger $-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + U(x)\psi = E\psi$

Inside the well $U(x)=0 \Rightarrow -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi$

has solution $\psi = N_1 \sin kx + N_2 \cos kx$ $\forall E = \frac{\hbar^2 k^2}{2m}$

Outside the well $U(x)=U_0 \Rightarrow -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = (E - U_0)\psi$

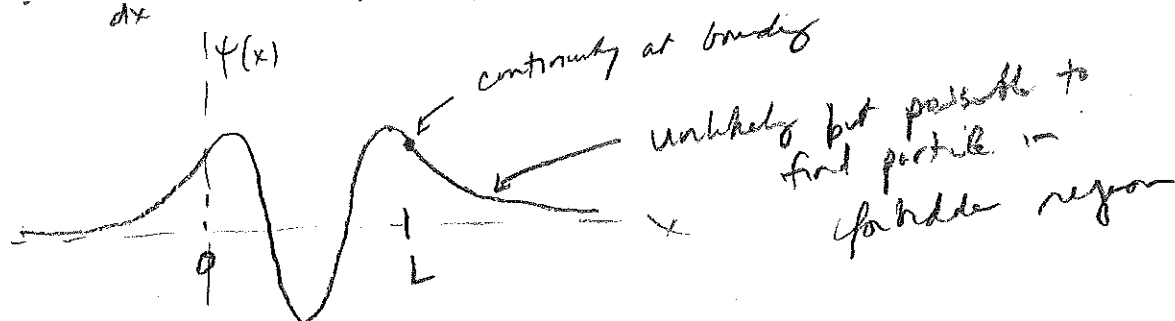
$\frac{d^2\psi}{dx^2} = \underbrace{\frac{2m}{\hbar^2} (U_0 - E)}_{\text{positive number, call it } \alpha^2} \psi$

$\Rightarrow \frac{d^2\psi}{dx^2} = \alpha^2 \psi$

Solution is not $\sin kx$ or $\cos kx$. Instead

$\psi = e^{+\alpha x}$ or $e^{-\alpha x}$

check: $\frac{d\psi}{dx} = \alpha e^{\alpha x}$, $\frac{d^2\psi}{dx^2} = \alpha^2 e^{\alpha x} = \alpha^2 \psi$ ✓



Talk about tunneling: $\Rightarrow$

- Scanning tunneling microscopy
- radioactive decay